

Counting $U(N)$ and $O(N)$ invariants and graph algebras of tensor models

Joseph Ben Geloun

LIPN, Paris Nord

Based on

JHEP **1711**, 092 (2017) and
Ann. Inst. HP D Comb. Phys. Interact. **1**, 77-138 (2014) with **S. Ramgoolam** (QMUL)

"Counting $O(N)$ invariants," work in Progress, with **R. C. Avouhou** (Jerusalem U & ICMPA, Benin), **N Dub** (LIPN)

November 1, 2018
Holographic Tensors,
OIST, Okinawa

Outline

- 1 Introduction: Permutation groups and Tensor models
- 2 Complex tensors and $U(N)$ invariants
- 3 Extending the results: Real tensors and $O(N)$ invariants

Outline

- 1 Introduction: Permutation groups and Tensor models
- 2 Complex tensors and $U(N)$ invariants
- 3 Extending the results: Real tensors and $O(N)$ invariants

Observables and Correlators in TM

- Key objects for tensor models: interactions/observables \equiv invariants of classical Lie groups, $U(N)$, $O(N)$.
- Invitation: new formulation of tensor model observables and correlators in terms of the symmetric group theoretic language and its representation theory.

Observables and Correlators in TM

- Key objects for tensor models: interactions/observables \equiv invariants of classical Lie groups, $U(N)$, $O(N)$.
- Invitation: new formulation of tensor model observables and correlators in terms of the symmetric group theoretic language and its representation theory.

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.
- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)
- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.

- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)

- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.
- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)
- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.
- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)
- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.
- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)
- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the Symmetric group and its representation theory

It allows us

- to embed our results in a different set theoretic formulation
- to gain confidence when implementing computations by software resources which could have been otherwise very difficult to handle.
- to shed a different light on our results
- to discover new and genuine effects
- to bridge theories and therefore discover new correspondences between countings, which from the outset look rather different (new bijections)
- to ask new questions
- to push towards new software experiments to guide our intuition

Switching to the symmetric group and its representation theory

• In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Switching to the symmetric group and its representation theory

- In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Switching to the symmetric group and its representation theory

• In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Switching to the symmetric group and its representation theory

• In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Switching to the symmetric group and its representation theory

• In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Switching to the symmetric group and its representation theory

- In practice, what symmetric group formulation has brought to the understanding of TM

⇒ **Mathematics/Computer Science side:**

→ Enumeration the observables/invariants (via computer softwares), generating functions are given (complex, and real)

→ Connection to topological field theory (TFT) with a geometrical interpretation of the counting (complex)

→ Understanding tensor invariants have a particularly rich algebra with interesting properties (matrix decomposition, bases, graduation)

→ Simplification and discovery new integer sequences [OEIS]

→ Discover computable sectors in TM

→ Establish a link with Computational Complexity Theory (via representation theoretic formulation of the counting): **Counting of tensor invariants is related to the problem of positivity of the Kronecker coefficient problem (NP vs P).**

→ Matrices & linguistic [S Ramgoolam et al].

⇒ **Physics side:**

→ For Matrix Models, much more [S Ramgoolam & R. de Mello Koch, Pablo Diaz,], from instance, understanding the half-BPS sector of $N = 4$ SYM.

→ Quantum information processing [Ramgoolam & Sedlak].

→ Highlight intriguing correspondences between countings in TM, Matrix Model (and therefore in String theory).

Goals

- Review some main results in **complex** TM using symmetric group formulation and its representation theory [work with Sanjaye, JHEP 2017; AIHP D 2014]
- Show preliminary results about the extension of these results to **real** TM, Carrozza-Tanasa-Klebanov-Tarnolposky. [wip with Avohou and Dub].

Outline

- 1 Introduction: Permutation groups and Tensor models
- 2 Complex tensors and $U(N)$ invariants
- 3 Extending the results: Real tensors and $O(N)$ invariants

Building blocks: Complex tensors and unitary invariants

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- Tensor contractions = unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

- An input: T is viewed as a $(d-1)$ -simplex. S_b^{int} “is” a gluing of simplexes and represents a d -polytope geometry
- In the following, all illustrations are made at fixed rank $d = 3$ but generally extends in any d .

Building blocks: Complex tensors and unitary invariants

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- Tensor contractions = unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

- An input: T is viewed as a $(d-1)$ -simplex. S_b^{int} “is” a gluing of simplexes and represents a d -polytope geometry
- In the following, all illustrations are made at fixed rank $d = 3$ but generally extends in any d .

Building blocks: Complex tensors and unitary invariants

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- Tensor contractions = unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

- An input: T is viewed as a $(d-1)$ -simplex. S_b^{int} “is” a gluing of simplexes and represents a d -polytope geometry
- In the following, all illustrations are made at fixed rank $d=3$ but generally extends in any d .

Building blocks: Complex tensors and unitary invariants

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- Tensor contractions = unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

- An input: T is viewed as a $(d-1)$ -simplex. S_b^{int} “is” a gluing of simplexes and represents a d -polytope geometry

- In the following, all illustrations are made at fixed rank $d = 3$ but generally extends in any d .

Building blocks: Complex tensors and unitary invariants

- A covariant complex tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} U_{p_1 q_1}^{(1)} \dots U_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad U^{(a)} \in U(N_a) \quad (1)$$

- Tensor contractions = unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T) \quad (2)$$

- An input: T is viewed as a $(d-1)$ -simplex. S_b^{int} “is” a gluing of simplexes and represents a d -polytope geometry
- In the following, all illustrations are made at fixed rank $d=3$ but generally extends in any d .

Unitary invariants

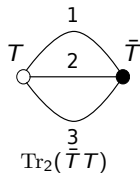
$$S_{\mathbf{b}}^{\text{int}}(T, \bar{T}) = \text{Tr}_{\mathbf{b}}(\bar{T} \cdot T \dots \bar{T} \cdot T)$$

- Coding unitary invariants: **b** bi-partite colored graphs
- Rank $D = 1$, Vectors: $\|\phi\| = \sum_a |\phi_a|^2$, 1 invariant.
- Rank $D = 2$, Matrices: $\text{Tr}[(M^{\dagger} M)^n]$, $\forall n \geq 1$, cyclic graphs (show on the board), $n \in \mathbb{N}$.

Unitary invariants

$$S_{\mathbf{b}}^{\text{int}}(T, \bar{T}) = \text{Tr}_{\mathbf{b}}(\bar{T} \cdot T \dots \bar{T} \cdot T)$$

- Coding unitary invariants: **b** bi-partite colored graphs

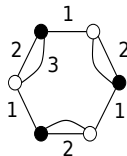
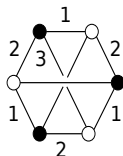
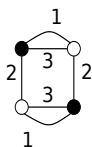
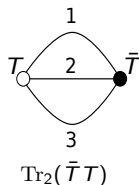


- Rank $D = 1$, Vectors: $\|\phi\| = \sum_a |\phi_a|^2$, 1 invariant.
- Rank $D = 2$, Matrices: $\text{Tr}[(M^\dagger M)^n]$, $\forall n \geq 1$, cyclic graphs (show on the board), $n \in \mathbb{N}$.

Unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T)$$

- Coding unitary invariants: **b** bi-partite colored graphs

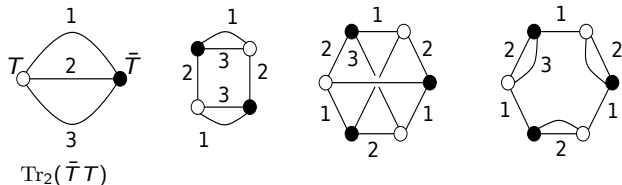


- Rank $D = 1$, Vectors: $\|\phi\| = \sum_a |\phi_a|^2$, 1 invariant.
- Rank $D = 2$, Matrices: $\text{Tr}[(M^\dagger M)^n]$, $\forall n \geq 1$, cyclic graphs (show on the board), $n \in \mathbb{N}$.

Unitary invariants

$$S_b^{\text{int}}(T, \bar{T}) = \text{Tr}_b(\bar{T} \cdot T \dots \bar{T} \cdot T)$$

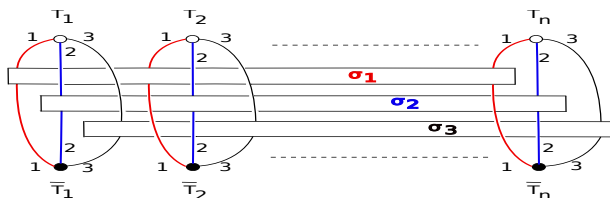
- Coding unitary invariants: **b** bi-partite colored graphs



- Rank $D = 1$, Vectors: $\|\phi\| = \sum_a |\phi_a|^2$, 1 invariant.
- Rank $D = 2$, Matrices: $\text{Tr}[(M^\dagger M)^n]$, $\forall n \geq 1$, cyclic graphs (show on the board), $n \in \mathbb{N}$.

Counting observables

[BG, Ramgoolam, AIHP D '14] Illustration in rank 3:



- Counting permutation triples $(\sigma_1, \sigma_2, \sigma_3) \in (S_n \times S_n \times S_n)$ up to the equivalence

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_2, \gamma_1 \sigma_3 \gamma_2), \quad \gamma_i \in S_n. \quad (3)$$

- Element of the double quotient $\text{Diag}(S_n) \backslash (S_n \times S_n \times S_n) / \text{Diag}(S_n)$.

Counting orbits: Apply Burnside's lemma

$$|H_1 \backslash G / H_2| = \frac{1}{|H_1||H_2|} \sum_{h_1 \in H_1} \sum_{h_2 \in H_2} \sum_{g \in G} \delta(h_1 g h_2 g^{-1}) \quad (4)$$

- Number of invariants

$$Z_3(n) = \frac{1}{(n!)^2} \sum_{\sigma_{1,2,3} \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1}) \quad (5)$$

→ Programming in Gap, and Mathematica [OEIS: A110143 (isomorphism of graph coverings)] Illustration at rank $d = 3$

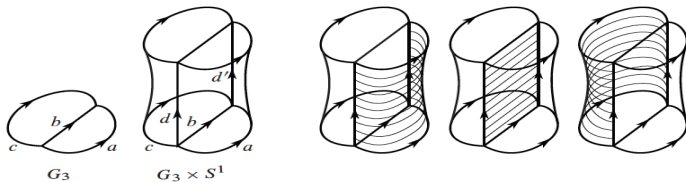
$$1; 4; 11; 43; 161; 901; 5579; 43206; 378360; 3742738, \dots \quad (6)$$

Topological Field Theory

→ TFT₂ on toroidal lattice

$$Z_3(n) = \frac{1}{(n!)^2} \sum_{\sigma_{1,2,3} \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1})$$

(7)



After some manipulations (gauge fixing one σ_i and introduce another variable), one arrives at

$$Z_3(n) = \frac{1}{n!} \sum_{\tau_0, \tau_1, \tau_2 \in S_n} \sum_{\gamma \in S_n} \delta(\gamma \tau_1 \gamma^{-1} \tau_1^{-1}) \delta(\gamma \tau_2 \gamma^{-1} \tau_2^{-1}) \delta(\gamma \tau_0 \gamma^{-1} \tau_0^{-1}) \delta(\tau_0 \tau_1 \tau_2) \quad (8)$$

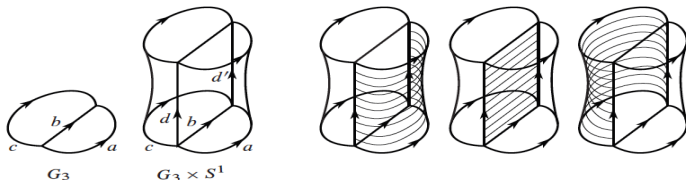
3 generators with a single relation, that is the fundamental group of S^2 with 3 punctures.

Topological Field Theory

→ TFT₂ on toroidal lattice

$$Z_3(n) = \frac{1}{(n!)^2} \sum_{\sigma_{1,2,3} \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1})$$

(7)



After some manipulations (gauge fixing one σ_i and introduce another variable), one arrives at

$$Z_3(n) = \frac{1}{n!} \sum_{\tau_0, \tau_1, \tau_2 \in S_n} \sum_{\gamma \in S_n} \delta(\gamma \tau_1 \gamma^{-1} \tau_1^{-1}) \delta(\gamma \tau_2 \gamma^{-1} \tau_2^{-1}) \delta(\gamma \tau_0 \gamma^{-1} \tau_0^{-1}) \delta(\tau_0 \tau_1 \tau_2) \quad (8)$$

3 generators with a single relation, that is the fundamental group of S^2 with 3 punctures.

Representation of the symmetric group: basics

- Irreps of symmetric group S_n are labelled by Young diagrams or $R \vdash n$ partition of n .

$$n = 7, \quad R = (1, 2, 4) = \begin{array}{|c|c|c|c|} \hline \square & \square & \square & \square \\ \hline \square & \square & & \\ \hline \square & & & \\ \hline \end{array} \quad (9)$$

- Expansion of the delta

$$\delta(\sigma) = \sum_R \frac{d(R)}{n!} \chi^R(\sigma), \quad d(R) = \frac{n!}{h(R)}, \quad h(R) = \prod_{i,j} \text{hook} - \text{Length}_{i,j}$$

$$\sum_{\sigma \in S_n} \chi^R(\sigma) \chi^S(\sigma) = n! \delta_{RS}; \quad \sum_{\gamma \in S_n} \delta(\gamma \sigma \gamma^{-1} \tau) = \sum_{R \vdash n} \chi^R(\sigma) \chi^R(\tau) \quad (10)$$

Revisiting the counting

- A small calculation

$$\begin{aligned}Z_3(n) &= \frac{1}{(n!)^2} \sum_{\sigma_i \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1}) \\ &= \frac{1}{(n!)^2} \sum_{\gamma_i \in S_n} \sum_{R_i \vdash n} \chi^{R_1}(\gamma_1) \chi^{R_1}(\gamma_2) \chi^{R_2}(\gamma_1) \chi^{R_2}(\gamma_2) \chi^{R_3}(\gamma_1) \chi^{R_3}(\gamma_2) \\ &= \sum_{R_1, R_2, R_3 \vdash n} (C(R_1, R_2, R_3))^2\end{aligned}\tag{11}$$

where the symbol

$$C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)\tag{12}$$

is the Kronecker coefficient.

→ The dimension of an algebra $\mathcal{K}(n)$.

Revisiting the counting

- A small calculation

$$\begin{aligned} Z_3(n) &= \frac{1}{(n!)^2} \sum_{\sigma_i \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1}) \\ &= \frac{1}{(n!)^2} \sum_{\gamma_i \in S_n} \sum_{R_i \vdash n} \chi^{R_1}(\gamma_1) \chi^{R_1}(\gamma_2) \chi^{R_2}(\gamma_1) \chi^{R_2}(\gamma_2) \chi^{R_3}(\gamma_1) \chi^{R_3}(\gamma_2) \\ &= \sum_{R_1, R_2, R_3 \vdash n} (C(R_1, R_2, R_3))^2 \end{aligned} \tag{11}$$

where the symbol

$$C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma) \tag{12}$$

is the Kronecker coefficient.

→ The dimension of an algebra $\mathcal{K}(n)$.

Revisiting the counting

- A small calculation

$$\begin{aligned} Z_3(n) &= \frac{1}{(n!)^2} \sum_{\sigma_i \in S_n} \sum_{\gamma_1, \gamma_2 \in S_n} \delta(\gamma_1 \sigma_1 \gamma_2^{-1} \sigma_1^{-1}) \delta(\gamma_1 \sigma_2 \gamma_2^{-1} \sigma_2^{-1}) \delta(\gamma_1 \sigma_3 \gamma_2^{-1} \sigma_3^{-1}) \\ &= \frac{1}{(n!)^2} \sum_{\gamma_i \in S_n} \sum_{R_i \vdash n} \chi^{R_1}(\gamma_1) \chi^{R_1}(\gamma_2) \chi^{R_2}(\gamma_1) \chi^{R_2}(\gamma_2) \chi^{R_3}(\gamma_1) \chi^{R_3}(\gamma_2) \\ &= \sum_{R_1, R_2, R_3 \vdash n} (C(R_1, R_2, R_3))^2 \end{aligned} \quad (11)$$

where the symbol

$$C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma) \quad (12)$$

is the Kronecker coefficient.

→ The dimension of an algebra $\mathcal{K}(n)$.

Kronecker coefficients $C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)$

Counts

→ multiplicity of the irreps R_3 in the tensor product of the irreps $R_1 \otimes R_2$

→ multiplicity of the one-dimensional (trivial) representation in the tensor product $R_1 \otimes R_2 \otimes R_3$.

• Link with Computational Complexity theory:

→ Finding a combinatorial rule to characterize them in general (Munurghan 1938, Stanley 2000): Find a positive combinatorial interp. i.e. find a family of combinatorial objects O_{R_1, R_2, R_3} such that $C(R_1, R_2, R_3) = \#O_{R_1, R_2, R_3}$

→ Burgisser and Ikenmeyer (2008), computing the Kroneckers is #P-hard and $\in \text{GapP}$

→ On the vanishing of Kronecker coefficients [Ikenmeyer, Mulmuley, Walter, 2015]

1507.02955: deciding positivity of Kronecker coefficients is NP-hard. (contradicts Mulmuley's conjecture that it could be in P; they are not just as the Littlewood-Richardson coefficients)"

My own challenge:

★ At fixed n , can we isolate bi-partites colored graphs corresponding to a non-vanishing $C(R_1, R_2, R_3)$ (or the square of it) ?

The problem is that a graph lives in the σ -space and $C(R_1, R_2, R_3)$ in the R -space.

Kronecker coefficients $C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)$

Counts

→ multiplicity of the irreps R_3 in the tensor product of the irreps $R_1 \otimes R_2$

→ multiplicity of the one-dimensional (trivial) representation in the tensor product $R_1 \otimes R_2 \otimes R_3$.

• Link with Computational Complexity theory:

→ Finding a combinatorial rule to characterize them in general (Munurghan 1938, Stanley 2000): Find a positive combinatorial interp. i.e. find a family of combinatorial objects O_{R_1, R_2, R_3} such that $C(R_1, R_2, R_3) = \#O_{R_1, R_2, R_3}$

→ Burgisser and Ikenmeyer (2008), computing the Kroneckers is $\#P$ -hard and $\in \text{GapP}$

→ On the vanishing of Kronecker coefficients [Ikenmeyer, Mulmuley, Walter, 2015]

1507.02955: deciding positivity of Kronecker coefficients is NP-hard. (contradicts Mulmuley's conjecture that it could be in P; they are not just as the Littlewood-Richardson coefficients)"

My own challenge:

★ At fixed n , can we isolate bi-partites colored graphs corresponding to a non-vanishing $C(R_1, R_2, R_3)$ (or the square of it) ?

The problem is that a graph lives in the σ -space and $C(R_1, R_2, R_3)$ in the R -space.

Kronecker coefficients $C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)$

Counts

→ multiplicity of the irreps R_3 in the tensor product of the irreps $R_1 \otimes R_2$

→ multiplicity of the one-dimensional (trivial) representation in the tensor product $R_1 \otimes R_2 \otimes R_3$.

• Link with Computational Complexity theory:

→ **Finding a combinatorial rule to characterize them in general** (Munurghan 1938, Stanley 2000): Find a positive combinatorial interp. i.e. **find a family of combinatorial objects** O_{R_1, R_2, R_3} such that $C(R_1, R_2, R_3) = \#O_{R_1, R_2, R_3}$

→ Burgisser and Ikenmeyer (2008), computing the Kroneckers is $\#P$ -hard and $\in \text{GapP}$

→ On the vanishing of Kronecker coefficients [Ikenmeyer, Mulmuley, Walter, 2015]

1507.02955: deciding positivity of Kronecker coefficients is NP-hard. (contradicts Mulmuley's conjecture that it could be in P; they are not just as the Littlewood-Richardson coefficients)"

My own challenge:

★ At fixed n , can we isolate bi-partites colored graphs corresponding to a non-vanishing $C(R_1, R_2, R_3)$ (or the square of it) ?

The problem is that a graph lives in the σ -space and $C(R_1, R_2, R_3)$ in the R -space.

Kronecker coefficients $C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)$

Counts

→ multiplicity of the irreps R_3 in the tensor product of the irreps $R_1 \otimes R_2$

→ multiplicity of the one-dimensional (trivial) representation in the tensor product $R_1 \otimes R_2 \otimes R_3$.

• Link with Computational Complexity theory:

→ **Finding a combinatorial rule to characterize them in general** (Munurghan 1938, Stanley 2000): Find a positive combinatorial interp. i.e. **find a family of combinatorial objects** O_{R_1, R_2, R_3} such that $C(R_1, R_2, R_3) = \#O_{R_1, R_2, R_3}$

→ Burgisser and Ikenmeyer (2008), computing the Kroneckers is **#P-hard** and $\in \text{GapP}$

→ On the vanishing of Kronecker coefficients [Ikenmeyer, Mulmuley, Walter, 2015]

1507.02955: **deciding positivity of Kronecker coefficients is NP-hard.** (contradicts Mulmuley's conjecture that it could be in P; they are not just as the Littlewood-Richardson coefficients)"

My own challenge:

★ At fixed n , can we isolate bi-partites colored graphs corresponding to a non-vanishing $C(R_1, R_2, R_3)$ (or the square of it) ?

The problem is that a graph lives in the σ -space and $C(R_1, R_2, R_3)$ in the R -space.

Kronecker coefficients $C(R_1, R_2, R_3) = \frac{1}{n!} \sum_{\sigma \in S_n} \chi^{R_1}(\sigma) \chi^{R_2}(\sigma) \chi^{R_3}(\sigma)$

Counts

→ multiplicity of the irreps R_3 in the tensor product of the irreps $R_1 \otimes R_2$

→ multiplicity of the one-dimensional (trivial) representation in the tensor product $R_1 \otimes R_2 \otimes R_3$.

• Link with Computational Complexity theory:

→ **Finding a combinatorial rule to characterize them in general** (Munurghan 1938, Stanley 2000): Find a positive combinatorial interp. i.e. **find a family of combinatorial objects** O_{R_1, R_2, R_3} such that $C(R_1, R_2, R_3) = \#O_{R_1, R_2, R_3}$

→ Burgisser and Ikenmeyer (2008), computing the Kroneckers is **#P-hard** and $\in \text{GapP}$

→ On the vanishing of Kronecker coefficients [Ikenmeyer, Mulmuley, Walter, 2015]

1507.02955: **deciding positivity of Kronecker coefficients is NP-hard.** (contradicts Mulmuley's conjecture that it could be in P; they are not just as the Littlewood-Richardson coefficients)"

My own challenge:

★ At fixed n , can we isolate bi-partites colored graphs corresponding to a non-vanishing $C(R_1, R_2, R_3)$ (or the square of it) ?

The problem is that a graph lives in the σ -space and $C(R_1, R_2, R_3)$ in the R -space.

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_\sigma \sigma$, $\lambda_\sigma \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i{}^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_{\sigma} \sigma$, $\lambda_{\sigma} \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i{}^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_{\sigma} \sigma$, $\lambda_{\sigma} \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i{}^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_{\sigma} \sigma$, $\lambda_{\sigma} \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i{}^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_{\sigma} \sigma$, $\lambda_{\sigma} \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i{}^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

$\mathcal{K}(n)$, the double coset graph algebra

- Group algebra $\mathbb{C}(S_n)$, i.e. an element of which writes $a = \sum_{\sigma \in S_n} \lambda_{\sigma} \sigma$, $\lambda_{\sigma} \in \mathbb{C}$
- Double coset formulation in $\mathbb{C}(S_n)^{\otimes 3}$,

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma_2, \gamma_1 \sigma_2 \gamma_3, \gamma_1 \sigma_3 \gamma_2) \quad (13)$$

- Define $\mathcal{K}(n) \subset \mathbb{C}(S_n)^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2, \sigma_1, \sigma_2, \sigma_3 \in S_n \right\} \quad (14)$$

→ A fact $\dim_{\mathbb{C}} \mathcal{K}(n) = Z_3(n)$.

→ $\mathcal{K}(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$ which is semi-simple with the pairing

$$\delta_3(\otimes_{i=1}^3 \sigma_i; \otimes_{i=1}^3 \sigma'_i) = \prod_{i=1}^3 \delta(\sigma_i \sigma'_i^{-1}) \quad (15)$$

→ Explain the sum of squares: the Wedderburn-Artin theorem

$$\sum_{R_1, R_2, R_3 \vdash n} (\mathbb{C}(R_1, R_2, R_3))^2 \quad (16)$$

Decomposition of $\mathcal{K}(n)$ in direct subspaces...

Fourier basis of $\mathcal{K}(n)$

- Consider $D_{ij}^R(\sigma) = \langle R, j | \sigma | R, i \rangle$ the matrix representation of σ in the irrep $R \vdash n$ (dimension $d(R)$)

$$\text{Orthogonality : } \sum_{\sigma \in S_n} D_{ij}^R(\sigma) D_{kl}^S(\sigma) = \frac{n!}{d(R)} \delta^{RS} \delta_{ik} \delta_{jl} ;$$

$$\text{Reality : } D_{ij}^R(\sigma^{-1}) = D_{ji}^R(\sigma) ; \quad (17)$$

$$\text{Clebsch - Gordan : } \sum_{\sigma \in S_n} D_{i_1 j_1}^{R_1}(\sigma) D_{i_2 j_2}^{R_2}(\sigma) D_{i_3 j_3}^{R_3}(\sigma) = \frac{n!}{d(R_3)} \sum_{\tau} C_{i_1, i_2; i_3}^{R_1, R_2; R_3, \tau} C_{j_1, j_2; j_3}^{R_1, R_2; R_3, \tau}$$

$$\tau \in \llbracket 1, C(R_1, R_2, R_3) \rrbracket$$

Fourier basis of $\mathcal{K}(n)$ algebra

- Introduce the Fourier basis of $\mathbb{C}(S_n)$

$$Q_{ij}^R = \frac{\kappa_R}{n!} \sum_{\sigma \in S_n} D_{ij}^R(\sigma) \sigma \quad (18)$$

$$\sum_{i_1, j_1, k} C_{i_1, i_2; i_3}^{R_1, R_2; R_3, \tau} C_{j_1, j_2; j_3}^{R_1, R_2; R_3, \tau'} \sum_{\sigma_1, \sigma_2} \rho_L(\sigma_1) \rho_R(\sigma_2) Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} = Q_{\tau, \tau'}^{R_1, R_2, R_3} \quad (19)$$

- The set $\{Q_{\tau, \tau'}^{R_1, R_2, R_3}\}$ forms an orthogonal basis of $\mathcal{K}(n)$.
- At fixed $[R_1, R_2, R_3]$, $Q_{\tau, \tau'}^{R_1, R_2, R_3}$ is matrix with $C(R_1, R_2, R_3)^2$ entries.

→ This is the Wedderburn-Artin basis for $\mathcal{K}(n)$.

Fourier basis of $\mathcal{K}(n)$ algebra

- Introduce the Fourier basis of $\mathbb{C}(S_n)$

$$Q_{ij}^R = \frac{\kappa_R}{n!} \sum_{\sigma \in S_n} D_{ij}^R(\sigma) \sigma \quad (18)$$

$$\sum_{i_1, j_1, k} C_{i_1, i_2; i_3}^{R_1, R_2; R_3, \tau} C_{j_1, j_2; j_3}^{R_1, R_2; R_3, \tau'} \sum_{\sigma_1, \sigma_2} \rho_L(\sigma_1) \rho_R(\sigma_2) Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} = Q_{\tau, \tau'}^{R_1, R_2, R_3} \quad (19)$$

- The set $\{Q_{\tau, \tau'}^{R_1, R_2, R_3}\}$ forms an orthogonal basis of $\mathcal{K}(n)$.
- At fixed $[R_1, R_2, R_3]$, $Q_{\tau, \tau'}^{R_1, R_2, R_3}$ is matrix with $C(R_1, R_2, R_3)^2$ entries.

→ This is the Wedderburn-Artin basis for $\mathcal{K}(n)$.

Fourier basis of $\mathcal{K}(n)$ algebra

- Introduce the Fourier basis of $\mathbb{C}(S_n)$

$$Q_{ij}^R = \frac{\kappa_R}{n!} \sum_{\sigma \in S_n} D_{ij}^R(\sigma) \sigma \quad (18)$$

$$\sum_{i_1, j_1, k} C_{i_1, i_2; i_3}^{R_1, R_2; R_3, \tau} C_{j_1, j_2; j_3}^{R_1, R_2; R_3, \tau'} \sum_{\sigma_1, \sigma_2} \rho_L(\sigma_1) \rho_R(\sigma_2) Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} = Q_{\tau, \tau'}^{R_1, R_2, R_3} \quad (19)$$

- The set $\{Q_{\tau, \tau'}^{R_1, R_2, R_3}\}$ forms an orthogonal basis of $\mathcal{K}(n)$.
- At fixed $[R_1, R_2, R_3]$, $Q_{\tau, \tau'}^{R_1, R_2, R_3}$ is matrix with $C(R_1, R_2, R_3)^2$ entries.

→ This is the Wedderburn-Artin basis for $\mathcal{K}(n)$.

The basis $Q_{\tau, \tau'}^{R_1, R_2, R_3}$

- Invariant $(\gamma_1 \otimes \gamma_2) \cdot Q_{\tau_1, \tau_2}^{R, S, T} \cdot (\gamma_1 \otimes \gamma_2) = Q_{\tau_1, \tau_2}^{R, S, T}$
- Multiply like matrices:

$$Q_{\tau_1, \tau_2}^{R, S, T} Q_{\tau_2', \tau_3}^{R', S', T'} = \delta^{RR'} \delta^{SS'} \delta^{TT'} \delta_{\tau_2 \tau_2'} Q_{\tau_1, \tau_3}^{R, S, T} \quad (20)$$

- Orthogonal under the pairing

$$\delta_2(Q_{\tau_1, \tau_1'}^{R, S, T}; Q_{\tau_2, \tau_2'}^{R', S', T'}) = \kappa_{R, S} d(T) \delta^{RR'} \delta^{SS'} \delta^{TT'} \delta_{\tau_1 \tau_2} \delta_{\tau_1' \tau_2'} \quad (21)$$

→ This guarantees that there is a non degenerate bilinear mapping from a trace form. Hence, the Wedderburn-Artin theorem applies, the algebra is semi-simple.

The center of $\mathcal{K}(n)$

- Trace over Q -basis

$$P^{R,S,T} = \sum_{\tau} Q_{\tau,\tau}^{R,S,T} \quad (22)$$

- $\{P^{R,S,T}\}$ is a basis of the center $\mathcal{Z}(\mathcal{K}(n))$

$$P^{R,S,T} Q_{\tau,\tau'}^{R',S',T'} = Q_{\tau,\tau'}^{R',S',T'} P^{R,S,T} \quad (23)$$

- $\{P^{R,S,T}\}$ is orthogonal

$$\delta_2(P^{R,S,T}; P^{R',S',T'}) = \kappa_{R,S} d(T) C(R, S, T) \delta^{RR'} \delta^{SS'} \delta^{TT'} \quad (24)$$

Proposition

The set $\{P^{R,S,T}\}$ is a basis of $\mathcal{Z}(\mathcal{K}(n))$ and

$$\dim \mathcal{Z}(\mathcal{K}(n)) = \text{number of non vanishing Kronecker coefficients} \quad (25)$$

- For (R, S, T) , $P^{R,S,T}$ is the sum of diagonal elements of $Q_{\tau,\tau'}^{R,S,T}$. Collecting all possible diagonals hence $P^{R,S,T}$ spans the centre $\mathcal{Z}(\mathcal{K}(n))$.

- A triple (R, S, T) , such that $C(R, S, T) \neq 0$ yields a non vanishing $Q_{\tau,\tau'}^{R,S,T}$ and contributes to a single $P^{R,S,T}$. The result on the dimension of $\mathcal{Z}(\mathcal{K}(n))$ follows.

The graph algebra

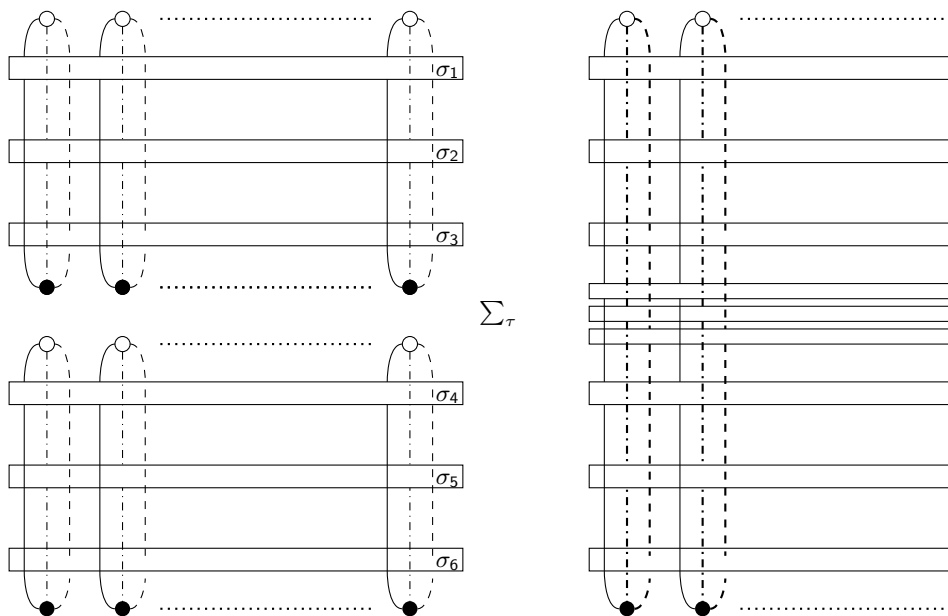
- Convenient normalization

$$A_{\sigma_1, \sigma_2, \sigma_3} = \frac{1}{(n!)^2} \sum_{\gamma_1, \gamma_2 \in S_n} \gamma_1 \sigma_1 \gamma_2 \otimes \gamma_1 \sigma_2 \gamma_2 \otimes \gamma_1 \sigma_3 \gamma_2 \quad (26)$$

- Multiplication

$$A_{\sigma_1, \sigma_2, \sigma_3} A_{\sigma_4, \sigma_5, \sigma_6} = \frac{1}{n!} \sum_{\tau \in S_n} A_{\sigma_1 \tau \sigma_4, \sigma_2 \tau \sigma_5, \sigma_3 \tau \sigma_6} \quad (27)$$

Multiplication of graphs



The algebra $\mathcal{K}(n = 2)$

$$\begin{array}{l}
 \begin{array}{c} \text{3} \\ \text{Diagram 1} \end{array} \cdot \begin{array}{c} \text{A} \cdot \text{E} \\ \text{Diagram 2} \end{array} = \begin{array}{c} \text{A} \\ \text{Diagram 3} \end{array} \\
 \begin{array}{c} \text{3} \\ \text{Diagram 4} \end{array} \cdot \begin{array}{c} \text{A} \cdot \text{A} \\ \text{Diagram 5} \end{array} = \begin{array}{c} \text{E} \\ \text{Diagram 6} \end{array} \\
 \begin{array}{c} \text{3} \\ \text{Diagram 7} \end{array} \cdot \begin{array}{c} \text{A} \cdot \text{B} \\ \text{Diagram 8} \end{array} = \begin{array}{c} \text{D} \\ \text{Diagram 9} \end{array} \\
 \begin{array}{c} \text{3} \\ \text{Diagram 10} \end{array} \cdot \begin{array}{c} \text{A} \cdot \text{D} \\ \text{Diagram 11} \end{array} = \begin{array}{c} \text{B} \\ \text{Diagram 12} \end{array}
 \end{array}$$

(29)

$$\mathcal{K}(2) = \mathbb{C}(S_2) \otimes \mathbb{C}(S_2)$$

(30)

Correlators

At rank $d = 3$, consider the Gaussian model

$$\mathcal{Z} = \int d\Phi d\bar{\Phi} e^{-\frac{1}{2} \sum_{ij} \Phi_{i_1 i_2 i_3} \bar{\Phi}_{i_1 i_2 i_3}} \quad (31)$$

$$\mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} = \sum_{i_l j_l k_l} \Phi_{i_1 j_1 k_1} \Phi_{i_2 j_2 k_2} \cdots \Phi_{i_n j_n k_n} \bar{\Phi}_{i_{\sigma_1(1)} j_{\sigma_2(1)} k_{\sigma_3(1)}} \bar{\Phi}_{i_{\sigma_1(2)} j_{\sigma_2(2)} k_{\sigma_3(2)}} \cdots \bar{\Phi}_{i_{\sigma_1(n)} j_{\sigma_2(n)} k_{\sigma_3(n)}} \quad (32)$$

- The Wick theorem

$$\begin{aligned} \langle \mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} \rangle &= \frac{1}{\mathcal{Z}} \int d\Phi d\bar{\Phi} e^{-\frac{1}{2} \sum_{i,j,k} \Phi_{ijk} \bar{\Phi}_{ijk}} \mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} \\ &= \sum_{i_l j_l k_l} \sum_{\mu \in S_n} \delta_{i_1 i_{\mu(\sigma_1(1))}} \delta_{i_2 i_{\mu(\sigma_1(2))}} \cdots \delta_{i_n i_{\mu(\sigma_1(n))}} \\ &\quad \times \delta_{j_1 j_{\mu(\sigma_2(1))}} \delta_{j_2 j_{\mu(\sigma_2(2))}} \cdots \delta_{j_n j_{\mu(\sigma_2(n))}} \delta_{k_1 k_{\mu(\sigma_3(1))}} \delta_{k_2 k_{\mu(\sigma_3(2))}} \cdots \delta_{k_n k_{\mu(\sigma_3(n))}} \\ &= \sum_{\mu \in S_n} N^{c(\mu\sigma_1) + c(\mu\sigma_2) + c(\mu\sigma_3)} = N^{\#\text{Faces}} \end{aligned} \quad (33)$$

$c(\alpha)$ is the number of cycles of α .

Computable correlators in representation

$$\begin{aligned}\langle \mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} \rangle &= \sum_{\gamma} N^{c(\gamma\sigma_1) + c(\gamma\sigma_2) + c(\gamma\sigma_3)} \\ &= \sum_{\gamma} \sum_{\alpha_1, \alpha_2, \alpha_3} N^{c(\alpha_1) + c(\alpha_2) + c(\alpha_3)} \delta(\gamma\sigma_1\alpha_1) \delta(\gamma\sigma_2\alpha_2) \delta(\gamma\sigma_3\alpha_3) \\ &= \sum_{\gamma} \sum_{R_l} \left[\prod_{l=1}^3 \text{Dim}_N(R_l) \right] \chi^{R_1}(\gamma\sigma_1) \chi^{R_2}(\gamma\sigma_2) \chi^{R_3}(\gamma\sigma_3)\end{aligned}\tag{34}$$

$$\begin{aligned}\mathcal{O}_{S_1, S_2, S_3} &= \frac{1}{(n!)^3} \sum_{\sigma_l \in S_n} \chi^{S_1}(\sigma_1) \chi^{S_2}(\sigma_2) \chi^{S_3}(\sigma_3) \langle \mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} \rangle \\ &= \left[\prod_{l=1}^3 \frac{\text{Dim}_N(S_l)}{d(S_l)} \right] \sum_{\gamma} \chi^{S_1}(\gamma) \chi^{S_2}(\gamma) \chi^{S_3}(\gamma) \\ &= n! \left[\prod_{l=1}^3 \frac{\text{Dim}_N(S_l)}{d(S_l)} \right] C(S_1, S_2, S_3)\end{aligned}\tag{35}$$

$\mathcal{O}_{S_1, S_2, S_3} \propto$ the Kronecker coefficients.

Outline

- 1 Introduction: Permutation groups and Tensor models
- 2 Complex tensors and $U(N)$ invariants
- 3 Extending the results: Real tensors and $O(N)$ invariants

Non-orientable invariants

(Work in progress, with [Nicolas Dub](#) and [Avohou R. Cocou](#))

- Recent interest: SYK-Type models, Carrozza-Tanasa, Klebanov-Tarnolpolsky $O(N)$ models.

- A real tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} O_{p_1 q_1}^{(1)} \cdots O_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad O^{(a)} \in O(N_a) \quad (36)$$

- Tensor contractions = orthogonal invariants

$$S_b^{\text{int}}(T) = \text{Tr}_b(T \cdot T \dots T \cdot T) \quad (37)$$

- Coding orthogonal invariants **b** colored graphs

This one is new!

Non-orientable invariants

(Work in progress, with [Nicolas Dub](#) and [Avohou R. Cocou](#))

- Recent interest: SYK-Type models, Carrozza-Tanasa, Klebanov-Tarnolpolsky $O(N)$ models.

- A real tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} O_{p_1 q_1}^{(1)} \cdots O_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad O^{(a)} \in O(N_a) \quad (36)$$

- Tensor contractions = orthogonal invariants

$$S_b^{\text{int}}(T) = \text{Tr}_b(T \cdot T \dots T \cdot T) \quad (37)$$

- Coding orthogonal invariants **b** colored graphs

This one is new!

Non-orientable invariants

(Work in progress, with [Nicolas Dub](#) and [Avohou R. Cocou](#))

- Recent interest: SYK-Type models, Carrozza-Tanasa, Klebanov-Tarnolpolsky $O(N)$ models.

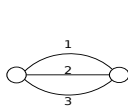
- A real tensor T_{p_1, \dots, p_d} with transformation rule

$$T_{p_1, \dots, p_d} = \sum_{q_k} O_{p_1 q_1}^{(1)} \dots O_{p_d q_d}^{(d)} T_{q_1, \dots, q_d}, \quad O^{(a)} \in O(N_a) \quad (36)$$

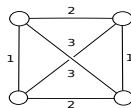
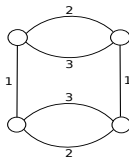
- Tensor contractions = **orthogonal invariants**

$$S_b^{\text{int}}(T) = \text{Tr}_b(T \cdot T \dots T \cdot T) \quad (37)$$

- Coding orthogonal invariants **b colored graphs**



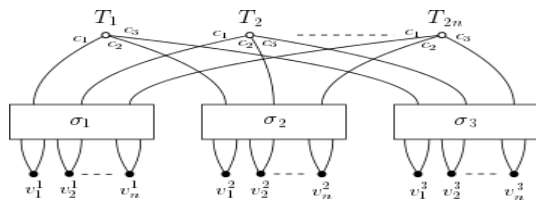
$\text{Tr}_2(T^2)$



This one is new!

Counting orthogonal invariants

Illustration in rank 3:



- Counting permutation triples $(\sigma_1, \sigma_2, \sigma_3) \in (S_{2n} \times S_{2n} \times S_{2n})$ up to the equivalence

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma, \gamma_2 \sigma_2 \gamma, \gamma_3 \sigma_3 \gamma), \quad \gamma_i \in S_n[S_2], \gamma \in S_n. \quad (38)$$
- Element of the double quotient $S_n[S_2] \times S_n[S_2] \times S_n[S_2] \setminus (S_n \times S_n \times S_n) / \text{Diag}(S_n)$.
- Number of invariants

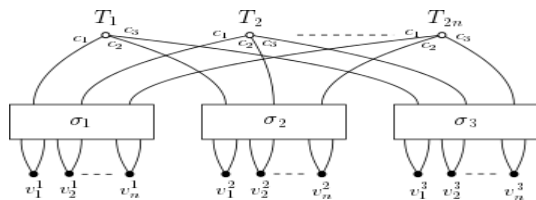
$$Z_{0;3}(2n) = \frac{1}{((n!(2!)^n)^3(2n!))} \sum_{\sigma_i, \gamma \in S_{2n}} \sum_{\gamma_i \in S_n[S_2]} \delta(\gamma_1 \sigma_1 \gamma^{-1} \sigma_1^{-1}) \delta(\gamma_2 \sigma_2 \gamma^{-1} \sigma_2^{-1}) \delta(\gamma_3 \sigma_3 \gamma^{-1} \sigma_3^{-1}) \quad (39)$$

[R.C. Read, "The enumeration of locally restricted graphs," Journal London Math.Soc. 34 (1959), 417-436]

1; 5; 16; 86; 448; 3580; 34981; 448628; 6854130; 121173330

Counting orthogonal invariants

Illustration in rank 3:



- Counting permutation triples $(\sigma_1, \sigma_2, \sigma_3) \in (S_{2n} \times S_{2n} \times S_{2n})$ up to the equivalence

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma, \gamma_2 \sigma_2 \gamma, \gamma_3 \sigma_3 \gamma), \quad \gamma_i \in S_n[S_2], \gamma \in S_n. \quad (38)$$
- Elément of the double quotient $S_n[S_2] \times S_n[S_2] \times S_n[S_2] \setminus (S_n \times S_n \times S_n) / \text{Diag}(S_n)$.
- Number of invariants

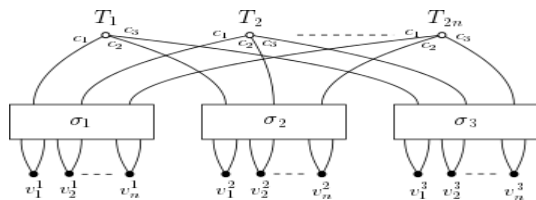
$$Z_{0;3}(2n) = \frac{1}{((n!(2!)^n)^3(2n!))} \sum_{\sigma_i, \gamma \in S_{2n}} \sum_{\gamma_i \in S_n[S_2]} \delta(\gamma_1 \sigma_1 \gamma^{-1} \sigma_1^{-1}) \delta(\gamma_2 \sigma_2 \gamma^{-1} \sigma_2^{-1}) \delta(\gamma_3 \sigma_3 \gamma^{-1} \sigma_3^{-1}) \quad (39)$$

[R.C. Read, "The enumeration of locally restricted graphs," Journal London Math.Soc. 34 (1959), 417-436]

1; 5; 16; 86; 448; 3580; 34981; 448628; 6854130; 121173330

Counting orthogonal invariants

Illustration in rank 3:



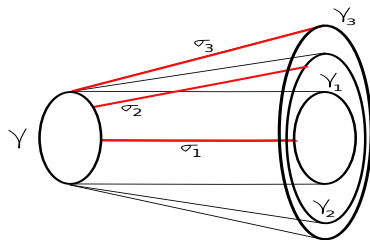
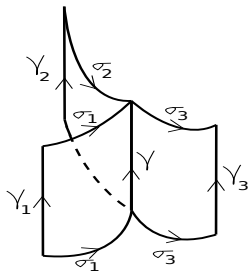
- Counting permutation triples $(\sigma_1, \sigma_2, \sigma_3) \in (S_{2n} \times S_{2n} \times S_{2n})$ up to the equivalence

$$(\sigma_1, \sigma_2, \sigma_3) \sim (\gamma_1 \sigma_1 \gamma, \gamma_2 \sigma_2 \gamma, \gamma_3 \sigma_3 \gamma), \quad \gamma_i \in S_n[S_2], g \in S_n. \quad (38)$$
- Elément of the double quotient $S_n[S_2] \times S_n[S_2] \times S_n[S_2] \setminus (S_n \times S_n \times S_n) / \text{Diag}(S_n)$.
- Number of invariants

$$Z_{0;3}(2n) = \frac{1}{((n!(2!)^n)^3(2n!))} \sum_{\sigma_i, \gamma \in S_{2n}} \sum_{\gamma_i \in S_n[S_2]} \delta(\gamma_1 \sigma_1 \gamma^{-1} \sigma_1^{-1}) \delta(\gamma_2 \sigma_2 \gamma^{-1} \sigma_2^{-1}) \delta(\gamma_3 \sigma_3 \gamma^{-1} \sigma_3^{-1}) \quad (39)$$

[R.C. Read, "The enumeration of locally restricted graphs," Journal London Math.Soc. 34 (1959), 417-436]

1; 5; 16; 86; 448; 3580; 34981; 448628; 6854130; 121173330



The algebra $\mathcal{K}_o(n)$

- Define $\mathcal{K}_o(n) \subset \mathbb{C}(S_{2n})^{\otimes 3}$ is the vector space over \mathbb{C}

$$\mathcal{K}(n) = \text{Span}_{\mathbb{C}} \left\{ \sum_{\gamma_i \in S_n[S_2]; \gamma \in S_{2n}} \gamma_1 \sigma_1 \gamma \otimes \gamma_2 \sigma_2 \gamma \otimes \gamma_3 \sigma_3 \gamma, \sigma_1, \sigma_2, \sigma_3 \in S_{2n} \right\} \quad (40)$$

- Dimension:

$$Z_{o;3}(n) = \frac{1}{(n!(2!)^n)^3} \times \sum_{R_i \vdash 2n} C(R_1, R_2, R_3) \left[\sum_{\gamma_1 \in S_n[S_2]} \chi^{R_1}(\gamma_1) \right] \left[\sum_{\gamma_2 \in S_n[S_2]} \chi^{R_2}(\gamma_2) \right] \left[\sum_{\gamma_3 \in S_n[S_2]} \chi^{R_3}(\gamma_3) \right] \quad (41)$$

→ $\mathcal{K}_o(n)$ is an associative unital subalgebra of $\mathbb{C}(S_n)^{\otimes 3}$.

Question

Given $R \vdash 2n$, how to expand R in irreps representation of $S_n[S_2]$?

The algebra $\mathcal{K}_o(n)$

- Introducing the notation $\vec{X} = (X_1, X_2, X_3)$, $\vec{D}_{ij}^R(\vec{\sigma}) = \prod_{l=1}^3 D_{ijl}^{R_l}(\sigma_l)$,
 $\vec{\sigma}^{\otimes} = \sigma_1 \otimes \sigma_2 \otimes \sigma_3$
- Invariant generators:

$$\begin{aligned}
 Q_{\vec{i}}^{\vec{R};\tau} &= \kappa_{\vec{R};\vec{i}} \sum_{j_l} C_{j_1 j_2 j_3}^{R_1, R_2; R_3, \tau'} \sum_{\gamma_l, \gamma} (P_{\gamma_1}^L \otimes P_{\gamma_2}^L \otimes P_{\gamma_2}^L) P_{\gamma}^R \left[Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} \right] \\
 &= \kappa_{\vec{R};\vec{i}} \sum_{\sigma_l \in S_{2n}} \sum_{p_l, q_l} W_{\vec{i}; \vec{p}}^{\vec{R}} C_{q_1, q_2; q_3}^{R_1, R_2; R_3, \tau'} \vec{D}_{pq}^R(\vec{\sigma}) \vec{\sigma}^{\otimes}
 \end{aligned}$$

$$W_{\vec{i}; \vec{p}}^{\vec{R}} = \left[\sum_{\gamma_1 \in S_n[S_2]} D_{i_1 p_1}^{R_1}(\gamma_1) \right] \left[\sum_{\gamma_2 \in S_n[S_2]} D_{i_2 p_2}^{R_2}(\gamma_2) \right] \left[\sum_{\gamma_3 \in S_n[S_2]} D_{i_3 p_3}^{R_3}(\gamma_3) \right] \quad (42)$$

- The set $\{Q_{\vec{i}}^{\vec{R};\tau}\}$ forms an invariant generating system of $\mathcal{K}_o(n)$:

$$(\gamma_1 \otimes \gamma_2 \otimes \gamma_3) \cdot Q_{\vec{i}}^{\vec{R};\tau} \cdot \gamma^{\otimes 3} = Q_{\vec{i}}^{\vec{R};\tau} \quad (43)$$

The algebra $\mathcal{K}_o(n)$

- Introducing the notation $\vec{X} = (X_1, X_2, X_3)$, $\vec{D}_{ij}^R(\vec{\sigma}) = \prod_{l=1}^3 D_{ijl}^{R_l}(\sigma_l)$,
 $\vec{\sigma}^{\otimes} = \sigma_1 \otimes \sigma_2 \otimes \sigma_3$
- Invariant generators:

$$\begin{aligned}
 Q_{\vec{i}}^{\vec{R};\tau} &= \kappa_{\vec{R};\vec{i}} \sum_{j_l} C_{j_1, j_2, j_3}^{R_1, R_2, R_3, \tau'} \sum_{\gamma_1, \gamma_2} (P_{\gamma_1}^L \otimes P_{\gamma_2}^L \otimes P_{\gamma_2}^L) P_{\gamma}^R \left[Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} \right] \\
 &= \kappa_{\vec{R};\vec{i}} \sum_{\sigma_l \in S_{2n}} \sum_{p_l, q_l} W_{\vec{i};\vec{p}}^{\vec{R}} C_{q_1, q_2, q_3}^{R_1, R_2, R_3, \tau'} \vec{D}_{pq}^R(\vec{\sigma}) \vec{\sigma}^{\otimes}
 \end{aligned}$$

$$W_{\vec{i};\vec{p}}^{\vec{R}} = \left[\sum_{\gamma_1 \in S_n[S_2]} D_{i_1 p_1}^{R_1}(\gamma_1) \right] \left[\sum_{\gamma_2 \in S_n[S_2]} D_{i_2 p_2}^{R_2}(\gamma_2) \right] \left[\sum_{\gamma_3 \in S_n[S_2]} D_{i_3 p_3}^{R_3}(\gamma_3) \right] \quad (42)$$

- The set $\{Q_{\vec{i}}^{\vec{R};\tau}\}$ forms an invariant generating system of $\mathcal{K}_o(n)$:

$$(\gamma_1 \otimes \gamma_2 \otimes \gamma_3) \cdot Q_{\vec{i}}^{\vec{R};\tau} \cdot \gamma^{\otimes 3} = Q_{\vec{i}}^{\vec{R};\tau} \quad (43)$$

The algebra $\mathcal{K}_o(n)$

- Introducing the notation $\vec{X} = (X_1, X_2, X_3)$, $\vec{D}_{ij}^R(\vec{\sigma}) = \prod_{l=1}^3 D_{ijl}^{R_l}(\sigma_l)$,
 $\vec{\sigma}^{\otimes} = \sigma_1 \otimes \sigma_2 \otimes \sigma_3$
- Invariant generators:

$$\begin{aligned}
 Q_{\vec{i}}^{\vec{R};\tau} &= \kappa_{\vec{R};\vec{i}} \sum_{j_l} C_{j_1, j_2, j_3}^{R_1, R_2, R_3, \tau'} \sum_{\gamma_l, \gamma} (P_{\gamma_1}^L \otimes P_{\gamma_2}^L \otimes P_{\gamma_2}^L) P_{\gamma}^R \left[Q_{i_1 j_1}^{R_1} \otimes Q_{i_2 j_2}^{R_2} \otimes Q_{i_3 j_3}^{R_3} \right] \\
 &= \kappa_{\vec{R};\vec{i}} \sum_{\sigma_l \in S_{2n}} \sum_{p_l, q_l} W_{\vec{i}; \vec{p}}^{\vec{R}} C_{q_1, q_2, q_3}^{R_1, R_2, R_3, \tau'} \vec{D}_{pq}^R(\vec{\sigma}) \vec{\sigma}^{\otimes}
 \end{aligned}$$

$$W_{\vec{i}; \vec{p}}^{\vec{R}} = \left[\sum_{\gamma_1 \in S_n[S_2]} D_{i_1 p_1}^{R_1}(\gamma_1) \right] \left[\sum_{\gamma_2 \in S_n[S_2]} D_{i_2 p_2}^{R_2}(\gamma_2) \right] \left[\sum_{\gamma_3 \in S_n[S_2]} D_{i_3 p_3}^{R_3}(\gamma_3) \right] \quad (42)$$

- The set $\{Q_{\vec{i}}^{\vec{R};\tau}\}$ forms an invariant generating system of $\mathcal{K}_o(n)$:

$$(\gamma_1 \otimes \gamma_2 \otimes \gamma_3) \cdot Q_{\vec{i}}^{\vec{R};\tau} \cdot \gamma^{\otimes 3} = Q_{\vec{i}}^{\vec{R};\tau} \quad (43)$$

The algebra $\mathcal{K}_o(n)$

- Checking orthogonality

$$\delta_3(Q_{\vec{i}}^{\vec{R};\tau}; Q_{\vec{i}'}^{\vec{R}';\tau'}) = \frac{(2n)!^3}{d(\vec{R})} \times \kappa_{\vec{R},\vec{i}} \kappa_{\vec{R},\vec{i}'} \delta_{\vec{R},\vec{R}'} \delta_{\tau,\tau'} W_{\vec{i},\vec{i}'}^{\vec{R}} \quad (44)$$

→ Hence depending on $W_{\vec{i},\vec{i}'}^{\vec{R}}$, and the pairing δ_3 might be degenerate on this basis.

→ We still have block decomposition in the \vec{R} and τ sector.

→ Block reduction is **tensor-like** and **not** matrix-like.

→ We might have lost the WA decomposition of the algebra: non-orientable diagrams break the matrix subalgebra decomposition.

The algebra $\mathcal{K}_o(n)$

- Checking orthogonality

$$\delta_3(Q_{\vec{i}}^{\vec{R};\tau}; Q_{\vec{i}'}^{\vec{R}';\tau'}) = \frac{(2n)!^3}{d(\vec{R})} \times \kappa_{\vec{R},\vec{i}} \kappa_{\vec{R},\vec{i}'} \delta_{\vec{R},\vec{R}'} \delta_{\tau,\tau'} W_{\vec{i},\vec{i}'}^{\vec{R}} \quad (44)$$

→ Hence depending on $W_{\vec{i},\vec{i}'}^{\vec{R}}$, and the pairing δ_3 might be degenerate on this basis.

→ We still have block decomposition in the \vec{R} and τ sector.

→ Block reduction is **tensor-like** and **not** matrix-like.

→ We might have lost the WA decomposition of the algebra: non-orientable diagrams break the matrix subalgebra decomposition.

The algebra $\mathcal{K}_o(n)$

- Checking orthogonality

$$\delta_3(Q_{\vec{i}}^{\vec{R};\tau}; Q_{\vec{i}'}^{\vec{R}';\tau'}) = \frac{(2n)!^3}{d(\vec{R})} \times \kappa_{\vec{R},\vec{i}} \kappa_{\vec{R},\vec{i}'} \delta_{\vec{R},\vec{R}'} \delta_{\tau,\tau'} W_{\vec{i},\vec{i}'}^{\vec{R}} \quad (44)$$

→ Hence depending on $W_{\vec{i},\vec{i}'}^{\vec{R}}$, and the pairing δ_3 might be degenerate on this basis.

→ We still have block decomposition in the \vec{R} and τ sector.

→ Block reduction is **tensor-like** and **not** matrix-like.

→ We might have lost the WA decomposition of the algebra: non-orientable diagrams break the matrix subalgebra decomposition.

The algebra $\mathcal{K}_o(n)$

- Checking orthogonality

$$\delta_3(Q_{\vec{i}}^{\vec{R};\tau}; Q_{\vec{i}'}^{\vec{R}';\tau'}) = \frac{(2n)!^3}{d(\vec{R})} \times \kappa_{\vec{R},\vec{i}} \kappa_{\vec{R},\vec{i}'} \delta_{\vec{R},\vec{R}'} \delta_{\tau,\tau'} W_{\vec{i},\vec{i}'}^{\vec{R}} \quad (44)$$

→ Hence depending on $W_{\vec{i},\vec{i}'}^{\vec{R}}$, and the pairing δ_3 might be degenerate on this basis.

→ We still have block decomposition in the \vec{R} and τ sector.

→ Block reduction is **tensor-like** and **not** matrix-like.

→ We might have lost the WA decomposition of the algebra: non-orientable diagrams break the matrix subalgebra decomposition.

The algebra $\mathcal{K}_o(n)$

- Checking orthogonality

$$\delta_3(Q_{\vec{i}}^{\vec{R};\tau}; Q_{\vec{i}'}^{\vec{R}';\tau'}) = \frac{(2n)!^3}{d(\vec{R})} \times \kappa_{\vec{R},\vec{i}} \kappa_{\vec{R},\vec{i}'} \delta_{\vec{R},\vec{R}'} \delta_{\tau,\tau'} W_{\vec{i},\vec{i}'}^{\vec{R}} \quad (44)$$

→ Hence depending on $W_{\vec{i},\vec{i}'}^{\vec{R}}$, and the pairing δ_3 might be degenerate on this basis.

→ We still have block decomposition in the \vec{R} and τ sector.

→ Block reduction is **tensor-like** and **not** matrix-like.

→ We might have lost the WA decomposition of the algebra: non-orientable diagrams break the matrix subalgebra decomposition.

Gaussian correlators

- Consider observables

$$\mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} = \sum_{i_l, j_l, k_l} \Phi_{i_{\sigma_1(1)} j_{\sigma_2(1)} k_{\sigma_3(1)}} \Phi_{i_{\sigma_1(2)} j_{\sigma_2(2)} k_{\sigma_3(2)}} \cdots \Phi_{i_{\sigma_1(2n)} j_{\sigma_2(2n)} k_{\sigma_3(2n)}} \quad (45)$$

$\sigma_i^2 = id$, pairing.

- The Wick theorem for an observable $\mathcal{O}_{\sigma_1, \sigma_2, \sigma_3}$:

$$\langle \mathcal{O}_{\sigma_1, \sigma_2, \sigma_3} \rangle = \sum_{\mu \in \mathcal{S}_n} N^{c(\mu\sigma_1) + c(\mu\sigma_2) + c(\mu\sigma_3)} = N^{\#Faces} \quad (46)$$

μ is again a pairing, $\mu^2 = id$.

Conclusion

- The counting of unitary invariants gives a sum of squares of Kroneckers.
- Unitary invariants span a semi-simple graph algebra, that is an algebra which decomposes in matrix subalgebras (WA theorem).
- WA orthogonal bases of the graph algebra have been found using representation theory, and the dimension of the algebra well reflects in the base label.
- Correlators can be computed in the same formulation. Fourier transform of the correlators gives function which are exactly computable (via softwares).
- Much more results in [JBG & Ramgoolam, JHEP 17]: for instance, the algebra become graded by irreps representations of S_d , where d is the tensor rank.
- Orthogonal invariants: More numerous and can be enumerated using similar symmetric group techniques.
- In the representation, the counting expresses as Kroneckers convoluted with characters (a bit more/less involved).
- The underlying TFT of the counting: a TFT on toroidal lattice with defects (effects of $S_n[S_2]$ subgroup).
- They span a graph algebra that might be not semi-simple. We have the simplest a set of invariant generators that we conjecture to be overcomplete.

Conclusion

- The counting of unitary invariants gives a sum of squares of Kroneckers.
 - Unitary invariants span a semi-simple graph algebra, that is an algebra which decomposes in matrix subalgebras (WA theorem).
 - WA orthogonal bases of the graph algebra have been found using representation theory, and the dimension of the algebra well reflects in the base label.
 - Correlators can be computed in the same formulation. Fourier transform of the correlators gives function which are exactly computable (via softwares).
 - Much more results in [JBG & Ramgoolam, JHEP 17]: for instance, the algebra become graded by irreps representations of S_d , where d is the tensor rank.
-
- Orthogonal invariants: More numerous and can be enumerated using similar symmetric group techniques.
 - In the representation, the counting expresses as Kroneckers convoluted with characters (a bit more/less involved).
 - The underlying TFT of the counting: a TFT on toroidal lattice with defects (effects of $S_n[S_2]$ subgroup).
 - They span a graph algebra that might be not semi-simple. We have the simplest a set of invariant generators that we conjecture to be overcomplete.

Conclusion

- The counting of unitary invariants gives a sum of squares of Kroneckers.
- Unitary invariants span a semi-simple graph algebra, that is an algebra which decomposes in matrix subalgebras (WA theorem).
- WA orthogonal bases of the graph algebra have been found using representation theory, and the dimension of the algebra well reflects in the base label.
- Correlators can be computed in the same formulation. Fourier transform of the correlators gives function which are exactly computable (via softwares).
- Much more results in [JBG & Ramgoolam, JHEP 17]: for instance, the algebra become graded by irreps representations of S_d , where d is the tensor rank.

- Orthogonal invariants: More numerous and can be enumerated using similar symmetric group techniques.
- In the representation, the counting expresses as Kroneckers convoluted with characters (a bit more/less involved).
- The underlying TFT of the counting: a TFT on toroidal lattice with defects (effects of $S_n[S_2]$ subgroup).
- They span a graph algebra that might be not semi-simple. We have the simplest a set of invariant generators that we conjecture to be overcomplete.

Conclusion

- The counting of unitary invariants gives a sum of squares of Kroneckers.
- Unitary invariants span a semi-simple graph algebra, that is an algebra which decomposes in matrix subalgebras (WA theorem).
- WA orthogonal bases of the graph algebra have been found using representation theory, and the dimension of the algebra well reflects in the base label.
- Correlators can be computed in the same formulation. Fourier transform of the correlators gives function which are exactly computable (via softwares).
- Much more results in [JBG & Ramgoolam, JHEP 17]: for instance, the algebra become graded by irreps representations of S_d , where d is the tensor rank.

- Orthogonal invariants: More numerous and can be enumerated using similar symmetric group techniques.
- In the representation, the counting expresses as Kroneckers convoluted with characters (a bit more/less involved).
- The underlying TFT of the counting: a TFT on toroidal lattice with defects (effects of $S_n[S_2]$ subgroup).
- They span a graph algebra that might be not semi-simple. We have the simplest a set of invariant generators that we conjecture to be overcomplete.