Ultimate-precision quantum clocks: from mathematical characterization towards light-matter implementations

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- [M.W. Quantum 5, 381 (2021)] What is a clock?
- [M.W., Ralph Silva, Gilles Pütz, Sandra Stupar, Renato Renner, PRX Quantum (2019)] What is the most accurate clock?
- [Arman Dost and M.W., arXiv:2303.10029 (2023)]

Can we realise it?





[M.W. Quantum 5, 381 (2021)]



 $\mathcal{M}_{\mathrm{CR}_{\mathrm{T}}\rightarrow\mathrm{CR}_{\mathrm{T}}}^{t}(\,\cdot\,)$ \mathcal{H}_{N_T} \mathcal{H}_d

$ho_{\mathrm{C}}(t)\otimes\left|0 ight angle\!\left\langle 0 ight|_{\mathrm{R}_{\mathrm{T}}}$



 $\mathcal{M}_{\mathrm{CR}_{\mathrm{T}}\rightarrow\mathrm{CR}_{\mathrm{T}}}^{t}(\,\cdot\,)$ \mathcal{H}_{N_T} \mathcal{H}_d

 $\rho_{\rm C}(t) \otimes |1\rangle \langle 1|_{\rm R_T}$



 $\mathcal{M}_{\mathrm{CR}_{\mathrm{T}}\rightarrow\mathrm{CR}_{\mathrm{T}}}^{t}(\,\cdot\,)$ \mathcal{H}_{N_T} \mathcal{H}_d

 $\rho_{\rm C}(t) \otimes |2\rangle \langle 2|_{\rm R_T}$



 $\mathcal{M}_{\mathrm{CR}_{\mathrm{T}}\rightarrow\mathrm{CR}_{\mathrm{T}}}^{t}(\,\cdot\,)$ \mathcal{H}_{N_T} \mathcal{H}_d

 $ho_{\mathrm{C}}(t)\otimes\left|3
ight
angle\!\left\langle 3
ight|_{\mathrm{R}_{\mathrm{T}}}$



 $\mathcal{M}_{\mathrm{CR}_{\mathrm{T}} \to \mathrm{CR}_{\mathrm{T}}}^{t}(\,\cdot\,)$ \mathcal{H}_{N_T} \mathcal{H}_d

 $\rho_{\mathrm{CR}_{\mathrm{T}}}(t) = \sum_{j=0}^{N_{\mathrm{T}}} p_j(t) \, \rho_{\mathrm{C}_j}(t) \otimes |j\rangle \langle j|_{\mathrm{R}_{\mathrm{T}}}$

[M.W. Quantum 5, 381 (2021)]

1) Self-timing condition:

 $\mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t_1+t_2} = \mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t_2} \circ \mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t_1}$

2) Continuity:

$$\lim_{t \to 0^+} \left\| \mathcal{M}_{\mathrm{CR} \to \mathrm{CR}}^t - \mathcal{I}_{\mathrm{CR}} \right\| = 0$$

$$p(t, l|k) :=$$

$$\operatorname{tr} \left[\mathcal{M}_{\mathrm{CR} \to \mathrm{CR}}^{t}(\rho_{\mathrm{C}} \otimes |k\rangle\langle k|_{\mathrm{R}}) |l\rangle\langle l|_{\mathrm{R}} \right]$$

$$3) \operatorname{No tick-skipping:}$$

$$\lim_{t \to 0^{+}} \frac{\sum_{\substack{l=0\\l \notin \{k,k+1\}}}^{N_{T}} p(t, k+1)}{p(t, k+1)}$$

$$\sum_{\substack{l=0\\l\notin\{k,k+1\}}}^{N_T} p(t,l|k)$$

$$\lim_{t \to 0^+} \frac{p(t, k+1)}{p(t, k+1|k)} = 0$$

4) Clockwork translation invariance:

$$\operatorname{tr}_{\mathrm{R}}\left[\mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t}\left(\rho_{\mathrm{C}}\otimes|k\rangle\!\langle k|_{\mathrm{R}}\right)|k+l\rangle\!\langle k+l|_{\mathrm{R}}\right] \text{ is } k \text{ independent for all } k,l,t$$

[M.W. Quantum 5, 381 (2021)]

Dynamical semi-group representation:

The pair $\rho_{\rm C}^0, (\mathcal{M}_{\rm CR\to CR}^t)_{t\geq 0}$ satisfy 1) to 4), if and only if there exists a Hermitian operator H and linear operators $(L_j)_{j=1}^{N_L}$ and $(J_j)_{j=1}^{N_L}$; which are all time independent, such that for all $t \ge 0$:

$$\begin{split} \mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t}(\cdot) &= \ \mathrm{e}^{t\mathcal{L}_{\mathrm{CR}}}(\cdot), \\ \mathcal{L}_{\mathrm{CR}}(\cdot) &= -\mathrm{i}[\tilde{H},(\cdot)] + \sum_{j=1}^{N_{L}} \tilde{L}_{j}(\cdot)\tilde{L}_{j}^{\dagger} - \frac{1}{2} \{\tilde{L}_{j}^{\dagger}\tilde{L}_{j},(\cdot)\} \\ &+ \sum_{j=1}^{N_{L}} \tilde{J}_{j}(\cdot)\tilde{J}_{j}^{\dagger} - \frac{1}{2} \{\tilde{J}_{j}^{\dagger}\tilde{J}_{j},(\cdot)\}, \end{split}$$
where $\tilde{H} = H \otimes \mathbb{1}_{\mathrm{R}}$, $\tilde{L}_{j} = L_{j} \otimes \mathbb{1}_{\mathrm{R}}$, $\tilde{J}_{j} = J_{j} \otimes O_{\mathrm{R}}$
 $O_{\mathrm{R}} := |1\rangle\langle 0|_{\mathrm{R}} + |2\rangle\langle 1|_{\mathrm{R}} + |3\rangle\langle 2|_{\mathrm{R}} + \ldots +$

 $|N_T\rangle\langle N_T - 1|_{\rm B} + |0\rangle\langle N_T|_{\rm B}$.

Examples:



1) Thermodynamic "heat-engine" clock [P. Erker et. al. PRX (2017)]

$$L_{1} = \sqrt{\gamma_{h}}\sigma_{h}, \qquad L_{2} = \sqrt{\gamma_{h}}e^{-\beta_{h}E_{h}}\sigma_{h}^{\dagger},$$
$$L_{3} = \sqrt{\gamma_{c}}\sigma_{c}, \qquad L_{4} = \sqrt{\gamma_{c}}e^{-\beta_{c}E_{c}}\sigma_{c}^{\dagger},$$
$$J_{1} = \sqrt{\Gamma}|0\rangle\langle d-1|_{w}, \quad J_{2} = J_{3} = J_{4} = 0,$$

2) Ladder clock [S. Stupar et.al, arXiv:1806.08812 (2018)] $L_j = |c_{j+1}\rangle\langle c_j|, \quad J_j = 0, \quad j = 0, 1, 2, ..., d-1$ $L_d = 0, \qquad \qquad J_d = |c_1\rangle\langle c_d|,$



Quasi-Ideal clock [M.W. et al, PRX Quantum (2019)] $L_j = 0, \quad J_j = \sqrt{2V_j} |\psi_C\rangle\langle t_j|,$ j = 0, 1, 2, ..., d [A. Peres, AJP (1980)]

[M.W. et al, Annales Henri Poincaré (2018)] How well can a clock measure time? [P. Erker et. al. PRX (2017)]

• Delay function of 1^{st} tick $P_{tick}^{(1)}(t)$:

 $P_{\text{tick}}^{(1)}(t)dt := \text{probability of not ticking in interval} [0, t)$ *followed by ticking in time interval* [t, t + dt]

 $P_{\rm tick}^{(k)}(t)$

• Accuracy of kth tick: $R_k := \frac{\mu_k^2}{\sigma_k^2}$ For reset clocks: $R_k := kR_1$

Complete classification of classical clocks

- Definition: Classical clock: a quantum clock without coherence.
 - Classical clocks obey stochastic dynamics
- Theorem (Optimal classical clock) [M.W. et al, PRX Quantum (2019)]

All d dimensional classical clocks satisfy

 $R_k \le kd$

-- Ladder clock is classical and $R_k = kd$

-- Thermodynamic "heat-engine" clock is quantum but

 $R_k \propto kd$



Quantum clock Advantage

[M.W. et al, PRX Quantum (2019)] [M.W. et al, Annales Henri Poincaré (2018)]

-Theorem:

There exists parameters for the quasi-ideal clock such that

$$R_k = k d^2$$
 as $d o \infty$

-Theorem:

[Y. Yang et. al. ArXiv: 2004.07857 (2020)]

All quantum clocks have accuracy

 $R_1 \le 2e\pi d^2$

Paper with Arman PT Dost:



What would be needed for an experimental realization? [Arman Dost and M.W., arXiv:2303.10029 (2023)]

Why this matters:

- Is it physical? Are we cheating?
- What is the thermodynamic cost?



1) Thermodynamic "heat-engine" clock [P. Erker et. al. PRX (2017)]

What does Wall-E propose?

What does Wall-E propose?



A quantum photonic clock



A quantum clock emitting a photon into the void

Is the quantum advantage achievable in low dimensions?

[Arman Dost and M.W., arXiv:2303.10029 (2023)]



What we need

$$\mathcal{M}_{\mathrm{CR}\to\mathrm{CR}}^{t}(\cdot) = \mathrm{e}^{t\mathcal{L}_{\mathrm{CR}}}(\cdot),$$
$$\mathcal{L}_{\mathrm{CR}}(\cdot) = -\mathrm{i}[\tilde{H},(\cdot)] + \sum_{j=1}^{N_{L}} \tilde{L}_{j}(\cdot)\tilde{L}_{j}^{\dagger} - \frac{1}{2} \{\tilde{L}_{j}^{\dagger}\tilde{L}_{j},(\cdot)\}$$
$$+ \sum_{j=1}^{N_{L}} \tilde{J}_{j}(\cdot)\tilde{J}_{j}^{\dagger} - \frac{1}{2} \{\tilde{J}_{j}^{\dagger}\tilde{J}_{j},(\cdot)\},$$

$$\tilde{H} = H \otimes \mathbb{1}_{\mathbf{R}} \qquad \tilde{L}_j = L_j \otimes \mathbb{1}_{\mathbf{R}}
\tilde{J}_j = J_j \otimes O_{\mathbf{R}}$$

• Recall Quasi-ideal clock:

$$L_j = 0, \quad J_j = \sqrt{2V_j} |\psi_{\mathcal{C}}\rangle \langle t_j|,$$

$$j = 0, 1, 2, ..., d$$

Initial clockwork state (Reset clock)

$$\left|t_{j}\right\rangle = \frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} e^{-i2\pi nk/d} \left|E_{n}\right\rangle$$

$$\hat{H} = \sum_{n=0}^{d-1} \omega_0 n \left| E_n \right\rangle \! \left\langle E_n \right|$$

 $H_{\rm total} = H_{\rm C} + H_{\rm E} + H_{\rm CER}$





 $H_{\rm total} = H_{\rm C} + H_{\rm E} + H_{\rm CE}$ Matter Light (bath) $-\vec{D} \cdot \vec{E}$ Weak Coupling – Born-Markov approximation $\tau_B \ll \tau_R$: $\frac{d}{dt}\rho_{\rm C}(t) =$ $\sum e^{it(\omega'-\omega)} \Gamma(\omega) \Big(D(\omega)\rho_{\rm C}(t) D^{\dagger}(\omega') - D^{\dagger}(\omega') D(\omega)\rho_{\rm C}(t) \Big)$ +h.c., $|E_n\rangle \to |E_{n+\omega}\rangle$ \blacktriangleright R.W.A. $\tau_S \ll \tau_R$:

$$\frac{d}{dt}\rho_{\rm C}(t) = \sum_{\omega} \Gamma(\omega) \Big(D(\omega)\rho_{\rm C}(t)D^{\dagger}(\omega) - D^{\dagger}(\omega)D(\omega)\rho_{\rm C}(t) \Big) + \text{h.c.},$$

 $|\psi_{\rm C}\rangle\!\langle\psi_{\rm C}|\langle t_j|\rho_{\rm C}(t)|t_j\rangle - \left\{V_j|t_j\rangle\!\langle t_j|,\rho_{\rm C}(t)\right\}$

$$\left|t_{j}\right\rangle = \frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} \mathrm{e}^{-\mathrm{i}2\pi nk/d} \left|E_{n}\right\rangle$$



– R.W.A.



– Phase engineering?

$$\left|t_{j}\right\rangle = \frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} \mathrm{e}^{-\mathrm{i}2\pi nk/d} \left|E_{n}\right\rangle$$

[Arman Dost and M.W., arXiv:2303.10029 (2023)]

The matter Hamiltonian:



Top ring state:

$$\psi_{\rm top}(z,r,\phi) = \frac{1}{\sqrt{2\pi R}} \,\mathrm{e}^{-\mathrm{i}\,\phi j}$$

- Bottom ring states:
 - Spatial localization: $\psi_n(z,r,\phi)\psi_k(z',r',\phi') \approx 0$
 - Rotational symmetry:

 $\psi_1(z, r, \phi - n2\pi/d) = \psi_n(z, r, \phi)$

$$[D_z]_{\psi_n,\psi_{\text{top}}} = q \int dz \int r dr \int d\phi \,\psi_n(z,r,\phi) \, z \,\psi_{\text{top}}^*(z,r,\phi)$$
$$= q \int dz \int r dr \int d\phi \,\psi_1(z,r,\phi-n2\pi/d) \, z \frac{\mathrm{e}^{-\mathrm{i}\,\phi j}}{\sqrt{2\pi R}}$$
$$= \mathrm{e}^{-\mathrm{i}2\pi n j/d} \, [D_z]_{\psi_1,\psi_{\text{top}}}$$

[Arman Dost and M.W., arXiv:2303.10029 (2023)]

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[Arman Dost and M.W., arXiv:2303.10029 (2023)]

The matter Hamiltonian:



> What we get:

$$H_{\rm total} = H_{\rm C} + H_{\rm E} - D_z E_z$$

$$\frac{d}{dt}\rho_{\rm C}(t) = -\mathrm{i}[H_{\rm C}, \rho_{\rm C}(t)]$$

+2V_j |\psi_{\rm C}\langle\psi_{\rm C}| \langle t_j |\rho_{\rm C}(t)|t_j \rangle - \{V_j |t_j\rangle t_j|, \rho_{\rm C}(t)\}
|t_j \rangle = \frac{1}{\sqrt{d}} \sum_{n=0}^{d-1} e^{-\mathrm{i}2\pi nj/d} |E_n \rangle

[Arman Dost and M.W., arXiv:2303.10029 (2023)]

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 $\succ \text{ C.f. what we want:}$ $\frac{d}{dt}\rho_{\rm C}(0) = -\mathrm{i}[H_{\rm C}, \rho_{\rm C}(t)]$ $+ \sum_{j=0}^{d-1} 2V_j |\psi_{\rm C}\rangle \langle \psi_{\rm C}| \langle t_j | \rho_{\rm C}(t) | t_j \rangle - \{V_j | t_j \rangle \langle t_j |, \rho_{\rm C}(t)\}$





How to construct 2 decay channels:



 $\frac{d}{dt}\rho_{\rm C}(0) = -\mathrm{i}[H_{\rm C}, \rho_{\rm C}(t)] + \sum_{j=\{j_1, j_2\}} 2V_j |\psi_{\rm C}\rangle\langle\psi_{\rm C}|\langle t_j|\rho_{\rm C}(t)|t_j\rangle - \{V_j |t_j\rangle\langle t_j|, \rho_{\rm C}(t)\}$



1) Thermodynamic "heat-engine" clock [P. Erker et. al. PRX (2017)]



$$\Delta S_{\text{tick}} \approx \beta_v (Q_h - Q_c)$$
$$Q_h := (d - 1)E_h$$
$$Q_c := (d - 1)E_c$$

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$$Q_h := (d - 1)E_h$$
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$$R = \frac{\Delta S_{\text{tick}}}{2}$$









$$\sigma(\rho(t)) := -\frac{d}{dt} S(\rho(t)||\tau_{\beta}) = \frac{d}{dt} S(\rho(t)) + J(\rho(t))$$
$$J(\rho) := \beta \operatorname{tr}[H\mathcal{D}\rho]$$
$$= \operatorname{tr}[\mathcal{L}(\rho)\ln(\tau_{\beta})]$$

Entropy production per tick?

 $\operatorname{tr}[\mathcal{L}_{\mathrm{C}}^{\mathrm{nt}}(\rho_{\mathrm{C}}^{\mathrm{nt}}(\tilde{t}))\ln(\tau_{\beta})]$





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$$\int_0^t \mathrm{d}\tilde{t} \, \mathrm{tr}[\mathcal{L}_{\mathrm{C}}^{\mathrm{nt}}(\rho_{\mathrm{C}}^{\mathrm{nt}}(\tilde{t})) \ln(\tau_\beta)]$$





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$$P_{\text{tick}}(t) \int_0^t \mathrm{d}\tilde{t} \, \text{tr}[\mathcal{L}_{\mathrm{C}}^{\mathrm{nt}}(\rho_{\mathrm{C}}^{\mathrm{nt}}(\tilde{t})) \ln(\tau_\beta)]$$





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$$\Delta S_{\text{tick}} := \int_0^\infty \mathrm{d}t \, P_{\text{tick}}(t) \int_0^t \mathrm{d}\tilde{t} \, \text{tr}[\mathcal{L}_{\mathrm{C}}^{\mathrm{nt}}(\rho_{\mathrm{C}}^{\mathrm{nt}}(\tilde{t})) \ln(\tau_\beta)]$$



$$\mathcal{L}(\cdot) = [H, \cdot] + \mathcal{D}(\cdot)$$

$$\sigma(\rho(t)) := -\frac{d}{dt} S(\rho(t)||\tau_{\beta}) = \frac{d}{dt} S(\rho(t)) + J(\rho(t))$$
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$$R = \frac{\Delta S_{\text{tick}}}{2}$$







$$R \approx \left(\frac{\Delta S_{\text{tick}}}{\beta} - \omega_0\right)^2 \frac{4}{\omega^2}$$





$$R \approx \left(\frac{\Delta S_{\text{tick}}}{\beta} - \omega_0\right)^2 \frac{4}{\omega^2}$$

Minimal entropy per tick quadratically improved!

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Is the most accurate quantum clock realizable?

- Weak coupling limit sufficient for quantum advantage
- Up to $d\!=\!8$ only require 2 decay channels
- Discreate Fourier transform realizable
 with flux loops
- Quantum entropy advantage too
- Experimental implementation via superconducting qubits?

QInfo Inria team

Location: ENS Lyon and Uni. Grenoble

Faculty members

- Alastair Abbott
- Guillaume Aubrun
- Omar Fawzi
- Daniel Stilck França
- Mischa Woods

Postdoctoral researchers

PhD students

- Cyril Elouard
- Mizanur Rahaman
- Ala Shayeghi

- Emily Beatty
- Paul Fermé
- Aadil Oufkir
- Hoang-Duy Ta

Multiple PhD and postdoc positions