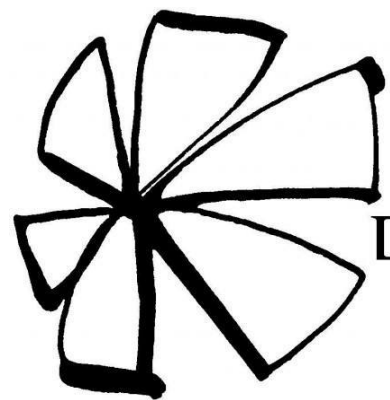


Tensor models from a functional Renormalization Group perspective

Astrid Eichhorn

based on work with Tim Koslowski, Johannes Lumma and Antonio Pereira

OIST mini-symposium on Holographic Tensors



Die Junge Akademie

October 31, 2018

University of Heidelberg



RUPRECHT-KARLS-
UNIVERSITÄT HEIDELBERG

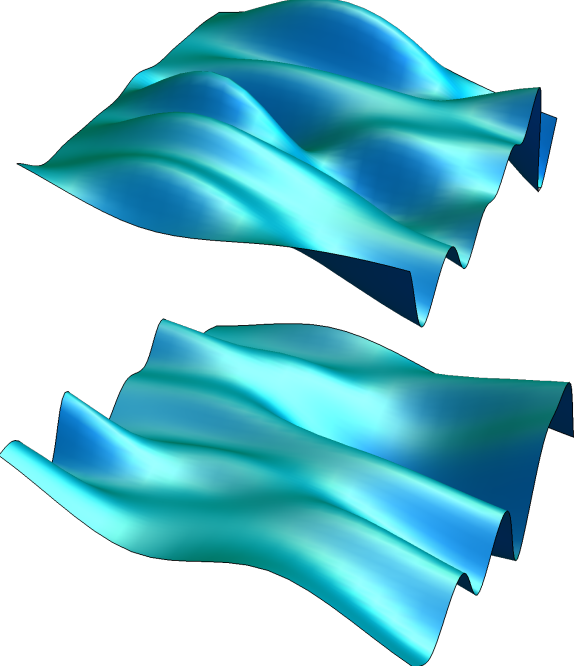
ZUKUNFT SEIT 1386

**Emmy
Noether-
Programm**

DFG Deutsche
Forschungsgemeinschaft



Motivation: the path integral for quantum gravity

$$\int_{\text{spacetimes}} e^{iS[\text{spacetime}]}$$
Two wavy, blue, translucent surfaces representing spacetime fluctuations. The surfaces are rendered with a 3D effect, showing peaks and valleys, and are colored with a gradient from light blue to dark blue. They are positioned to the right of the mathematical expression.

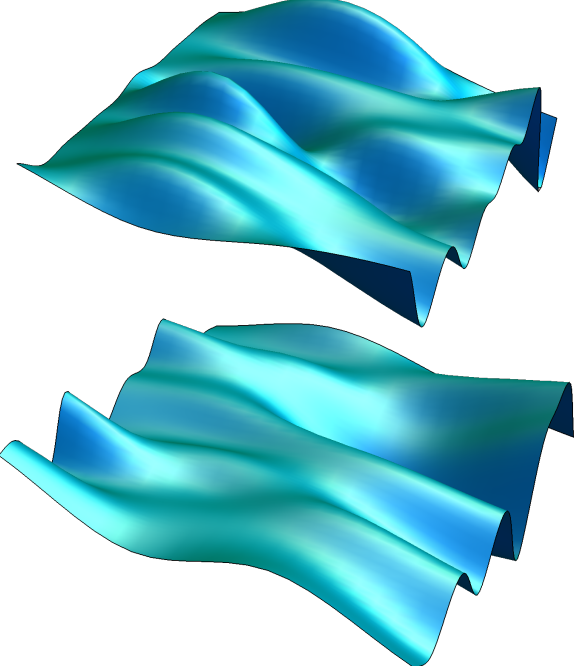
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→ **yes, for**

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→ **possible to go beyond?**

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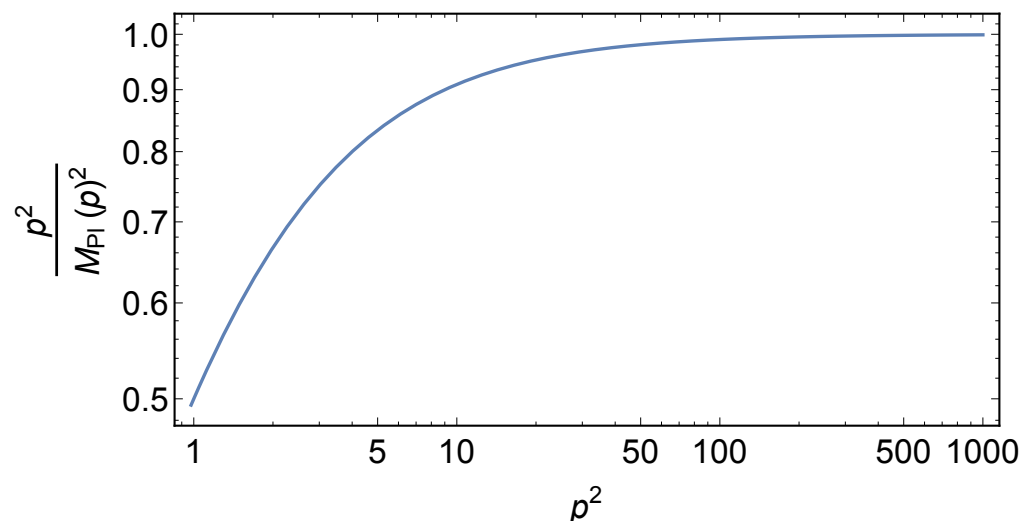
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continuum models can be discrete & discrete models continuous

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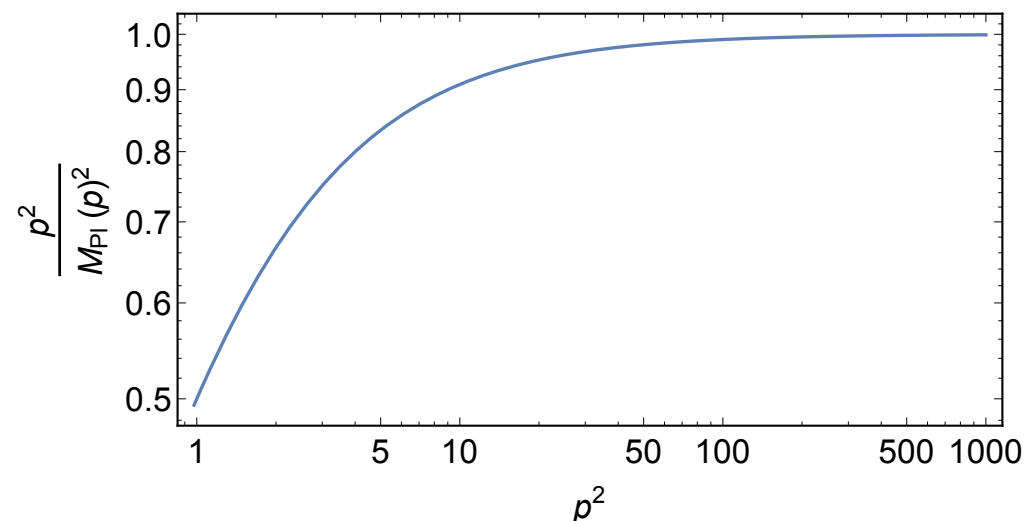
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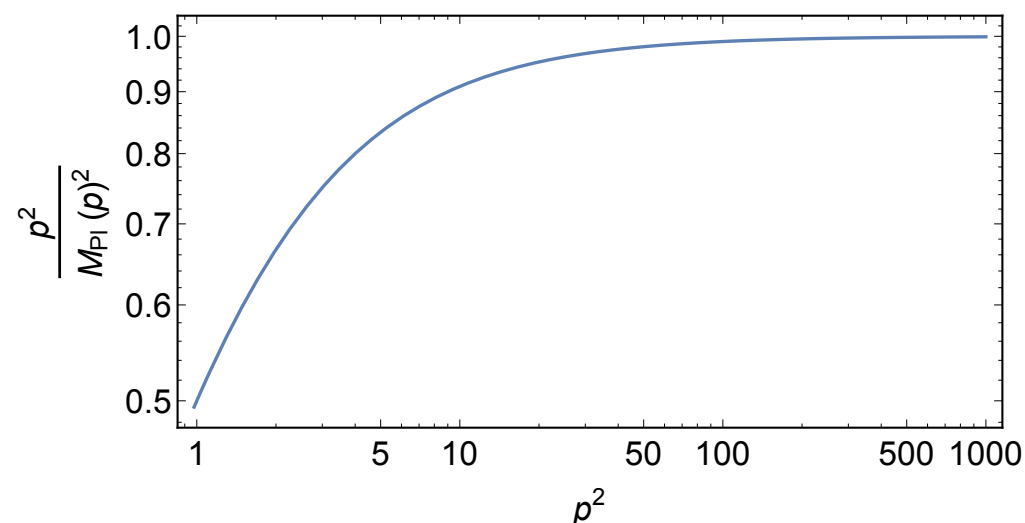
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continuum models can be discrete

&

discrete models continuous

- discrete spectra of geometric operators in LQG
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- perfect action
(lattice action provides continuum dynamics even at finite lattice spacing)
- kinematical discreteness (at the level of individual configurations) need not persist at the dynamical level
(path integral: sum over discrete config.'s yields continuum effective geometry)

Continuum limit: predictivity at all scales

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Models defined with a simple cutoff lose predictivity at the cutoff scale:

example: Effective field theory for gravity

$$S[g_{\mu\nu}] = \int d^4x \sqrt{g} \left(\frac{\Lambda}{16\pi G_N} - \frac{R}{16\pi G_N} + a_1 R^2 + b_1 R_{\mu\nu} R^{\mu\nu} + \dots \right)$$

Breaks down at cutoff scale, because the infinitely many choices for parameters matter for predictions at $E = M_{Pl}$

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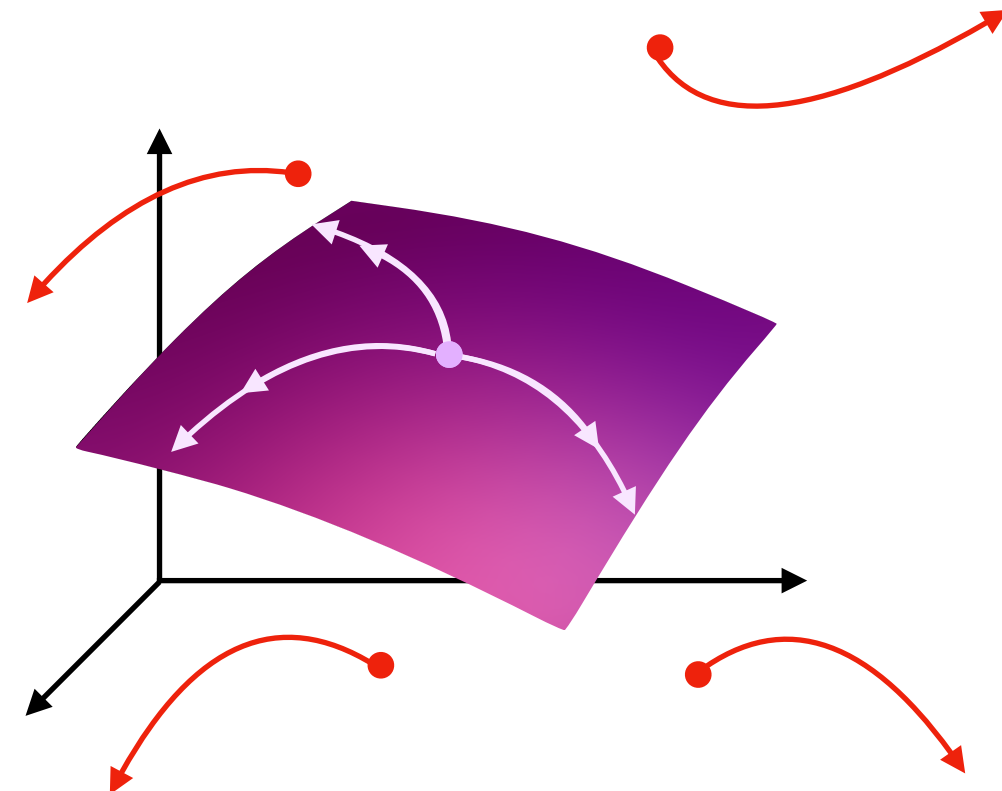
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- Continuum limit = Renormalization Group fixed point: finitely many relevant directions (general argument based on dimensional analysis)



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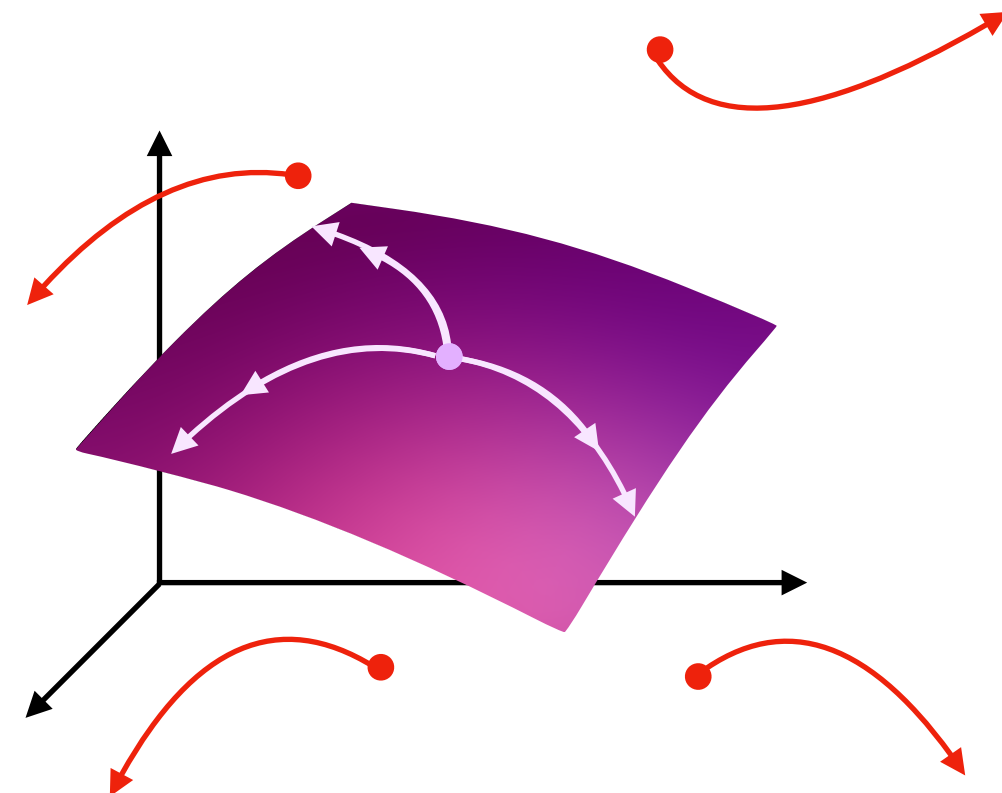
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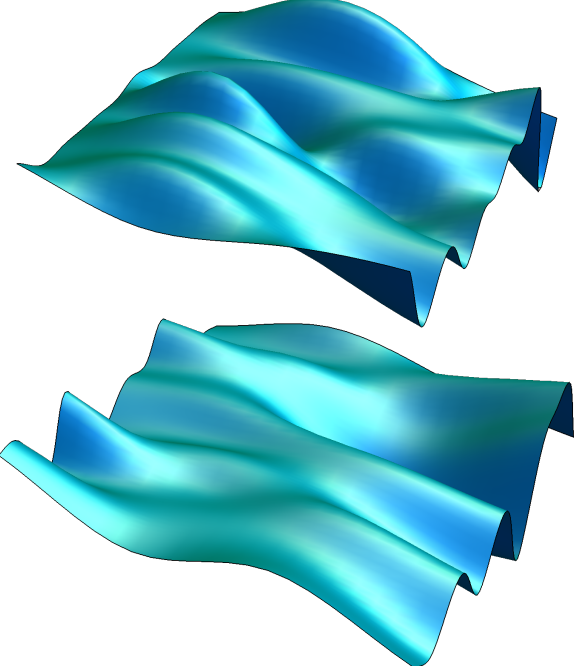
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- Continuum limit = Renormalization Group fixed point: finitely many relevant directions (general argument based on dimensional analysis)
- Continuum limit = mathematical phase transition: Physics at criticality determined by finitely many parameters



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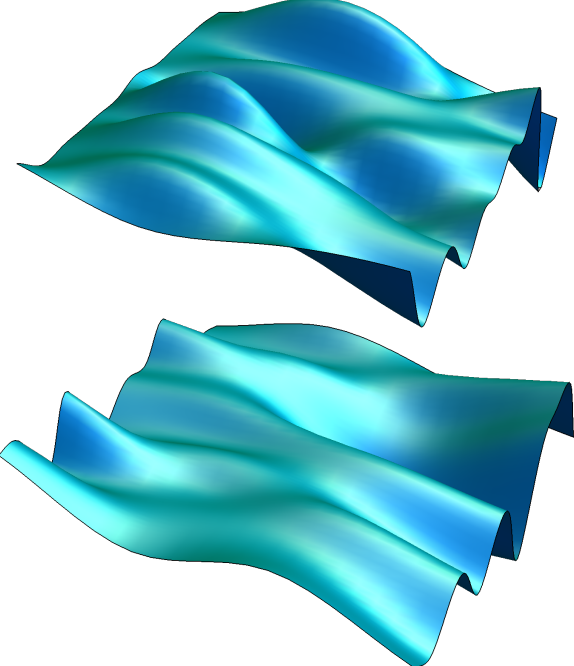
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asymptotically safe quantum gravity

tensor models

**(causal) dynamical triangulations
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indirect **asymptotically safe quantum gravity**

trade-off?

- continuum limit?

direct

tensor models

- background independence?

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The continuum limit from local coarse graining

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$$\int \mathcal{D}\gamma_{\mu\nu} e^{iS[\gamma_{\mu\nu}]} \rightarrow \int \mathcal{D}\gamma_{\mu\nu} e^{-S[\gamma_{\mu\nu}]} \rightarrow \int_{\text{"}p^2 < k^2\text{"}} \mathcal{D}\gamma_{\mu\nu} e^{-S[\gamma_{\mu\nu}]}$$

local coarse graining:
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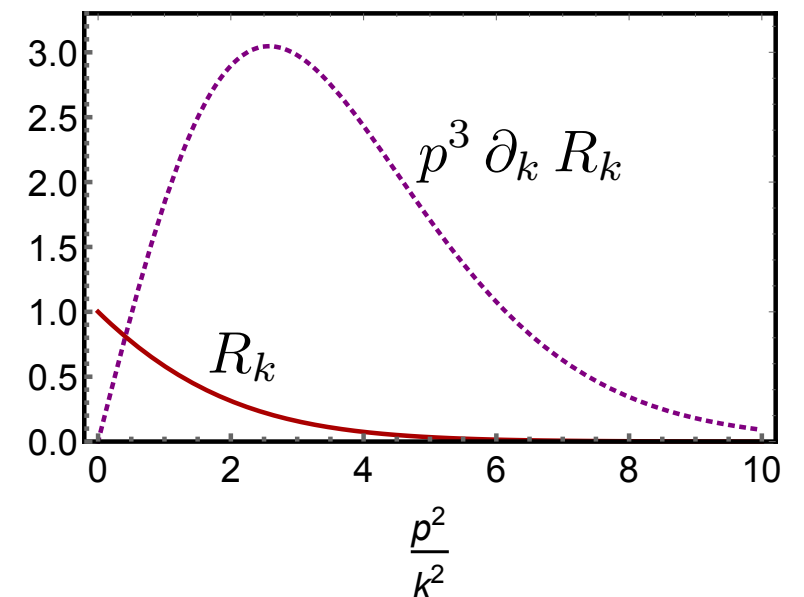
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quadratic suppression term



The continuum limit from local coarse graining

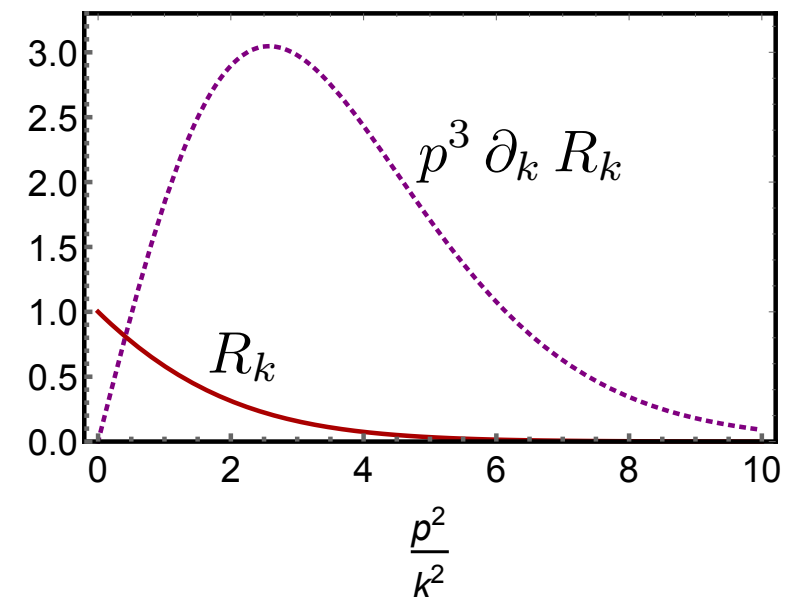
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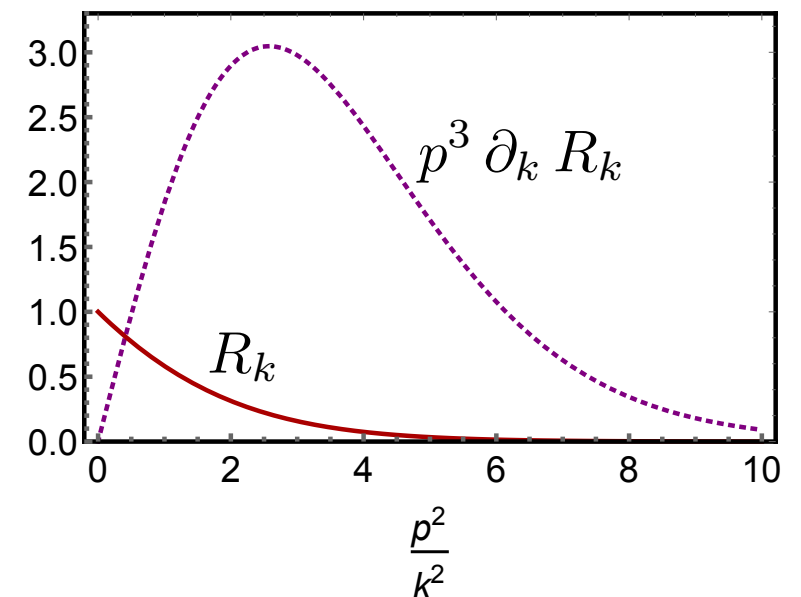
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**change of effective dynamics at k
driven by fluctuations at momenta k**

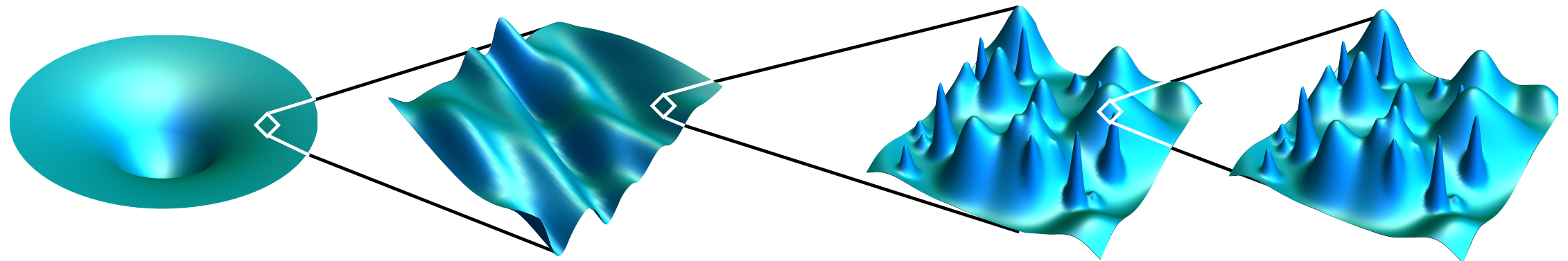
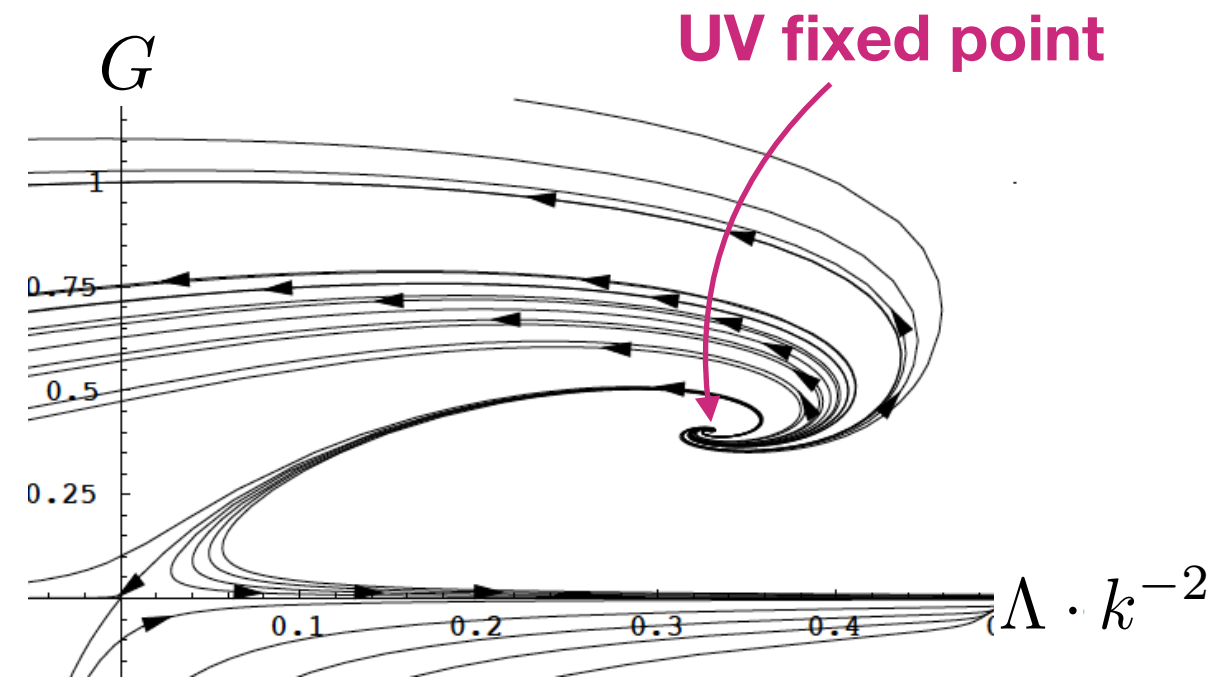
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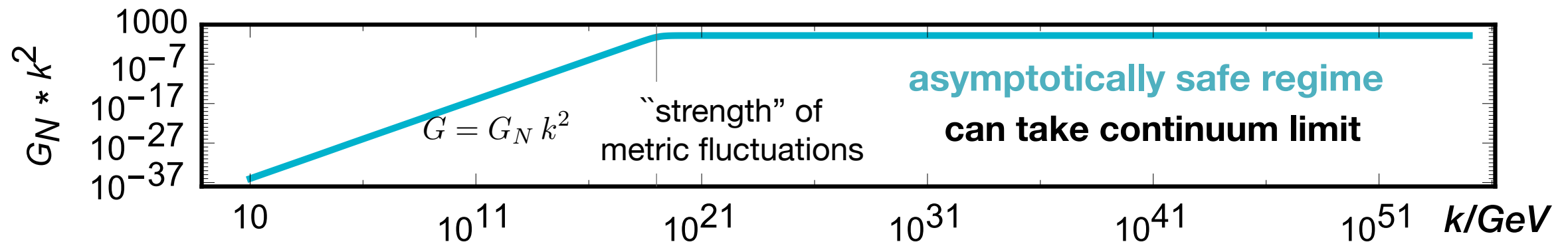
**Compelling evidence for Reuter fixed point
(with and without matter)
such that continuum limit can be taken**



classical gravity regime

Planck scale

quantum scale invariance

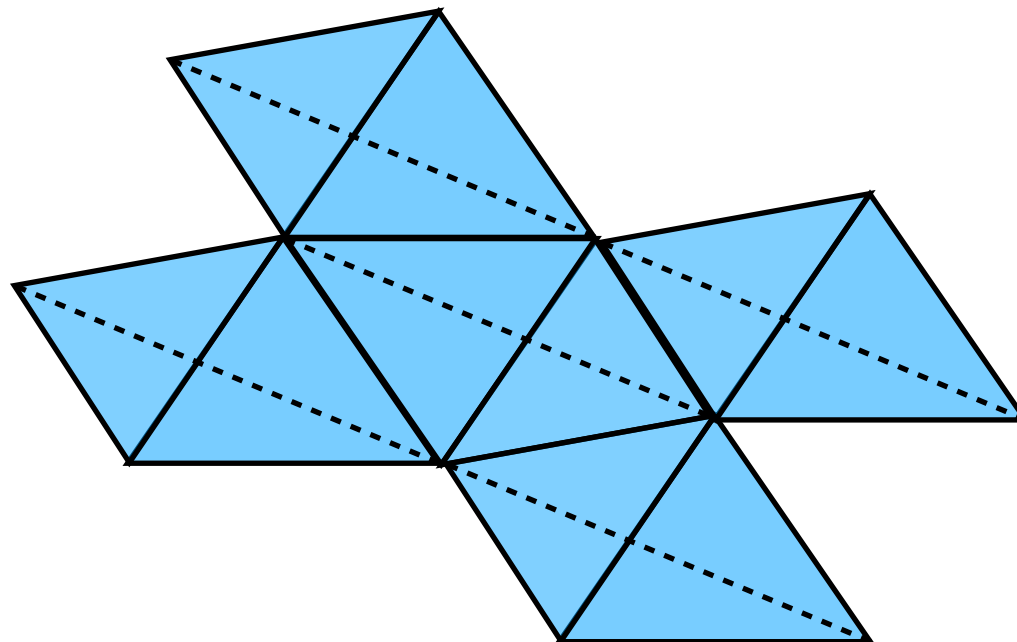
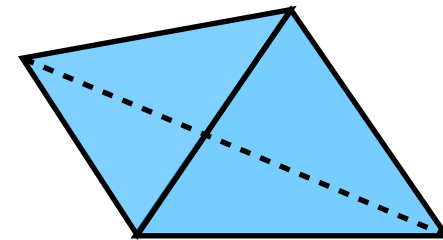
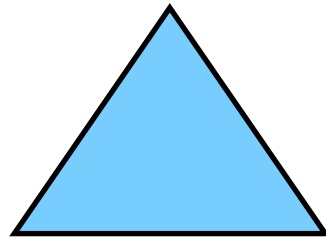


The continuum limit from non-local coarse graining

$$\int_{\text{spacetimes}} e^{-S} \xrightarrow{\sim} \sum_{\text{triangulations}} e^{-S_{\text{discrete}}}$$

tensor models: rank D , N^D tensor, $O(N)^D$ invariance [Sasakura '91; Gross '91; Ambjorn, Durhuus, Jonsson '91; Gurau '11; Gurau, Rivasseau '11]

$$\xrightarrow{\sim} \int \mathcal{D}T e^{-T_{ijk}T_{ijk} + g_{4,1}^2 T_{ijk}T_{ilm}T_{nlm}T_{njk} + g_{4,1} \underbrace{T_{ijk}T_{ilm}T_{njm}T_{nlk}}_{\text{quadrilateral}}}$$



**Feynman diagrams:
triangulations**

The continuum limit from non-local coarse graining

example: rank 3; generalization to rank D straightforward

$$Z_N = \int \mathcal{D}T e^{-S[T] + \text{Tr} T \cdot J - \frac{1}{2} T_{abc} R_N(a, b, c, d, e, f)^{abcdef} T_{def}}$$

suppresses “infrared” indices

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see also R. Toriumi, T. Krajewski for Polchinski equation

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tensor components at N

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tensor components at N

formal “trick”: can transform any integral into differential equation
what is the physical meaning?

The continuum limit from non-local coarse graining

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tensor components at N

How is this coarse graining?

Abstract view on RG flows:

RG flows go from many to few degrees of freedom (a-theorem)

(background dependent: local coarse graining; background independent: nonlocal “coarse graining”)

The continuum limit from non-local coarse graining

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search for universality at large N : fixed point of abstract RG flow

blueprint: matrix models for quantum gravity

double scaling limit $(g - g_{\text{crit}})^{\frac{1}{\theta}} N = \text{const}$

= continuum limit with contributions from all topologies

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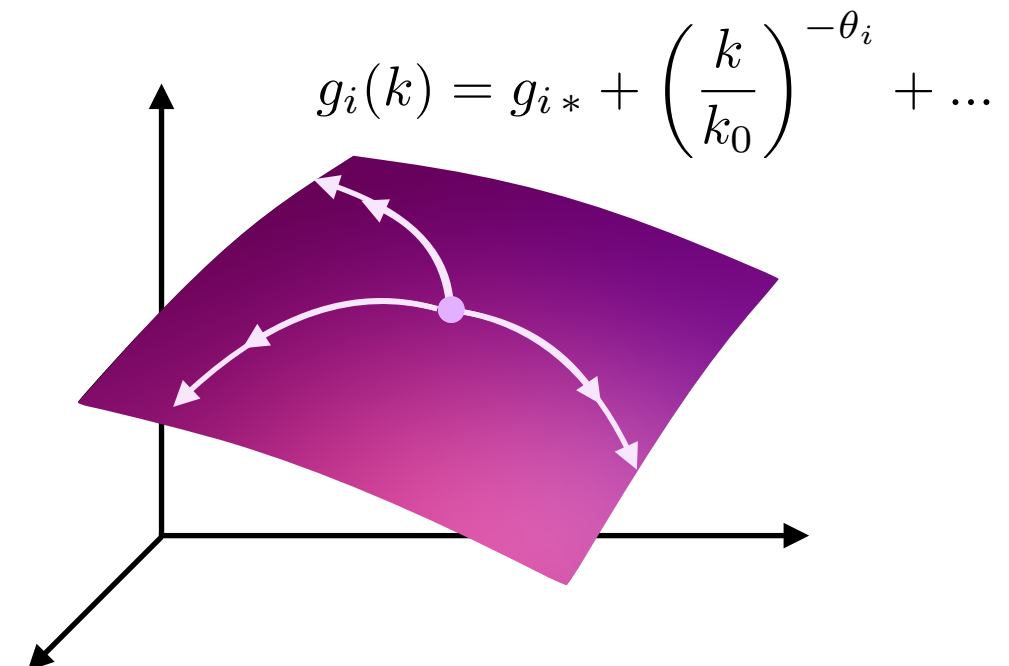
= continuum limit with contributions from all topologies

$$\longrightarrow g(N) = g_{\text{crit}} + N^{-\theta}$$

linearized “RG flow” in tensor size N

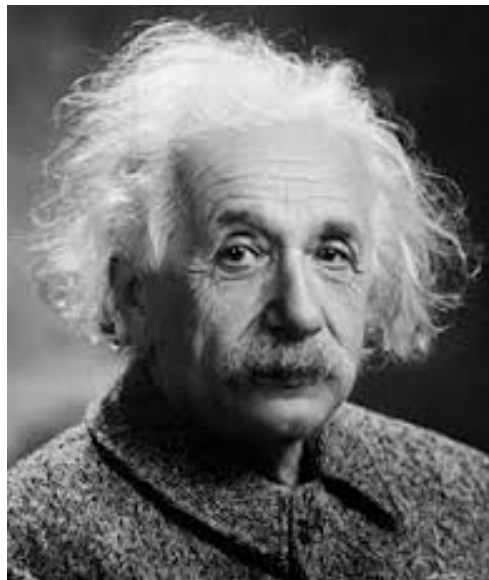
[Brezin, Zinn-Justin '92]

[AE, Koslowski '13]

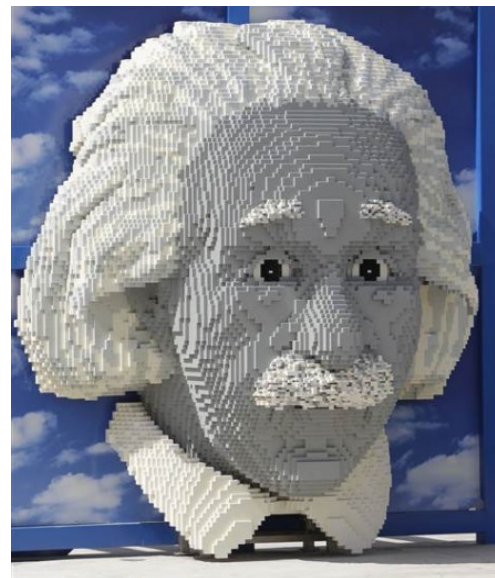


RG fixed point in tensor models

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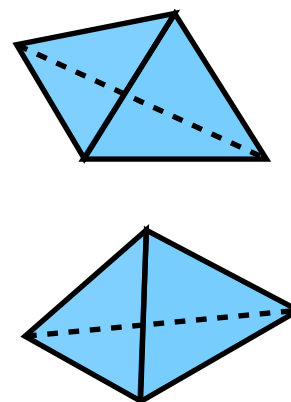
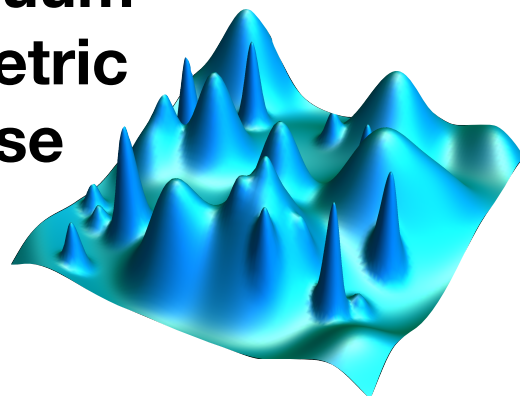


continuum



discrete

continuum
geometric
phase



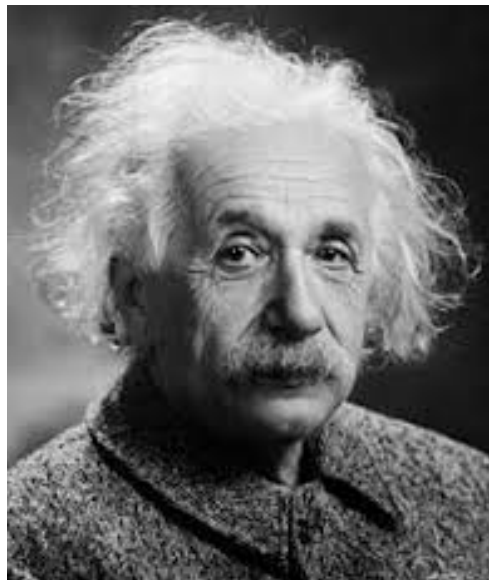
discrete
“pregeometric”
phase

transition

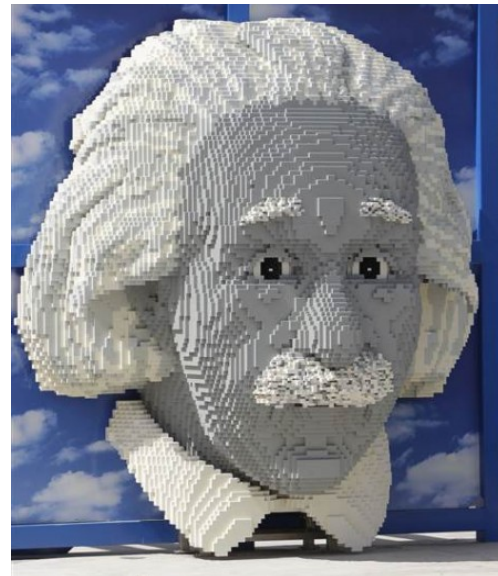
interaction strength

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continuum



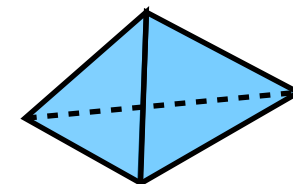
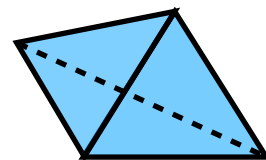
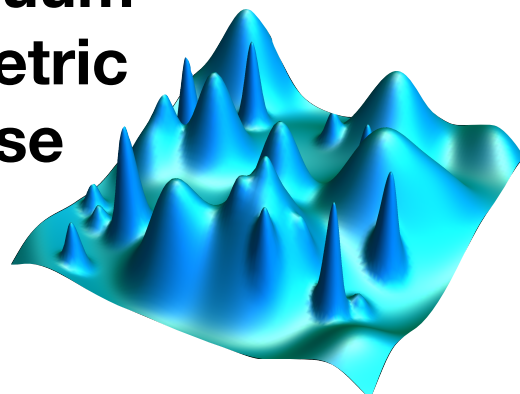
discrete

**Asymptotic safety & dynamical triangulations:
Two sides of the same (mathematical)
phase transition?**

(same configuration spaces?)

Testing:
Comparison of universal
scaling exponents

continuum
geometric
phase



discrete
“pregeometric”
phase

transition

interaction strength

FRG for tensor models in practice

$$\partial_N \Gamma_N = \frac{1}{2} \text{Tr} \left[\left(\frac{\partial^2 \Gamma_N}{\partial T^2} + R_N \right)^{-1} \partial_N R_N \right]$$

- **background independence implies no mass-dimensionality**
 - **rhs of flow equation determines “canonical” scaling dimension**
- **rhs of flow equation generates all couplings compatible with symmetries**
 - **truncations of theory space required in practice**
- **regulator breaks $O(N)$ or $U(N)$ symmetry (no free lunch for coarse graining in QG?)**
 - **projection onto couplings not based on combinatorial structure of interactions alone**
 - **Ward-identities imposed in enlarged theory space**

Canonical scaling dimension in background-independent RG flow

local flows

background-independent flows

Canonical scaling dimension in background-independent RG flow

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background-independent flows

RG step:
local averaging
scale transformation

RG step:
non-local averaging

canonical scaling dimensions
(at origin in coupling space)
defined by *mass* dimensionality

no units, no scales,
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**agrees with dimensions from
autonomous system of β functions**

$$\beta_{\bar{G}} = -\#k^{(d-2)} \bar{G}^2 + \dots$$

$$G = \bar{G}k^{d-2}$$

$$\beta_G = 2G - \#G^2 + \dots$$

Canonical scaling dimension in background-independent RG flow

local flows

RG step:
local averaging
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$$\beta_{\bar{G}} = -\#k^{(d-2)} \bar{G}^2 + \dots$$

$$G = \bar{G}k^{d-2}$$

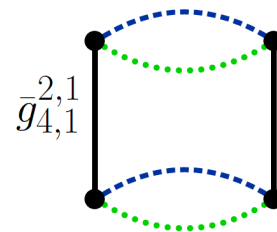
$$\beta_G = 2G - \#G^2 + \dots$$

background-independent flows

RG step:
non-local averaging

no units, no scales,
no a priori scaling dimensions

**canonical dimensions from
autonomous system of β functions
in large N limit**



$$\beta_{\bar{g}_{4,1}^{2,1}} = 2 N^2 \left(\bar{g}_{4,1}^{2,1} \right)^2 + \dots$$

$$g_{4,1}^{2,1} = \bar{g}_{4,1}^{2,1} N^2$$

$$\beta_{g_{4,1}^{2,1}} = 2g_{4,1}^{2,1} + 2 \left(g_{4,1}^{2,1} \right)^2 + \dots$$

A bootstrap strategy to choose truncations

$$\partial_N \Gamma_N = \frac{1}{2} \text{Tr} \left[\left(\frac{\partial^2 \Gamma_N}{\partial T^2} + R_N \right)^{-1} \partial_N R_N \right]$$

- **aim at quantitative determination of relevant critical exponents**

$\Rightarrow$ all relevant couplings at a fixed point should be included

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- **aim at quantitative determination of relevant critical exponents**

$\Rightarrow$ all relevant couplings at a fixed point should be included

- 1. assume ordering principle in theory space (e.g., near-canonical)**
- 2. include all couplings expected to be relevant & leading irrelevant ones**
- 3. check ordering principle at fixed points in truncation**

A lightning review of FRG results in 2 and 3 dimensions

Benchmarks:

2d gravity

2d gravity with matter: multicritical points

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Benchmarks:

2d gravity $\theta_1 = 0.8$ [Douglas, Shenker '90; Brezin, Kazakov '90; Gross, Migdal '90]

$$g_i \text{Tr}(\phi^i)$$

TABLE I: Critical exponents θ_i and fixed point values at $\mathcal{O}(n)$ in the $\mathcal{P}^{-1}\mathcal{F}$ expansion.

n_{max}	θ_1	θ_2	θ_3	θ_4	θ_5	θ_6	g_4^*	$g_6^* \cdot 10^3$	$g_8^* \cdot 10^4$	$g_{10}^* \cdot 10^5$	$g_{12}^* \cdot 10^5$	$g_{14}^* \cdot 10^6$
2	1.066						-0.1005					
3	1.046	-1.080					-0.0722	-3.8				
4	1.036	-1.053	-2.14				-0.0563	-4.6	-3.6			
5	1.029	-1.037	-2.115	-3.171			-0.0461	-4.7	-5.4	-5		
6	1.025	-1.027	-2.10	-3.137	-4.197		-0.0390	-4.5	-6.2	-8	-1	
7	1.022	-1.020	-2.093	-3.110	-4.172	-5.213	-0.0338	-4.2	-6.5	-10	-1	-1.5

[AE, Koslowski, '13]

2d gravity with matter: multicritical points [Kazakov, '89] $\theta_{m-1} = \frac{2}{m + \frac{1}{2}}$

$$g_{i,j,k} \text{Tr}(\phi^i) \text{Tr}(\phi^j) \text{Tr}(\phi^k)$$

g_4	g_6	g_8	g_{10}	g_{12}	g_{22}	g_{24}	g_{26}	g_{28}	g_{44}	g_{46}	g_{222}	g_{224}	θ_1	θ_2	θ_3	θ_4	θ_5
$-\frac{1}{4x}$	0	0	0	0	0	0	0	0	0	0	0	0	1	-1/2	-1	-1	-3/2
$-\frac{8}{21x}$	$\frac{2}{63x^2}$	0	0	0	$\frac{1}{21x}$	0	0	0	0	0	0	0	0.82	0.82	-0.33	-0.64	-0.67
$-\frac{51}{116x}$	$\frac{27}{464x^2}$	$-\frac{9}{3712x^3}$	0	0	$\frac{15}{232x}$	$-\frac{3}{464x^2}$	0	0	0	0	0	0	0.80	0.80	0.61	-0.25	-0.5
$\frac{-16448}{34535x}$	$\frac{13392}{172675x^2}$	$\frac{-4608}{863375x^3}$	$\frac{576}{4316875x^4}$	0	$\frac{2634}{34535x}$	$\frac{-2112}{172675x^2}$	$\frac{288}{863375x^3}$	0	$\frac{108}{172675x^2}$	0	$\frac{144}{863375x^3}$	0	0.81	0.81	0.68	0.45	-0.2

[AE, Koslowski, '14]

A lightning review of FRG results in 2 and 3 dimensions

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5	1.029	-1.037	-2.115	-3.171			-0.0461	-4.7	-5.4	-5		
6	1.025	-1.027	-2.10	-3.137	-4.197		-0.0390	-4.5	-6.2	-8	-1	
7	1.022	-1.020	-2.093	-3.110	-4.172	-5.213	-0.0338	-4.2	-6.5	-10	-1	-1.5

$$g_i \text{Tr}(\phi^i)$$

**reasonable results:
simple truncations sufficient**

[AE, Koslowski, '13]

2d gravity with matter: multicritical points [Kazakov, '89] $\theta_{m-1} = \frac{2}{m + \frac{1}{2}}$

g_4	g_6	g_8	g_{10}	g_{12}	g_{22}	g_{24}	g_{26}	g_{28}	g_{44}	g_{46}	g_{222}	g_{224}	θ_1	θ_2	θ_3	θ_4	θ_5
$-\frac{1}{4x}$	0	0	0	0	0	0	0	0	0	0	0	0	1	-1/2	-1	-1	-3/2
$-\frac{8}{21x}$	$\frac{2}{63x^2}$	0	0	0	$\frac{1}{21x}$	0	0	0	0	0	0	0	0.82	0.82	-0.33	-0.64	-0.67
$-\frac{51}{116x}$	$\frac{27}{464x^2}$	$-\frac{9}{3712x^3}$	0	0	$\frac{15}{232x}$	$-\frac{3}{464x^2}$	0	0	0	0	0	0	0.80	0.80	0.61	-0.25	-0.5
$\frac{-16448}{34535x}$	$\frac{13392}{172675x^2}$	$\frac{-4608}{863375x^3}$	$\frac{576}{4316875x^4}$	0	$\frac{2634}{34535x}$	$\frac{-2112}{172675x^2}$	$\frac{288}{863375x^3}$	0	$\frac{108}{172675x^2}$	0	$\frac{144}{863375x^3}$	0	0.81	0.81	0.68	0.45	-0.2

$$g_{i,j,k} \text{Tr}(\phi^i) \text{Tr}(\phi^j) \text{Tr}(\phi^k)$$

[AE, Koslowski, '14]

A lightning review of FRG results in 2 and 3 dimensions

rank 3 complex:

[AE, Koslowski, '17]

$$\Gamma_N = Z_N \left(\text{diagram} \right) + g_{4,1}^{2,1} \left(\text{diagram} \right) + g_{4,1}^{2,2} \left(\text{diagram} \right) + g_{4,1}^{2,3} \left(\text{diagram} \right) + g_{4,2}^2 \left(\text{diagram} \right)$$

$$+ g_{6,1}^{3,1} \left(\text{diagram} \right) + g_{6,1}^{3,2} \left(\text{diagram} \right) + g_{6,1}^{3,3} \left(\text{diagram} \right) + g_{6,2}^{3,1} \left(\text{diagram} \right) + g_{6,2}^{3,2} \left(\text{diagram} \right) + g_{6,2}^{3,3} \left(\text{diagram} \right) + g_{6,3}^3 \left(\text{diagram} \right) + g_{6,1}^{2,1} \left(\text{diagram} \right) + g_{6,1}^{2,2} \left(\text{diagram} \right) + g_{6,1}^{2,3} \left(\text{diagram} \right) + g_{6,1}^0 \left(\text{diagram} \right)$$

dynamical decoupling unless additional symmetry

$$g_{4,1}^{2,1} = g_{4,1}^{2,2} = g_{4,1}^{2,3} \text{ is imposed}$$

$$\beta_{g_{4,2}^2} = (3 + 2\eta)g_{4,2}^2 + \frac{6 - \eta}{15} \left((g_{4,2}^2)^2 + 2g_{4,2}^2 (g_{4,1}^{2,1} + g_{4,1}^{2,2} + g_{4,1}^{2,3}) + 2g_{4,1}^{2,1}g_{4,1}^{2,2} + 2g_{4,1}^{2,1}g_{4,1}^{2,3} + 2g_{4,1}^{2,2}g_{4,1}^{2,3} \right) + \dots$$

shifted Gaussian fixed point (= interacting)

A lightning review of FRG results in 2 and 3 dimensions

rank 3 complex: [AE, Koslowski, '17]

cyclic melonic

different approximations for anomalous dimension

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^2$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^3$	
full	-1.94	0	—	—	—	—	quartic
semi-pert	-2.14	0	—	—	—	—	
pert	-4.62	0	—	—	—	—	
full	-1.37	0	-2.14	0	0	0	hexic
semi-pert	-1.47	0	-2.46	0	0	0	
pert	-2.14	0	-6.12	0	0	0	

	scheme	θ_1	θ_2	θ_3
quartic	full	2.21	-0.24	-0.93
	semi-pert	2	-0.21	-0.68
	pert	2	0.69	-2
hexic	full	2.14	-0.59	-1.26
	semi-pert	2	-0.58	-1.26
	pert	2	-0.35	-2

A lightning review of FRG results in 2 and 3 dimensions

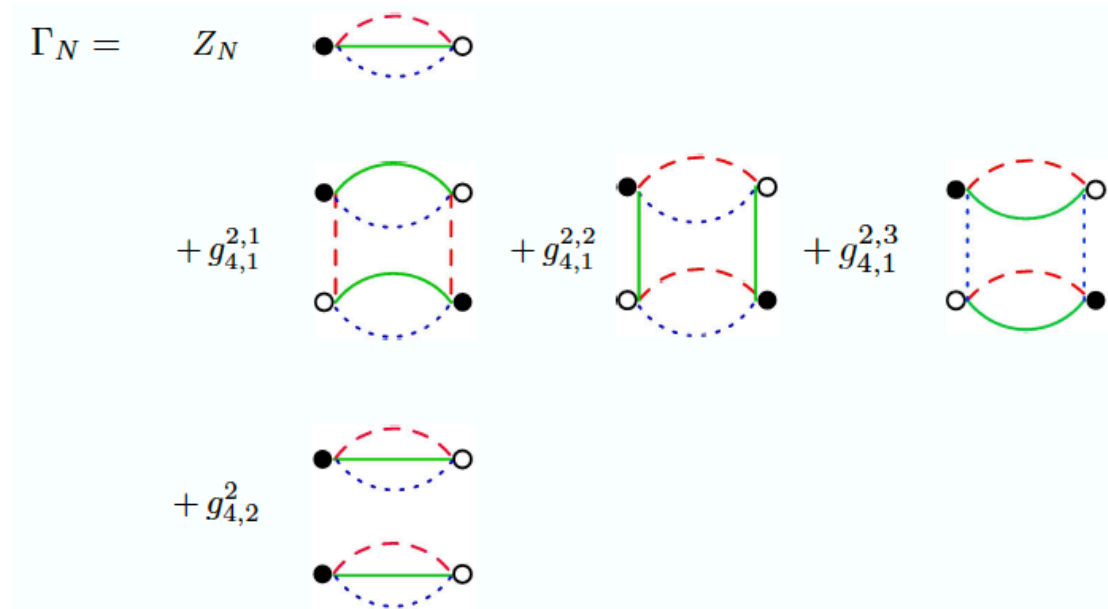
rank 3 complex: [AE, Koslowski, '17]

cyclic melonic (dimensional reduction)

different approximations for anomalous dimension

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^2$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^3$	
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A lightning review of FRG results in 2 and 3 dimensions

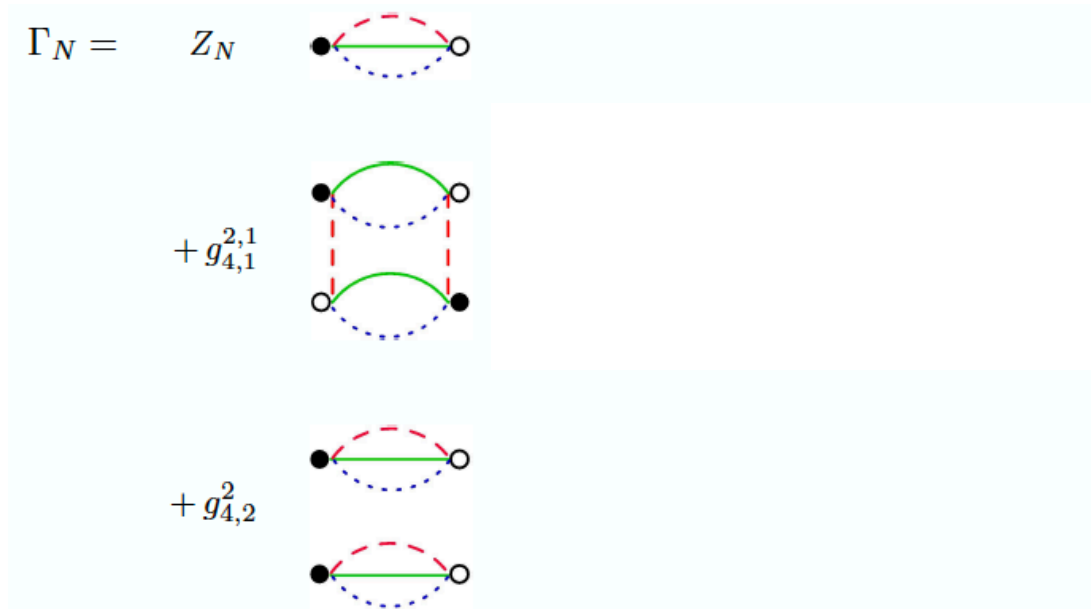
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	pert	2	-0.35	-2



matrix-model like

(matrix-model critical exponent from tensor model: AE, Koslowski, Lumma, Pereira: to appear)

A lightning review of FRG results in 2 and 3 dimensions

rank 3 complex: [AE, Koslowski, '17]

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	pert	2	-0.35	-2

cyclic melonic meets multitrace (dimensional reduction?)

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^2$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^3$
full	-1.05	-1.33	—	—	—	—
semi-pert	-1.58	-1.05	—	—	—	—
pert	-4.62	1.73	—	—	—	—
full	-0.53	-2.11	-0.14	0	-0.39	—
semi-pert	-0.63	-2.35	-0.20	0	-0.57	—
pert	-2.01	-1.64	-4.43	0	-4.84	—
full	-1.04	-0.84	-0.97	0	-0.64	-0.99
semi-pert	-1.15	-0.89	-1.21	0	-0.78	-1.14
pert	-2.10	-0.54	-5.54	0	-1.68	-0.47

scheme	θ_1	θ_2	θ_3
full	2.64	0.14	-0.65
semi-pert	2.28	0.14	-0.68
pert	2	-0.69	-2
full	3.10	0.25	-0.48
semi-pert	2.68	0.26	-0.51
pert	2.11	0.34	-2
full	2.56	0.44	-0.97
semi-pert	2.30	0.44	-0.98
pert	2.04	0.37	-2

A lightning review of FRG results in 2 and 3 dimensions

rank 3 complex: [AE, Koslowski, '17]

cyclic melonic (dimensional reduction)

different approximations for anomalous dimension

scheme	$g_{4,1}^{2,1}$	$g_{4,2}^2$	$g_{6,1}^{3,1}$	$g_{6,1}^{2,1}$	$g_{6,2}^{3,1}$	$g_{6,3}^3$	
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semi-pert	-1.58	-1.05	—	—	—	—
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full	-0.53	-2.11	-0.14	0	-0.39	—
semi-pert	-0.63	-2.35	-0.20	0	-0.57	—
pert	-2.01	-1.64	-4.43	0	-4.84	—
full	-1.04	-0.84	-0.97	0	-0.64	-0.99
semi-pert	-1.15	-0.89	-1.21	0	-0.78	-1.14
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pert	2.04	0.37	-2

rank 3 real: talk by Johannes Lumma

Summary

motivation: continuum limit generates predictivity

functional RG: tool to search for universal continuum limit

**local coarse graining
with background:**

local QFT of the metric

**non-local coarse graining
without background:**

tensor models

RG in background-independent case:

- **determines canonical scaling dimensions in large N limit**
- **reproduces benchmark results in 2d models**
- **provides candidates for universality classes in rank 3**

Outlook

- **rank 4 (real and complex)**
- **connection to CDTs/EDTs**
- **emergent geometries?**