

> Panorama B



> Fluxes and applications



Plan

Plan

> Flux compactifications

Plan

- > Flux compactifications
- > Moduli stabilization

Plan

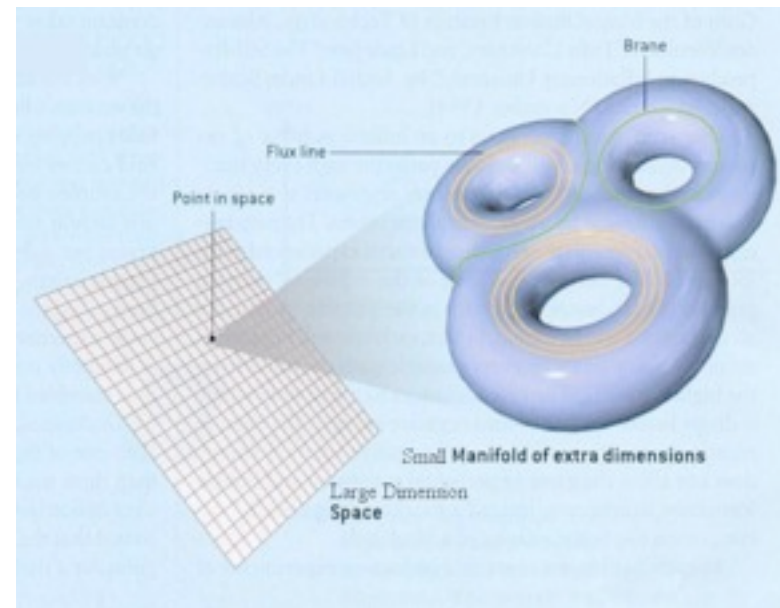
- > Flux compactifications
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- > Supersymmetry breaking

Plan

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- > Inflation

Moduli stabilization

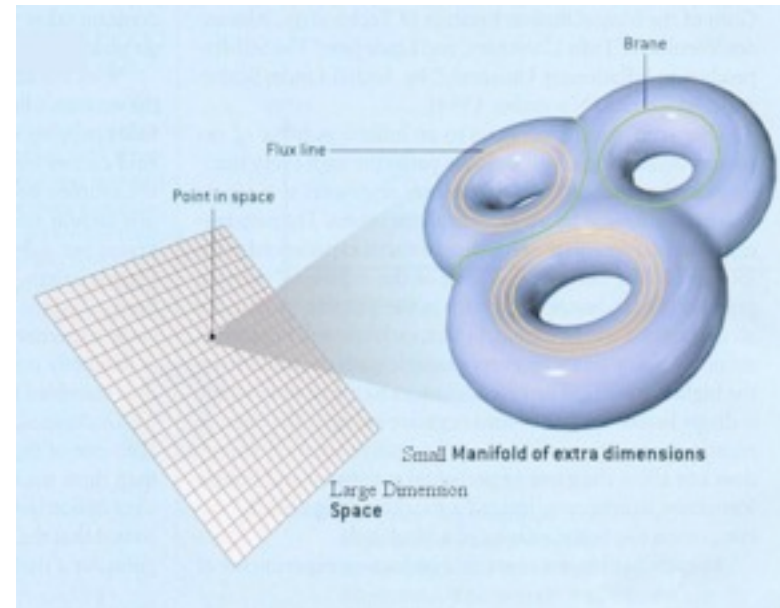
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- 📌 Cosmological problems unless massive enough
- 📌 Need mechanisms for moduli stabilization



Moduli stabilization

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Focus on flux
compactifications



Flux compactification

- 🔌 In addition to metric background,
introduce backgrounds for NSNS and RR p-form fields
- Due to gauge invariance, backgrounds for field strength

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 Topological sector defined by cohomology class

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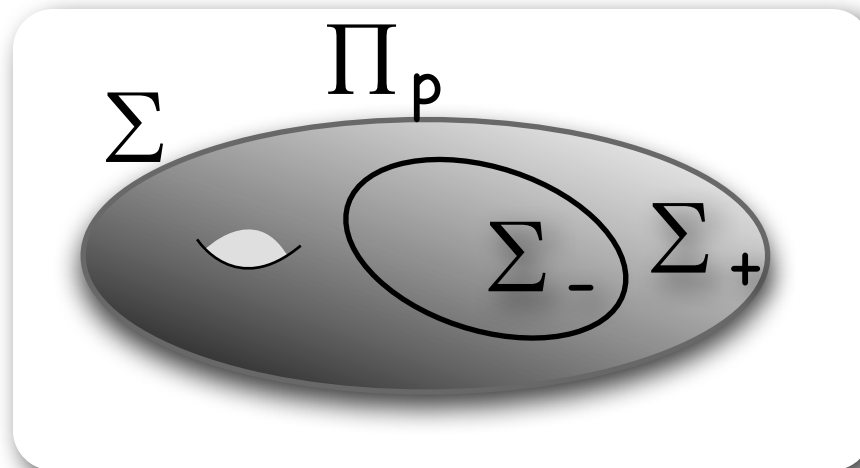
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📌 Topological sector defined by cohomology class

In the absence of sources, $dF_{p+1}=0$

Must specify fluxes over basis of cycles $\int_{\Lambda_k} F_{p+1} = N_k$

📌 Flux quantization



Flux compactification

 Fluxes introduce moduli dependence in potential energy

Closed string moduli are stabilized

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Closed string moduli are stabilized

- 📌 Dating back to Freund-Rubin

Ex: $AdS_5 \times S^5$

S^5 volume not a modulus, but sits at a minimum of potential

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Vacua need not be Minkowski, but also AdS (or dS??)

Focus on M_4 , or in (A)dS₄ with hierarchycal length scales

Type IIB with 3-form fluxes



Prototypical example: type IIB with NSNS and RR 3-form flux

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 10d action

$$S_{\text{IIB}} = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g} \left(R - \frac{\partial_M \bar{\tau} \partial^M \tau}{2(\text{Im } \tau)^2} - \frac{1}{2} |F_1|^2 - \frac{|G_3|^2}{2 \text{Im } \tau} - \frac{1}{4} |\tilde{F}_5|^2 \right) \\ + \frac{1}{2\kappa_{10}^2} \int C_4 \wedge \frac{G_3 \wedge \bar{G}_3}{4i \text{Im } \tau} + S_{\text{local}} .$$

Useful combination $G_3 = F_3 - \tau H_3$; $\tau = C_0 + ie^{-\phi}$

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Specify integrals, and local sources, satisfying RR tadpole condition

$$Q_{\text{flux}} + Q_{\text{D3}} + Q_{\text{O3}} = 0$$

Type IIB with 3-form fluxes



Flux “tension” and charge

$$\mathcal{L}_G = -\frac{1}{24\kappa_{10}^2} \int_{\mathbf{X}_6} d^6y \, g^{\frac{1}{2}} \frac{(G_3)_{mnp} (\bar{G}_3)^{mnp}}{\text{Im } \tau} = -\frac{1}{8\kappa_{10}^2} \int_{\mathbf{X}_6} d^6y \frac{G_3 \wedge *_6 \bar{G}_3}{\text{Im } \tau}$$

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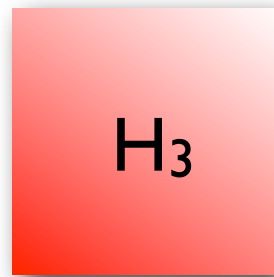
 Stabilizes complex structure moduli (and dilaton)

Type IIB with 3-form fluxes

$$F_3 = H_3$$



T^3



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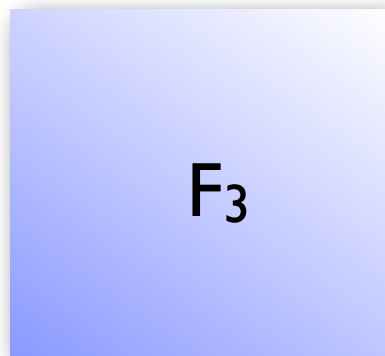
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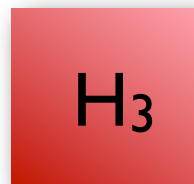
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Warped geometries



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$$ds^2 = Z(x^m)^{1/2} ds_{4d}^2 + Z(x^m)^{1/2} g_{mn}^{CY} dx^m dx^n$$

$$F_5 = (1 + *_{10d}) dZ^{-1} dx^0 dx^1 dx^2 dx^3$$

$$\nabla^2 Z(x^m) = g_s |G_3|^2 + g_s \sum \delta(D3/O3)$$

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Intuition: ISD flux works like “effective” D3-brane

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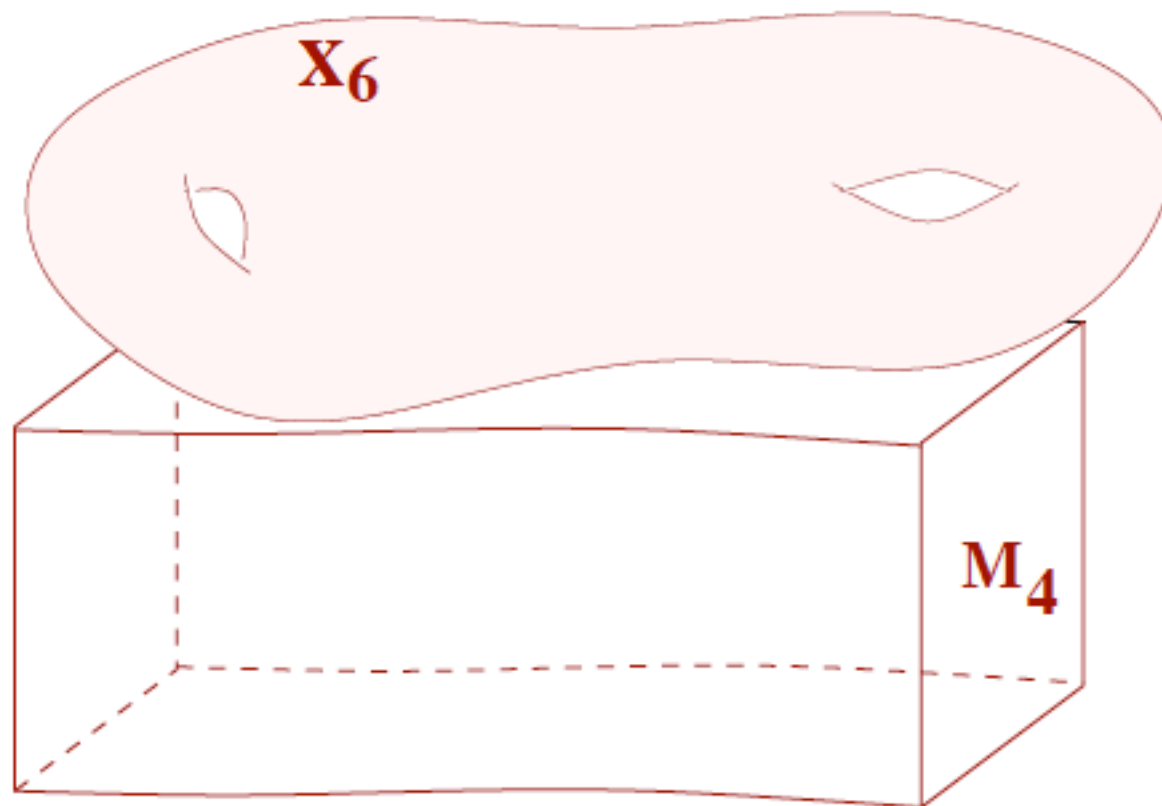
 Underlying CY makes life easy \Rightarrow Many explicit models

Flux superpotential

 Effect of flux on closed string moduli can be described in low energy effective field theory by W_{flux}

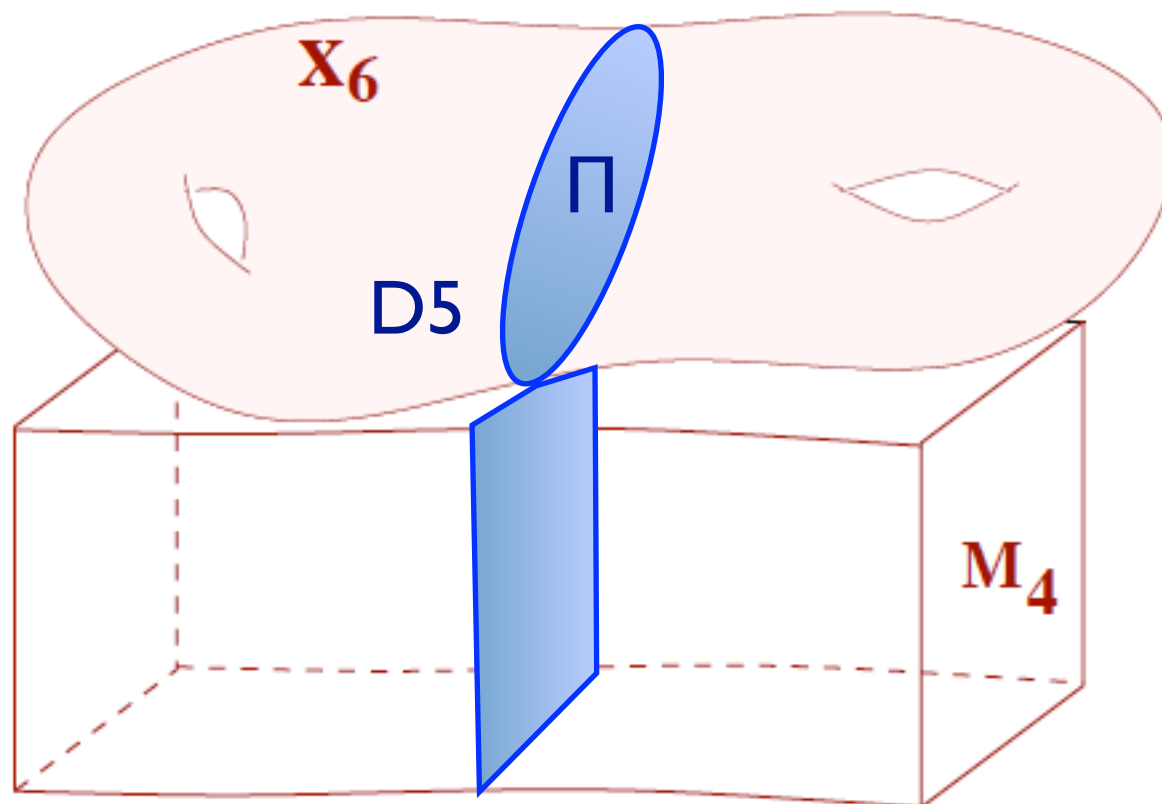
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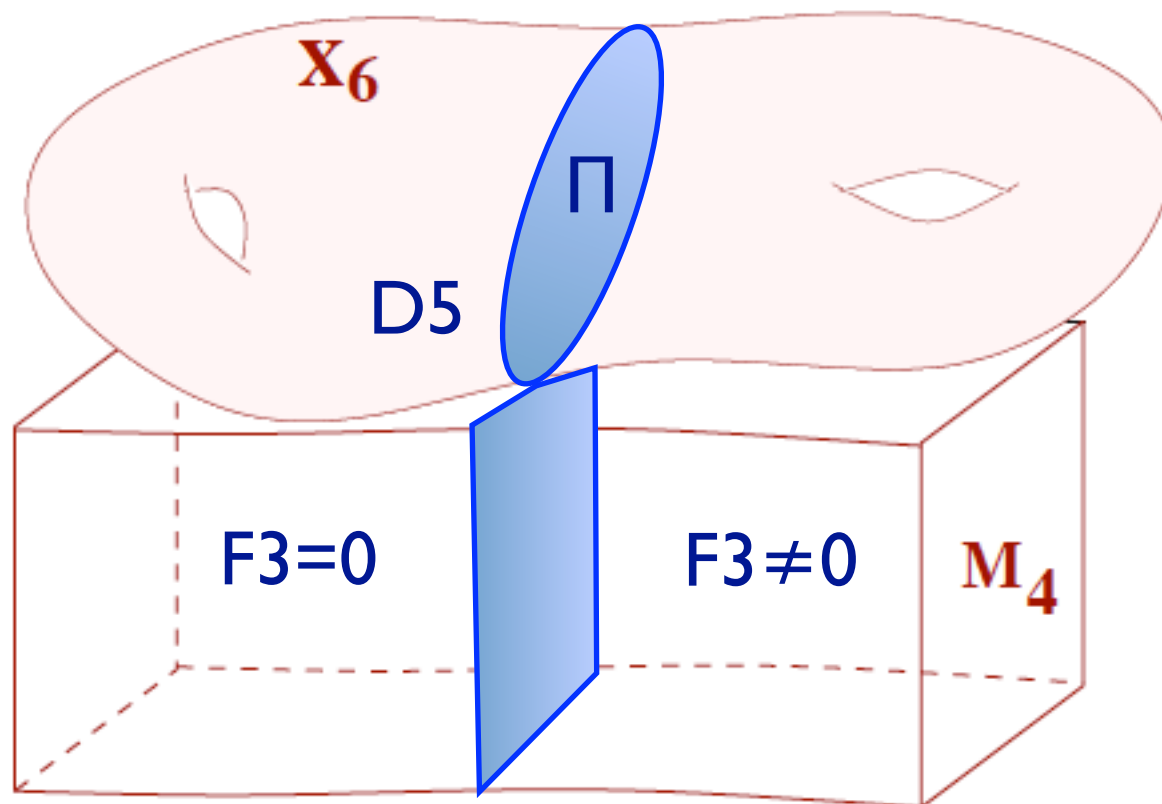
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- Compute W_{flux} from tension of domain wall introducing flux



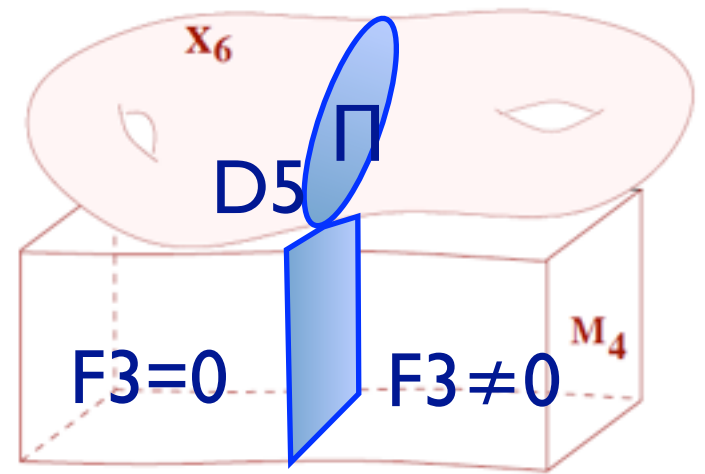
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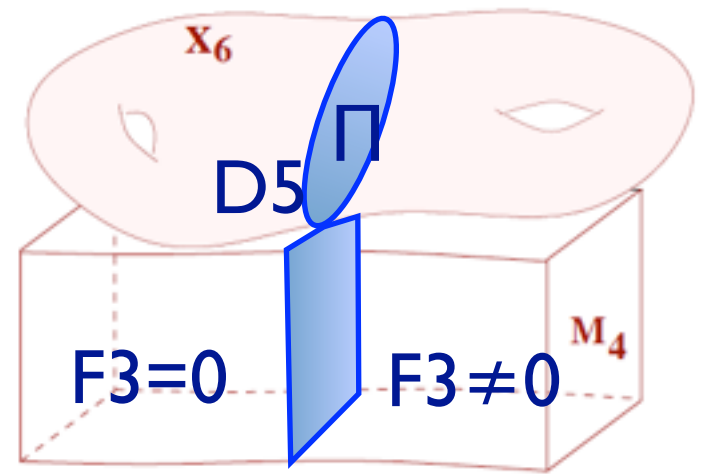


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D5 is magnetic source for C2

$$dF_3 = \delta_3(\Pi) \wedge \delta_0(x^3) dx^3,$$

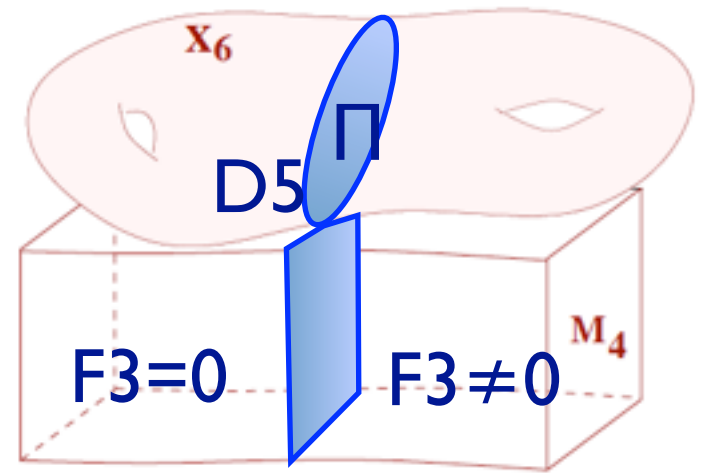


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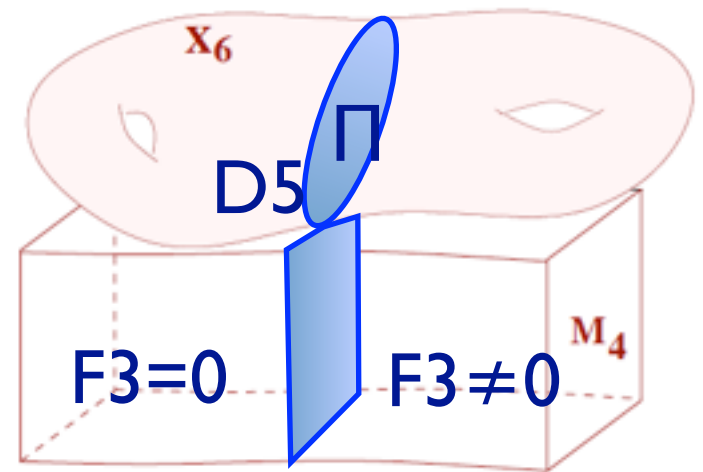
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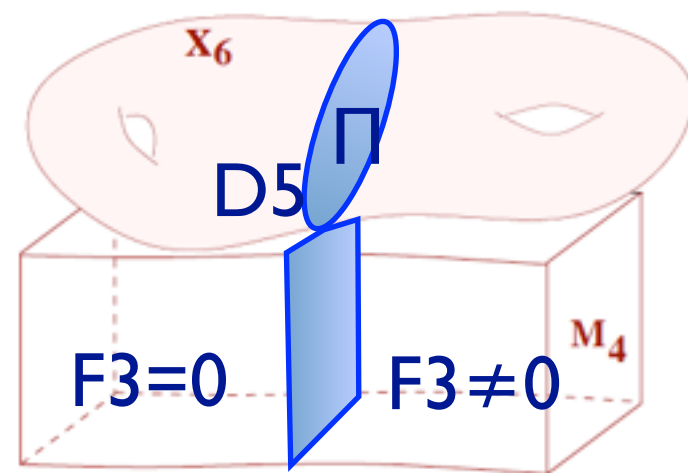
$$W_{F_3} = \int_{\Pi} \Omega_3 = \int_{\mathbf{x}_6} F_3 \wedge \Omega_3$$

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$$W_{\text{GVW}} = \int_{\mathbf{x}_6} G_3 \wedge \Omega_3$$

For general F_3 , H_3 fluxes

Supersymmetry

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 Also from 4d effective theory

$$K = -\log(\int_X \Omega \wedge \bar{\Omega}) - \log(\tau - \bar{\tau}) - 3\log(\rho + \bar{\rho})$$

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Scalar potential $V = e^K (g^{a\bar{b}} D_a W D_{\bar{b}} \bar{W} - 3|W|^2)$

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Vacuum: $D_i W_{\text{flux}} = 0$ $D_\tau W \sim \int_X \bar{G}_3 \wedge \Omega$ $G_3|_{(3,0)} = 0$

G3 is ISD $D_{z_i} W \sim \int_X G_3 \wedge \chi_{(2,1),i}$ $G_3|_{(1,2)} = 0$

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SUSY: G3 is (2,1) $D_\rho W \sim \int_X G_3 \wedge \Omega$ $G_3|_{(0,3)} = 0$

F-theory description



M-theory on elliptic CY4 in limit of vanishing fiber size

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Vacuum: G_4 is SD

Susy: G_4 is (2,2)

Towards the SM

MSSM from magnetized D7s

N_α	(n_α^1, m_α^1)	(n_α^2, m_α^2)	(n_α^3, m_α^3)	
$N_a = 3$	(0,1)	(3,1)	(-3,1)	D7 ₁
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(need few extra branes, adding few extra matter)



Flash
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Susy

$$G_3 = \frac{8}{\sqrt{3}} e^{-\frac{\pi i}{6}} (d\bar{z}_1 dz_2 dz_3 + dz_1 d\bar{z}_2 dz_3 + dz_1 dz_2 d\bar{z}_3)$$

Stabilizes at $\tau_1 = \tau_2 = \tau_3 = \tau = e^{2\pi i/3}$

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Non-susy

$$G_3 = 2(d\bar{z}_1 dz_2 dz_3 + dz_1 d\bar{z}_2 dz_3 + dz_1 dz_2 d\bar{z}_3 + d\bar{z}_1 d\bar{z}_2 d\bar{z}_3)$$

Stabilizes e.g. at $\tau_1 = \tau_2 = \tau_3 = \tau = i$

Full moduli stabilization & deSitter

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Corrections

Full moduli stabilization & deSitter



Corrections

- Earlier no-scale structure disappears upon including corrections
- Perturbative and non-perturbative, in α' and g_s
- Corrections may be small in large volume, small g_s regime, but not compared to zero!

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Full moduli stabilization & deSitter



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Generalized fluxes

- Above set of fluxes is not really the most general
- Geometric, non-geometric, U-dual fluxes
- Superpotentials depending on all moduli

Full moduli stabilization & deSitter

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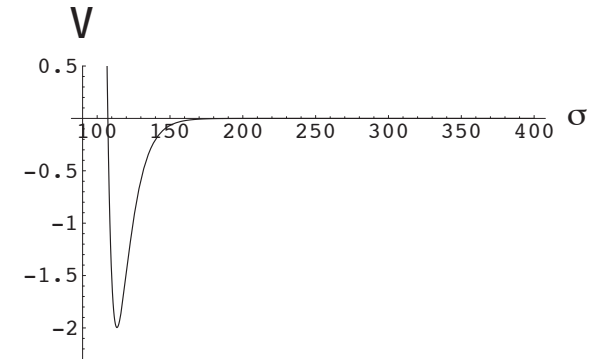
Use non-perturbative effects from D3-brane instantons

One-modulus toy model

$$W = W_{0,\text{flux}} + A(z_0)e^{-T}$$

Susy AdS vacua with stabilized moduli

[Kachru, Kallosh, Linde, Trivedi]



Full moduli stabilization & deSitter

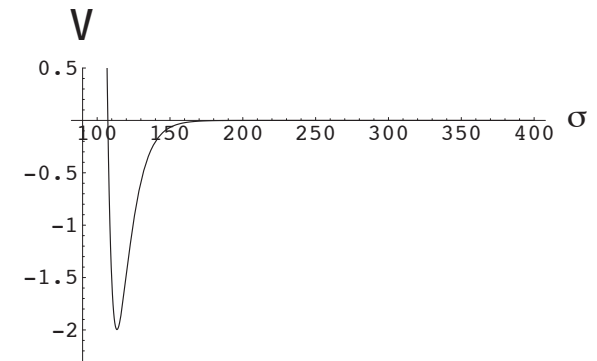
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- Can be relaxed in other scenarios (Large Volume Stabilization)

Full moduli stabilization & deSitter

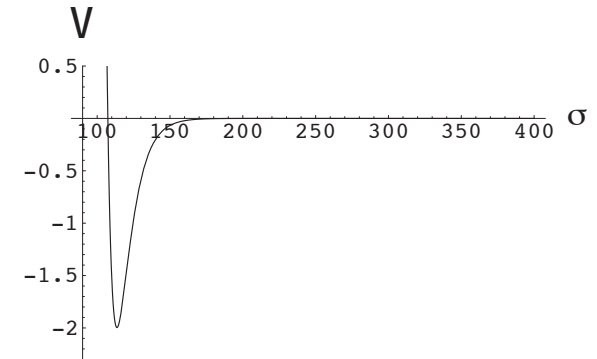
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[Kachru, Kallosh, Linde, Trivedi]



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Can be relaxed in other scenarios (Large Volume Stabilization)

- 📌 Generalization: D3-instantons for each independent Kahler modulus?

Full moduli stabilization & deSitter

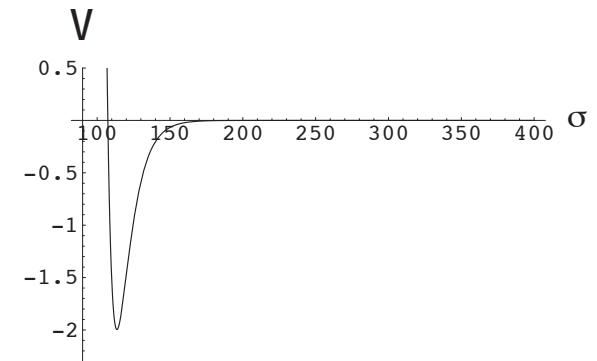
- 📌 Use non-perturbative effects from D3-brane instantons

One-modulus toy model

$$W = W_{0,\text{flux}} + A(z_0)e^{-T}$$

Susy AdS vacua with stabilized moduli

[Kachru, Kallosh, Linde, Trivedi]



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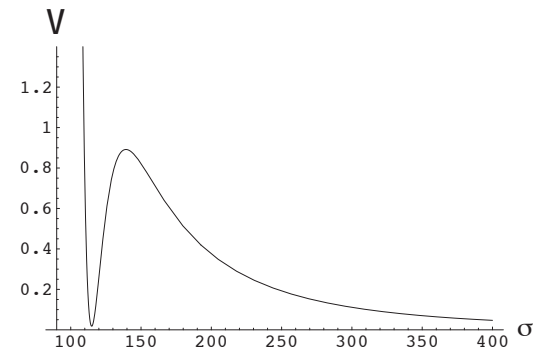
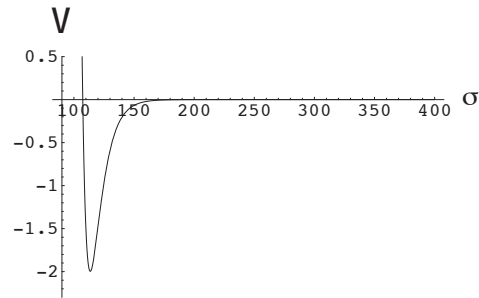
- 📌 Clash with SM arising from D7's

Intersections of instanton with D7's lead to susy involving SM fields

$$W = e^{-T} \rightarrow W = e^{-T} \Phi_1 \dots \Phi_n$$

Full moduli stabilization & deSitter

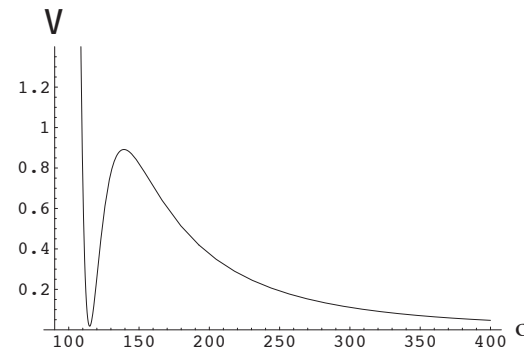
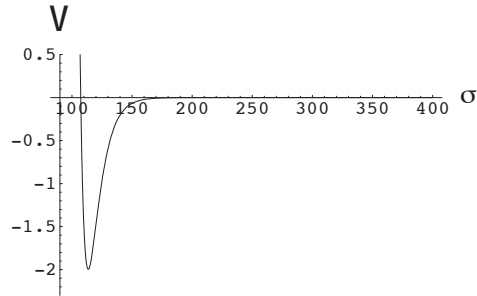
 Proposal to add sources of extra tension for uplifting to deSitter



- Anti D3-branes

Full moduli stabilization & deSitter

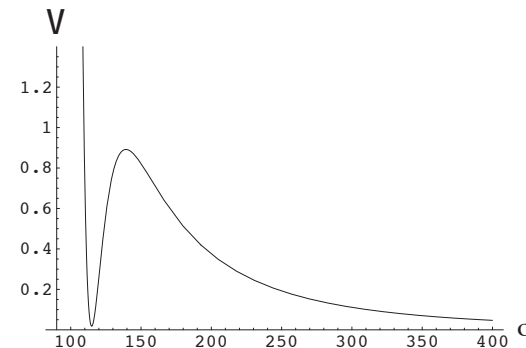
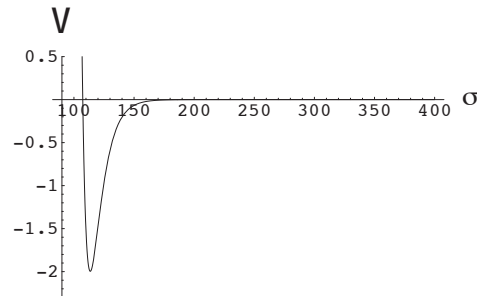
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- Anti D3-branes
- D-terms: anti-instantons on D7s
- F-terms for fluxes
- DSB sectors
- D-terms: anti-instantons on D7s

Full moduli stabilization & deSitter

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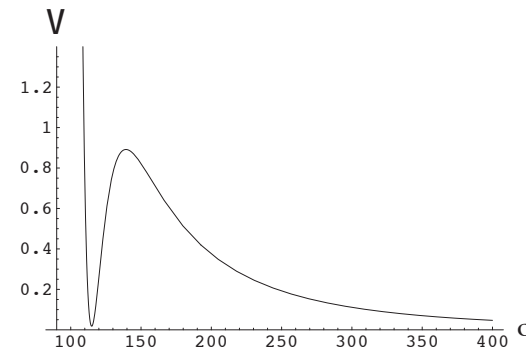
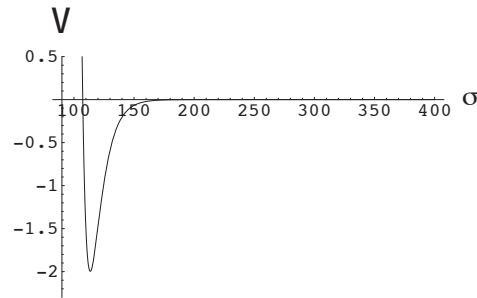


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 Hard to make anything explicit enough

Full moduli stabilization & deSitter

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 Hard to make anything explicit enough (new: nilpotent goldstino)

Generalized fluxes

- 📌 Flux compactifications in type IIA involve even and odd forms
⇒ Talk both to Kahler and complex structure moduli!

What about mirror symmetry?

- 📌 Using T-duality in local T^3 fibration, H_3 turns e.g. into geometric twist
⇒ Geometric fluxes

- 📌 Compactification on non-CY geometries, possibly 4d $N=1$

$SU(3)$ holonomy → $SU(3)$ structure

⇒ Generalized complex geometry, mirror symmetry, ...

- 📌 Painful lack of explicit compact examples

Generalized fluxes

 Generalized geometric and non-geometric fluxes from T-duality

Regard T^3 as T^2 (trivially) fibered over S^1

H_3 is monodromy $b \rightarrow b+1$ for $b = \int_{T^2} B$

Particular $SL(2, \mathbb{Z})$ monodromy on $T = A_{T^2} + ib$

One T-duality along T^2 gives $\tau \rightarrow \tau + 1$ in $SL(2, \mathbb{Z})$ of τ

\Rightarrow Geometric twisting, geometric flux

One S^1 non-trivially fibered over two directions $\omega^a{}_{bc}$

Two T-dualities give non-geometric $SL(2, \mathbb{Z})$ monodromy on T

\Rightarrow Non-geometric twisting, non-geometric flux $Q^{ab}{}_c$

Full T-duality covariance suggests

\Rightarrow locally non-geometric R^{abc}

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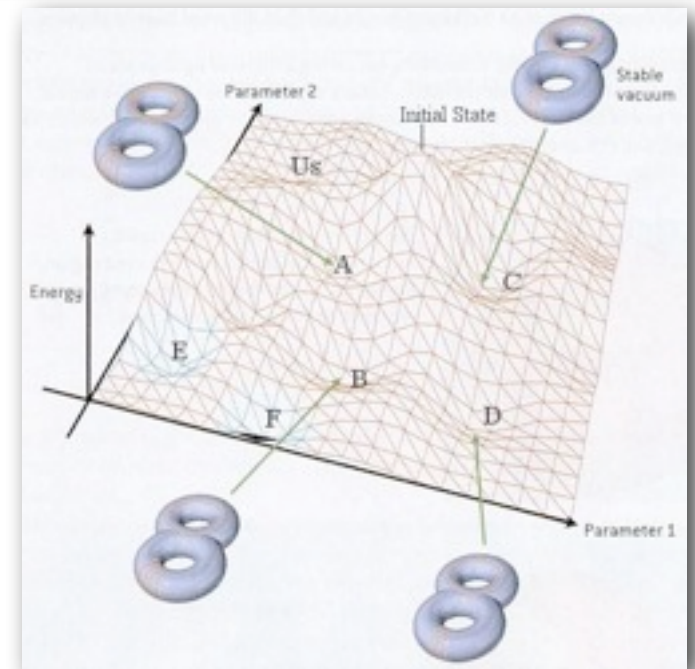
Generalized geometry. Double (&exceptional) field theory, ...

Flux landscape

 The general picture is compelling enough

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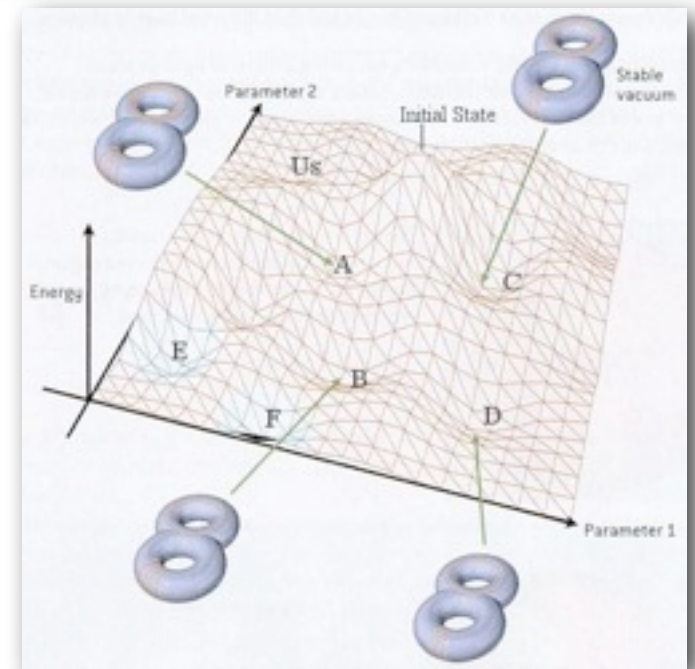


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Internal data determine 4d physics

symmetries, spectrum, couplings



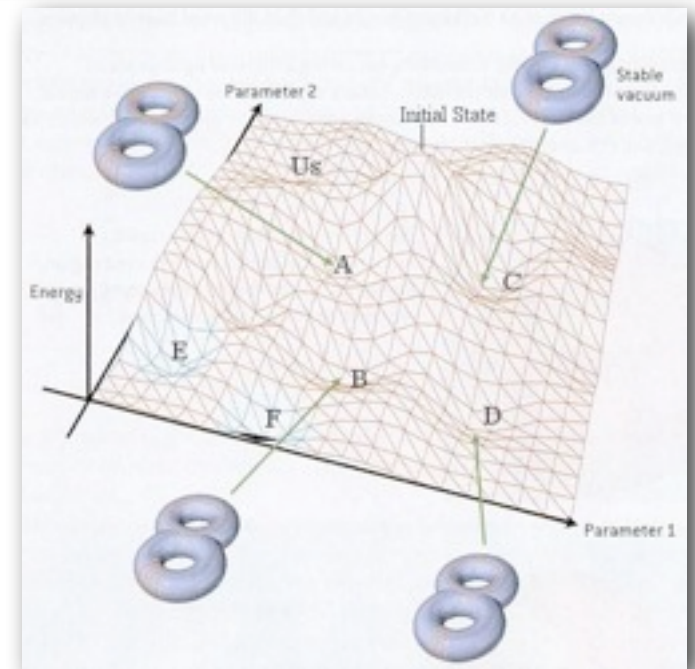
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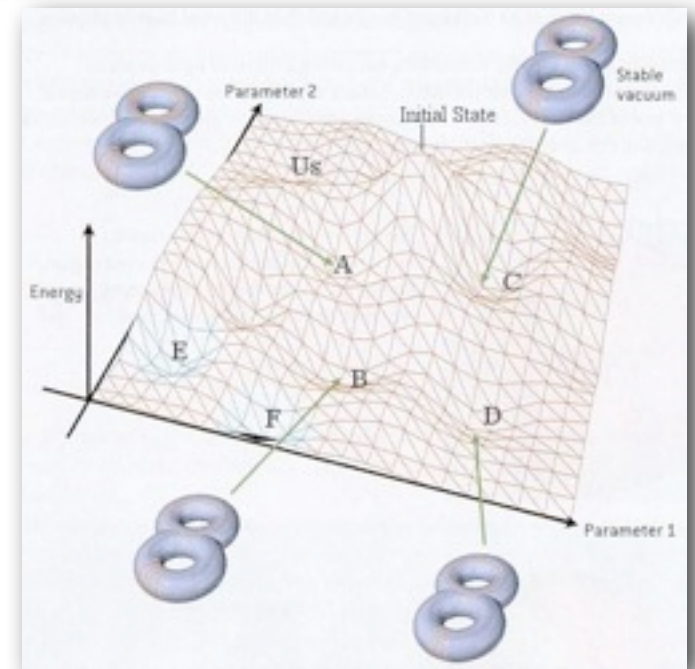
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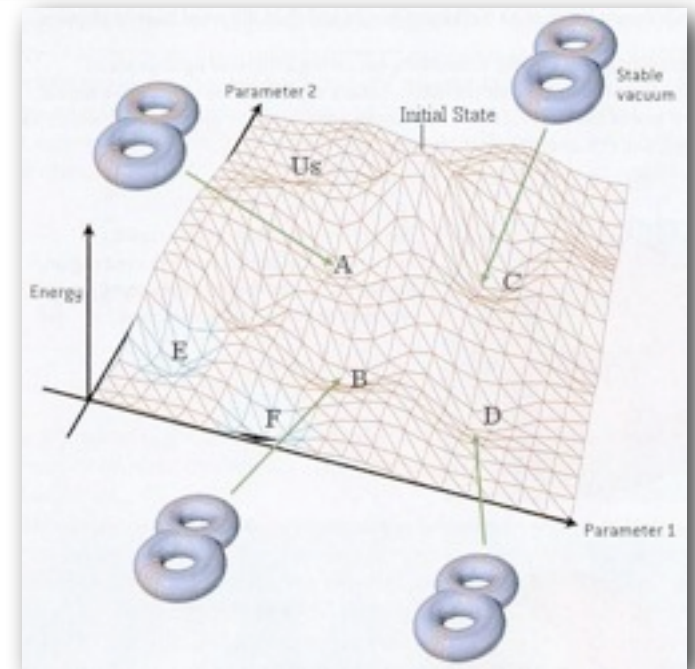
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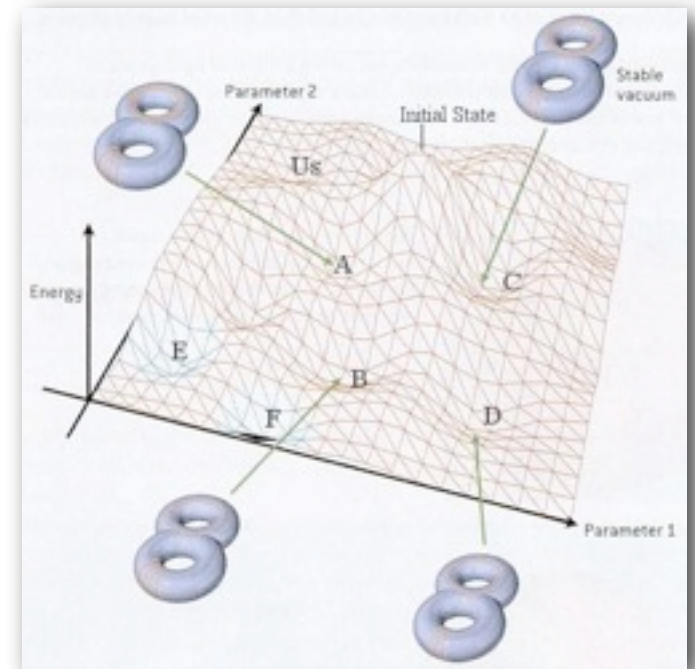
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Revision of “naturalness”: cosmological constant, hierarchy,...

Landscape of viewpoints

Landscape of viewpoints



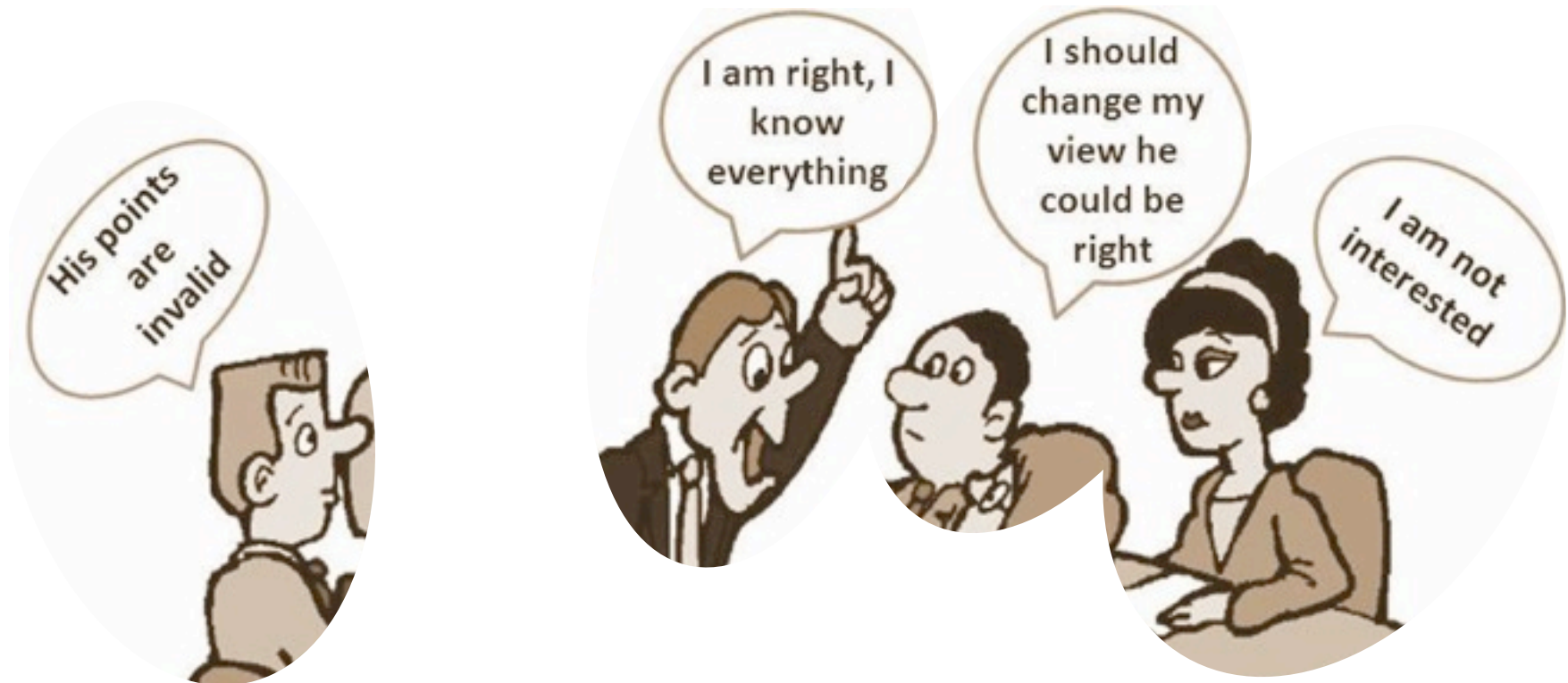
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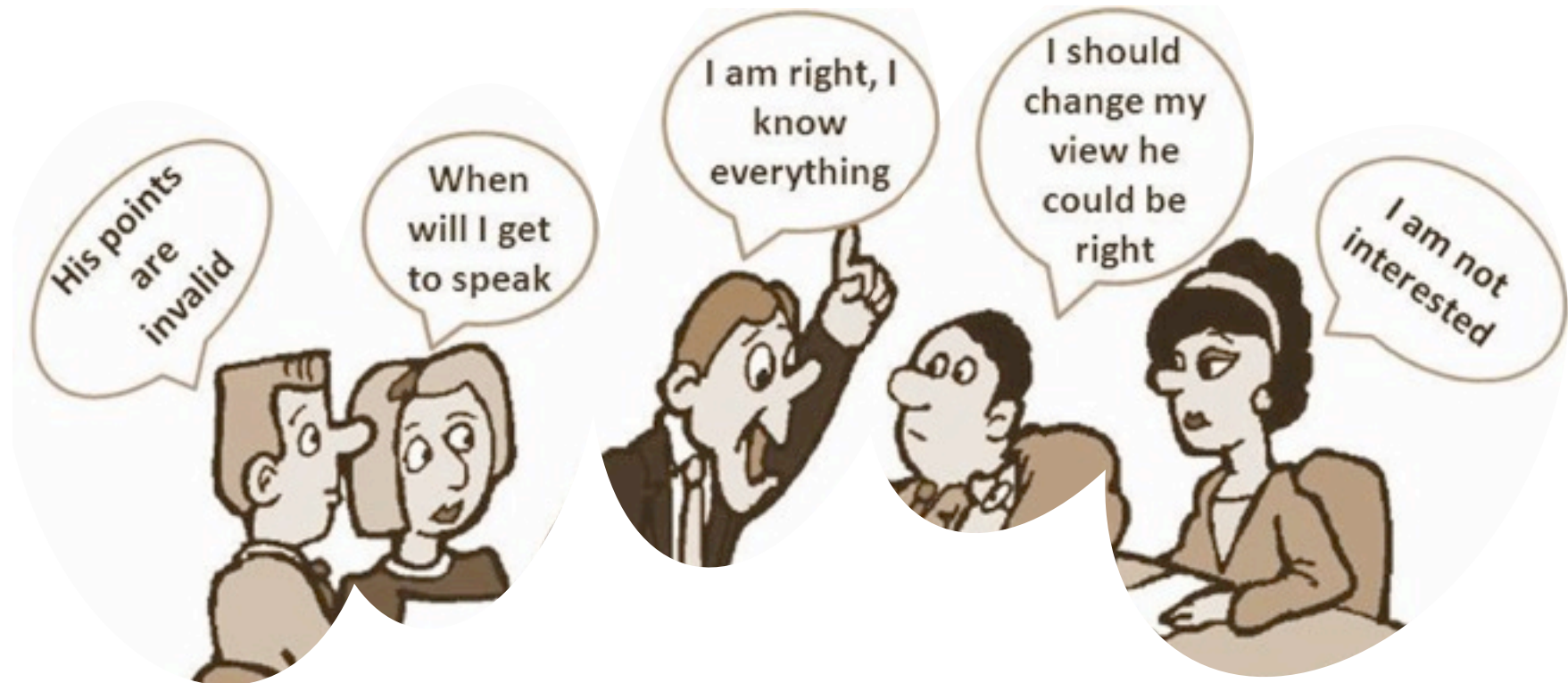
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Yeah, but...

Yeah, but...



What about some real Physics?

Landscape of viewpoints



What about some real Physics?





Keep rocking

After all, we are dealing with a theory
which is much more clever than any of us



Fluxes and D-branes

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Must consider the interplay of different ingredients

Fluxes and D-branes

-  Must consider the interplay of different ingredients
-  Focus on inter-relations between (field strength) fluxes and D-branes

Fluxes and D-branes



- 📌 Must consider the interplay of different ingredients
- 📌 Focus on inter-relations between (field strength) fluxes and D-branes
- 📌 Effects at several levels
 - Topological: Freed-Witten consistency conditions
 - (Susy) Open string moduli stabilization
 - Susy breaking

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Freed-Witten consistency conditions

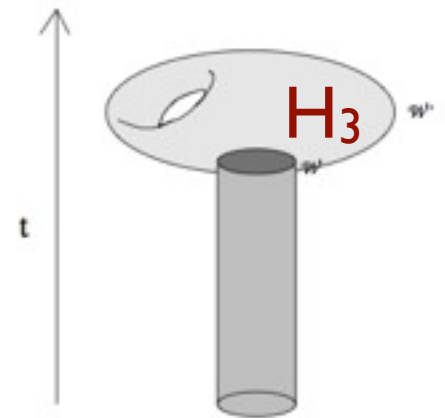
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Take D6-brane on 3-cycle with non-trivial H_3

$$\int_{D6} H_3 \wedge \tilde{A}_4$$

Flux is magnetic source for D6 U(1)

Must be cancelled by boundaries of outgoing D4's



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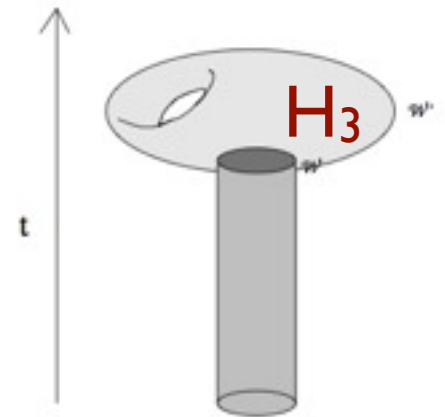
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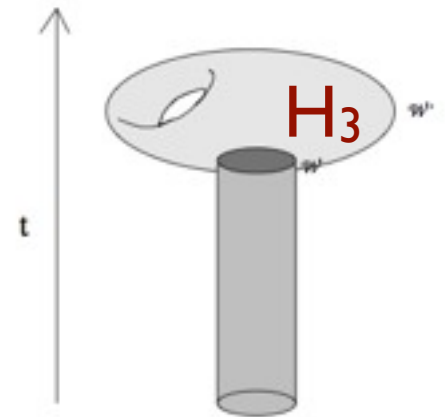
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- 📌 Analogous statements for RR fluxes

Freed-Witten and discrete symmetries

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Discrete gauge symmetries from flux catalysis

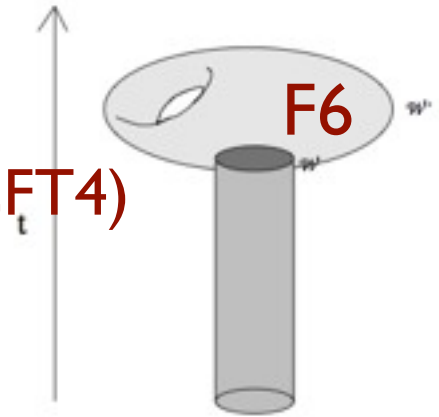
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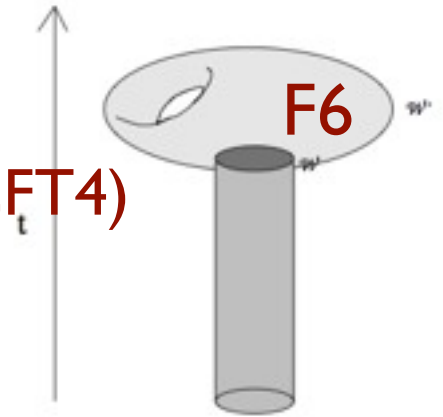
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Junction: D6 on X_6 . Emits k F1's (\mathbb{Z}_k strings)

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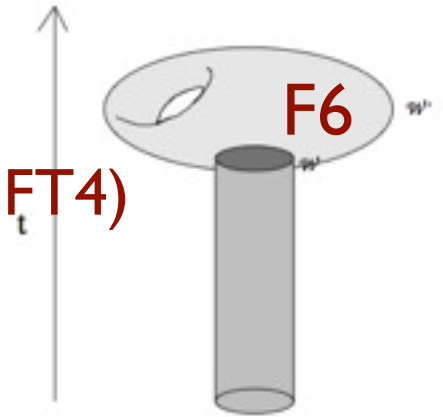
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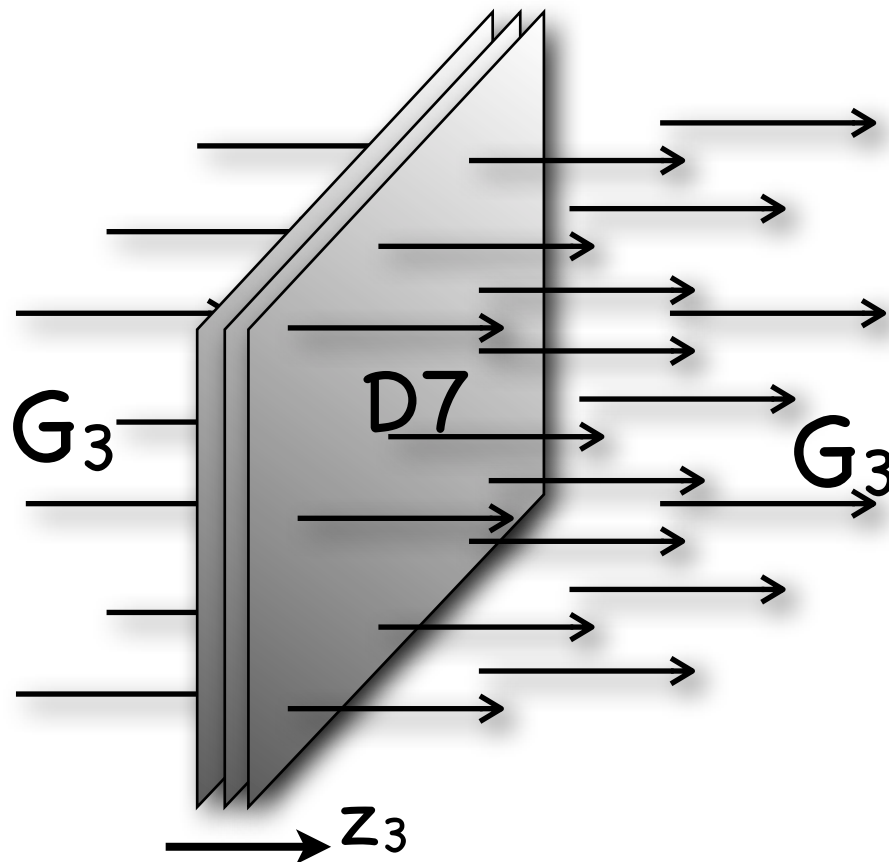
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📌 Can also do \mathbb{Z}_k domain walls. See later

Fluxes, susy breaking and soft terms

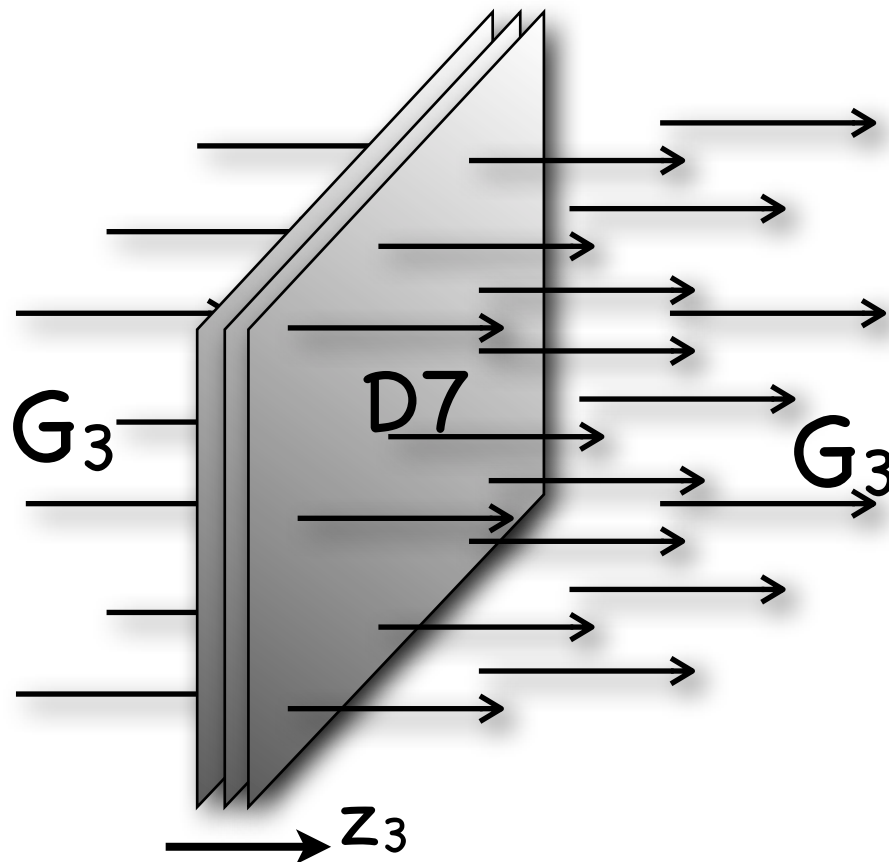
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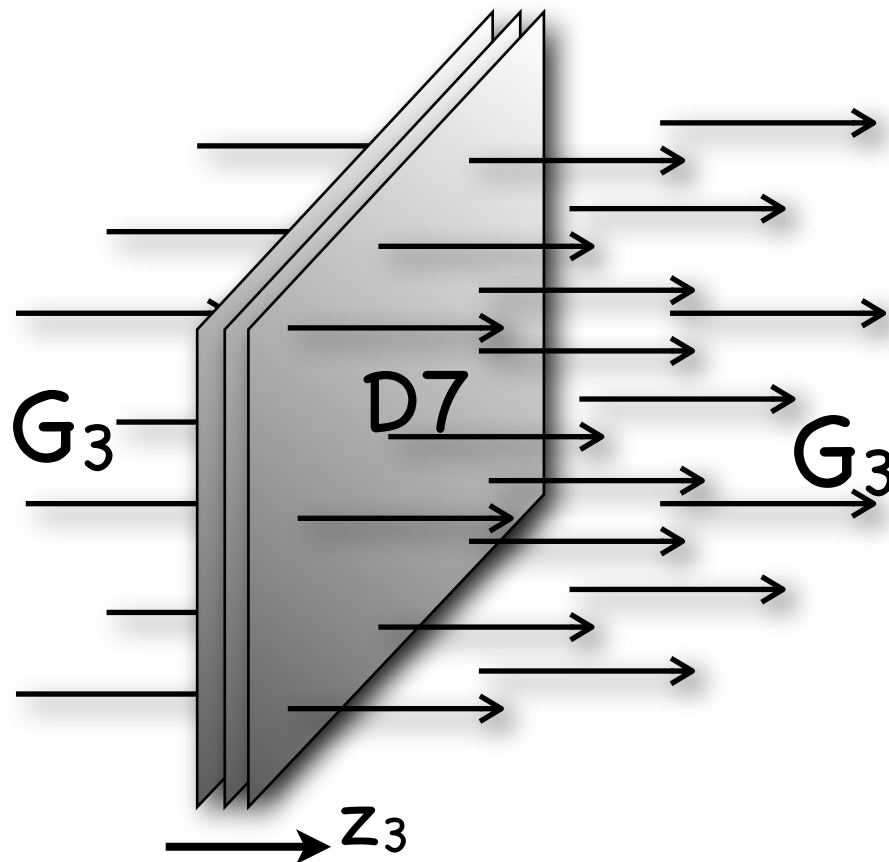
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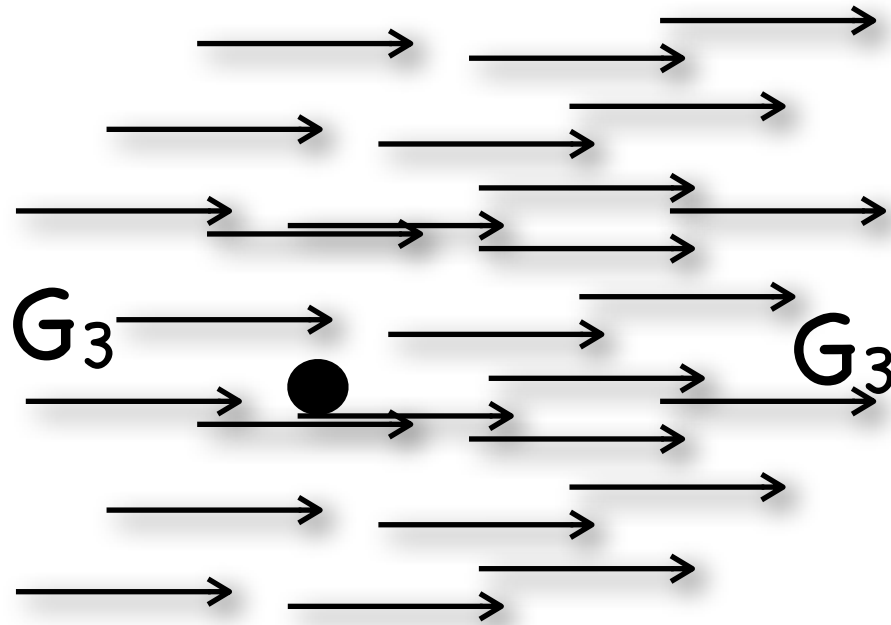
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Explicitly computable using D-brane world-volume action in general supergravity background, or using 4d effective theory approach



Fluxes, susy breaking and soft terms

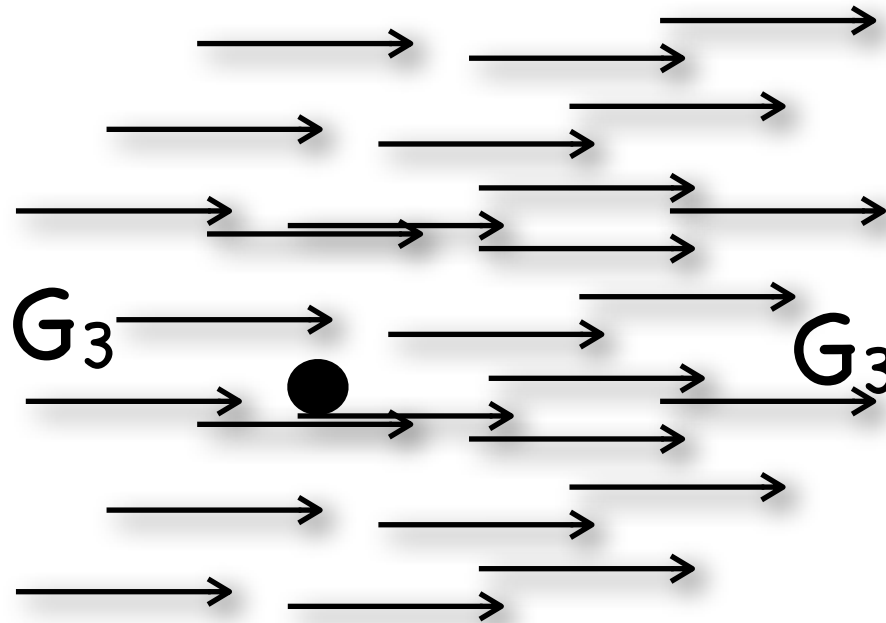
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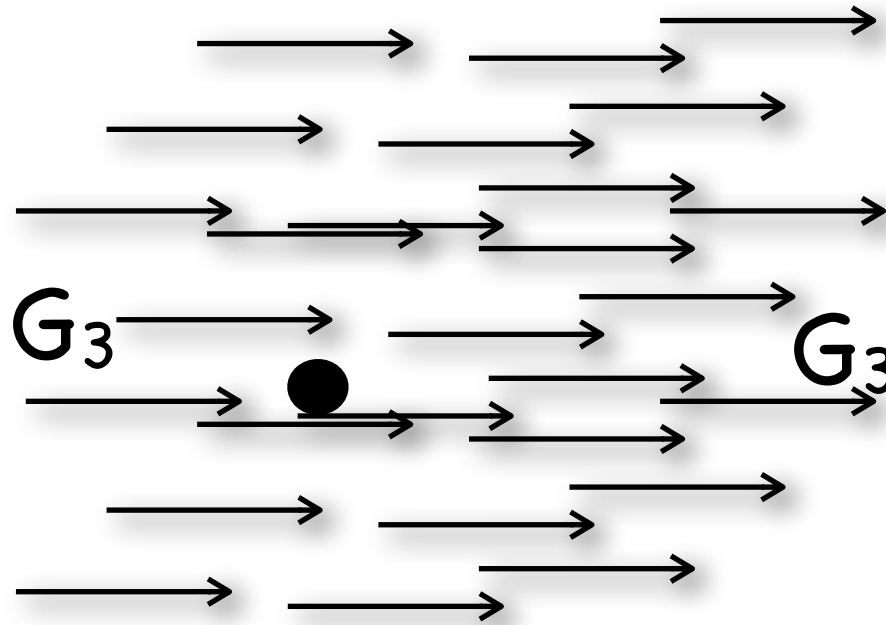
ISD flux background: same tension and 4-form charge as a D3-brane
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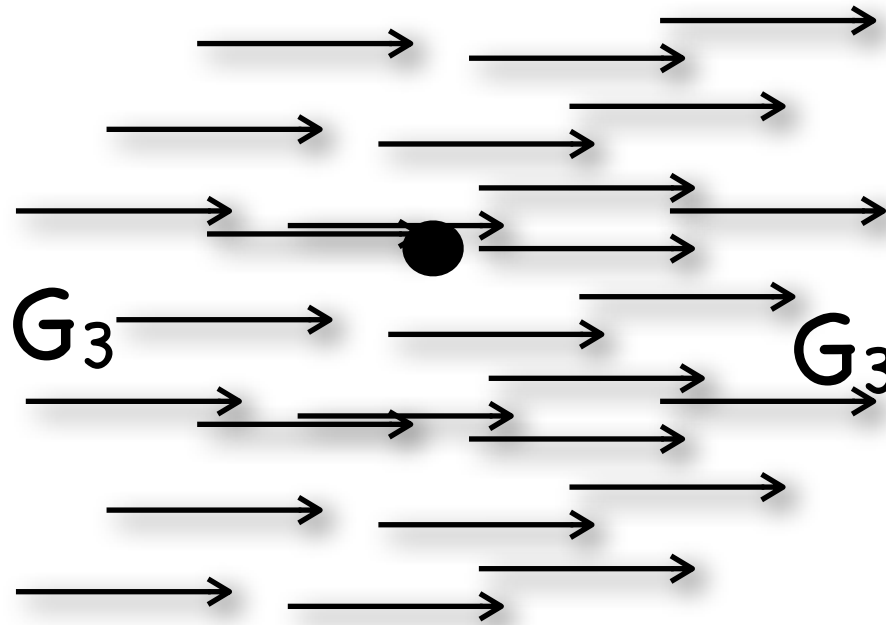
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$$m^2 \sim |G_{(0,3)}|^2$$

$$M \sim G_{(0,3)}$$

$$A \sim G_{(0,3)}$$

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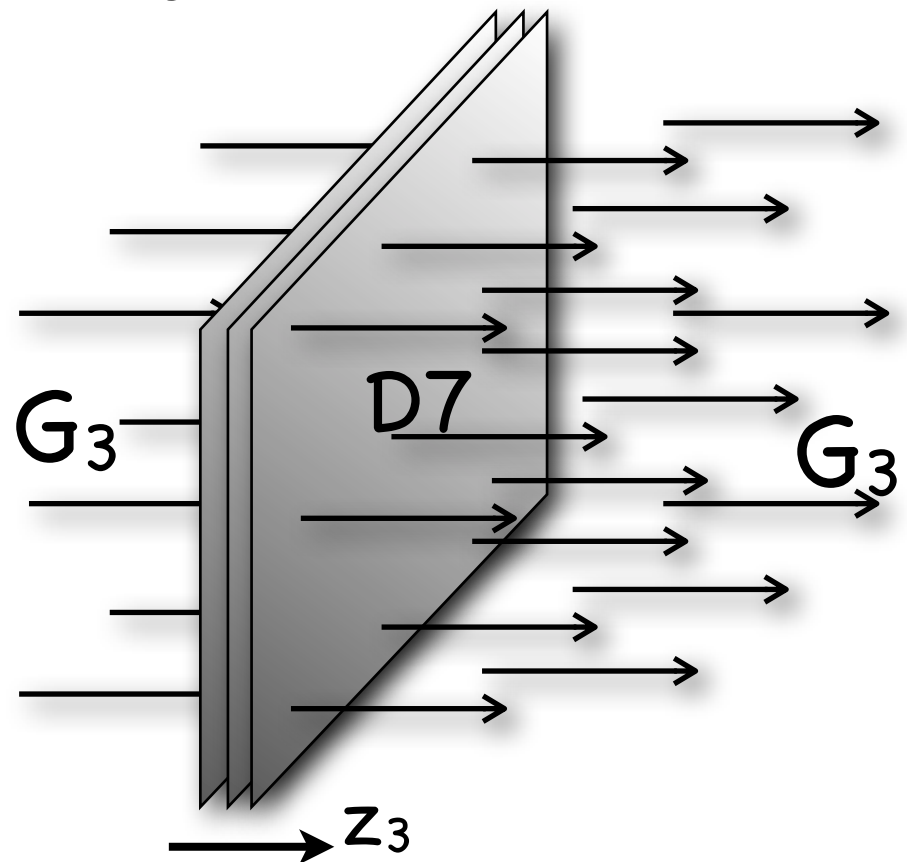
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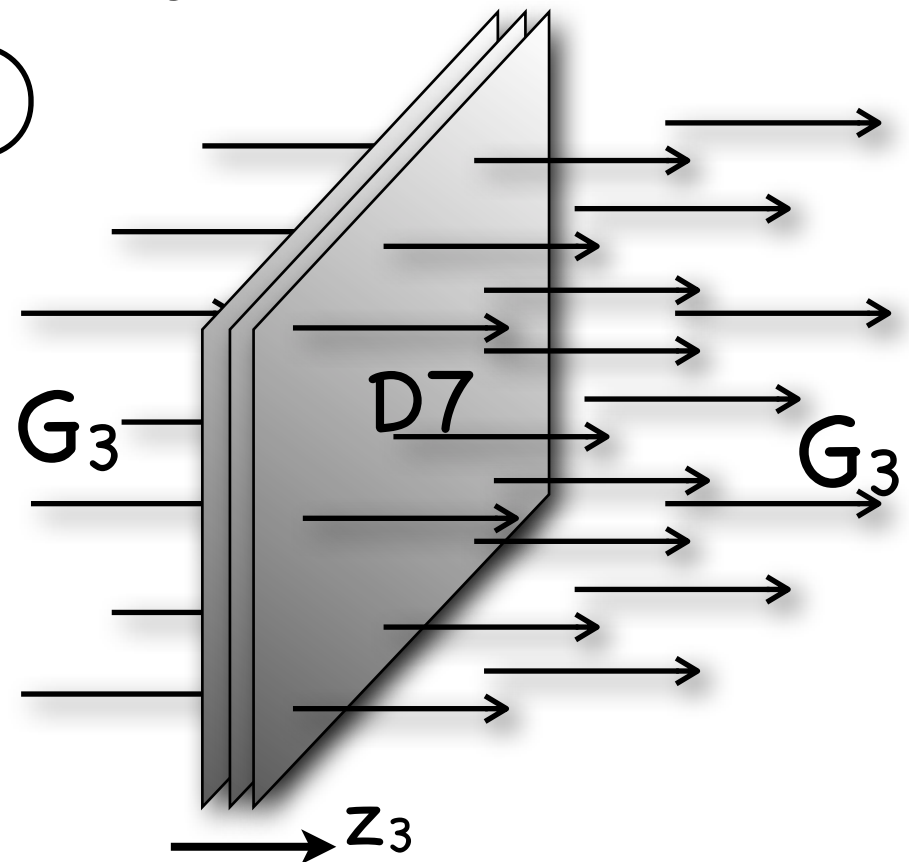


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Can recover both D3 and D7 results from effective theory

Flux components are vevs for moduli auxiliary fields: spurions

$$D_\tau W \sim \int_X \bar{G}_3 \wedge \Omega \quad D_{z_i} W \sim \int_X G_3 \wedge \chi_{(2,1),i} \quad D_\rho W \sim \int_X G_3 \wedge \Omega$$

Fluxes, susy breaking and soft terms


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Gravity mediation (in general, not universal, no mSUGRA)

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 Scales $M_{SUSY} \sim f \frac{\alpha'}{R^3} \sim f \frac{M_c^2}{M_p}$ f: possible local suppression

LowE Susy: TeV soft terms from $M_c \sim 10^{11} \text{GeV}$

~~LowE Susy:~~ Choose $M_c \sim 10^{14} \text{GeV}$ then $M_{SUSY} \sim 10^{10} \text{GeV}$

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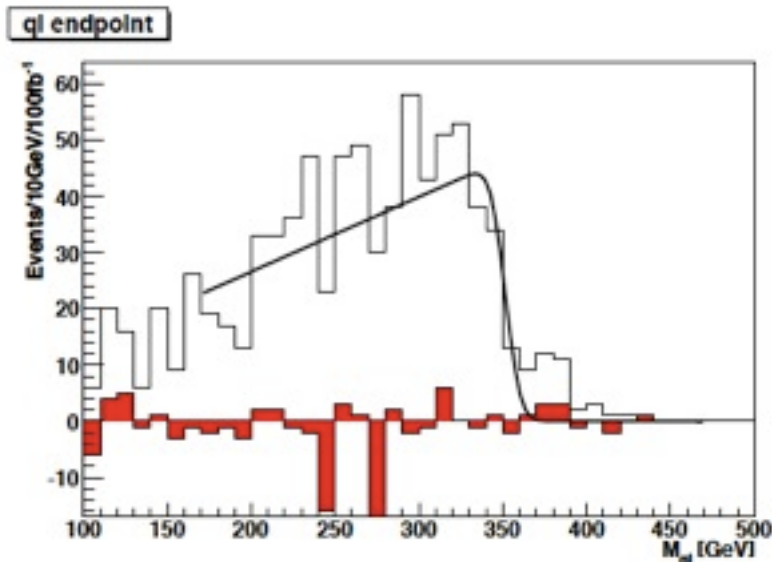
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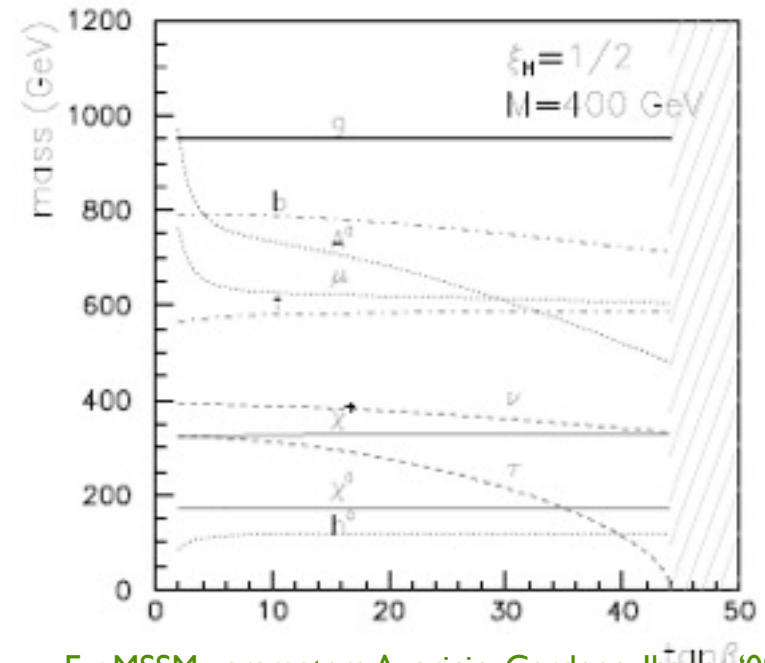
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📌 Can get to make plots



Ex: spectrum reconstruction from edges
Conlon, Kom, Suruliz, Allanach, Quevedo '07

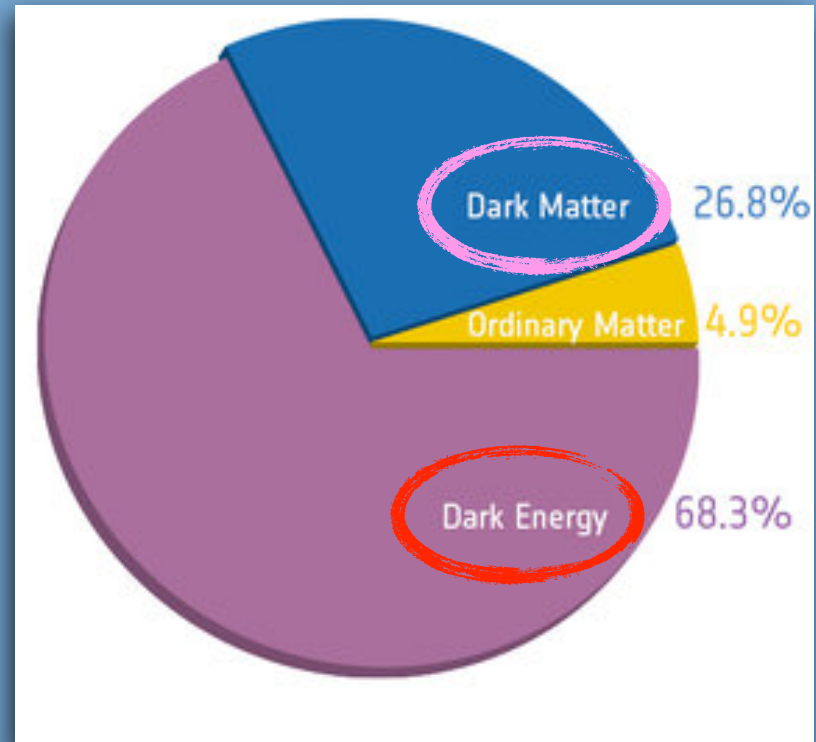
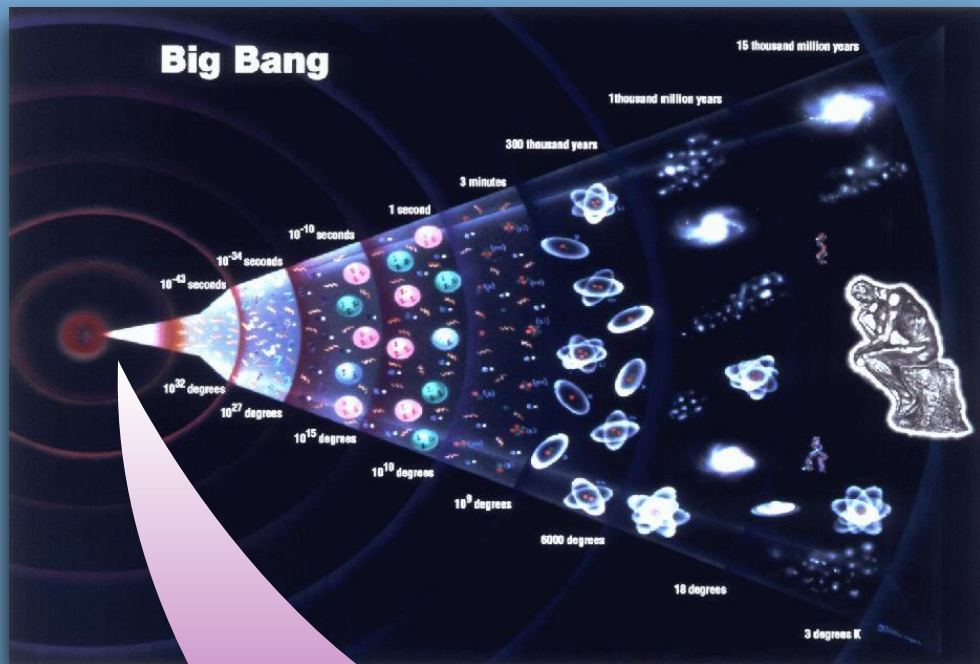


Ex: MSSM parameters Aparicio, Cerdeno, Ibanez '08

Cosmological Standard Model

(Λ CDM, "concordance model")

Flash
back





inflation
cf. Shiu's
lectures

Fluxes, 3-forms and axion monodromy



Recent interest in large field inflation

Fluxes, 3-forms and axion monodromy

-  Recent interest in large field inflation
-  Scalars with shift symmetry (axions) are well protected
continuous symmetry broken by non-pert effects to a discrete periodicity

Fluxes, 3-forms and axion monodromy

- Recent interest in large field inflation
- Scalars with shift symmetry (axions) are well protected
continuous symmetry broken by non-pert effects to a discrete periodicity
- String theory axions have sub-Planckian decay constant

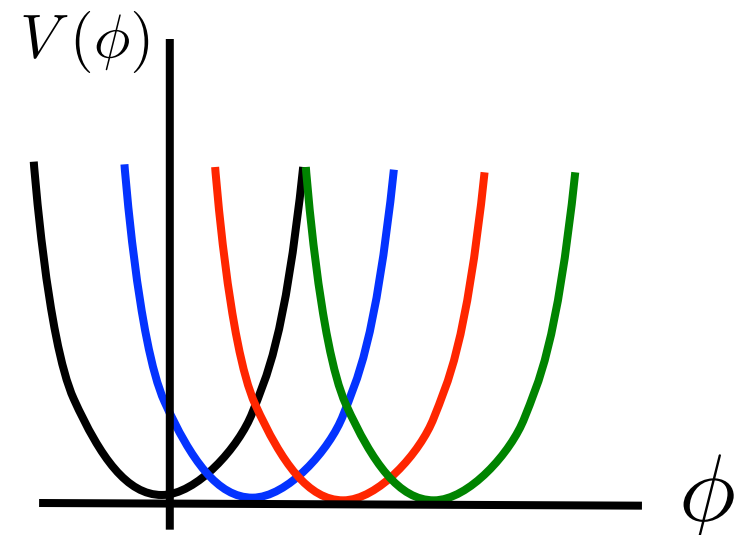
Axion monodromy:

Potential is periodic but multivalued

Field theory analogue:

theta dependent vacuum energy

in large N pure gluodynamics **Witten**



Discrete Z_n gauge symmetries



Flash
back

Discrete Z_n gauge symmetries



Discrete gauge symmetries

Flash
back

Discrete \mathbb{Z}_n gauge symmetries



Discrete gauge symmetries

\mathbb{Z}_n particles, \mathbb{Z}_n strings, ...

Flash
back

Discrete \mathbb{Z}_n gauge symmetries

Flash
back

📌 Discrete gauge symmetries

\mathbb{Z}_n particles, \mathbb{Z}_n strings, ...

📌 There are \mathbb{Z}_n gauge symmetries associated to 4d domain walls

Discrete \mathbb{Z}_n gauge symmetries

Flash
back

📌 Discrete gauge symmetries

\mathbb{Z}_n particles, \mathbb{Z}_n strings, ...

📌 There are \mathbb{Z}_n gauge symmetries associated to 4d domain walls

\mathbb{Z}_n symmetry of a 3-form that eats up a 2-form

$$\frac{1}{2}|F_4|^2 + |db_2 - n c_3|^2$$

Gauge invariance

$$c_3 \rightarrow c_3 + d\Lambda_2 \quad ; \quad b_2 \rightarrow b_2 + n\Lambda_2$$

Discrete \mathbb{Z}_n gauge symmetries

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📌 Dualizing b_2 to an axion, get Kaloper-Sorbo description of axion monodromy models.

Discrete \mathbb{Z}_n gauge symmetries

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- Can consider other \mathbb{Z}_n charged objects in 4d
- Lagrangian for 3-form eating up a 2-form

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Massive axion

Can arise in D-brane & flux models

Fluxes, 3-forms and axion monodromy

 Structure is automatic in flux compactifications

After all, fluxes produce the stabilization of axions in moduli!!

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Integrating over fluxed CY with $\phi = \int_{\Sigma_2} B_2$, $M = \int_{\Pi_p} F_p$

Change in axion induces extra flux

$$\Delta\phi \rightarrow \Delta \int_{\Sigma_2 \times \Pi_p} \tilde{F}_{p+2} = \phi M$$

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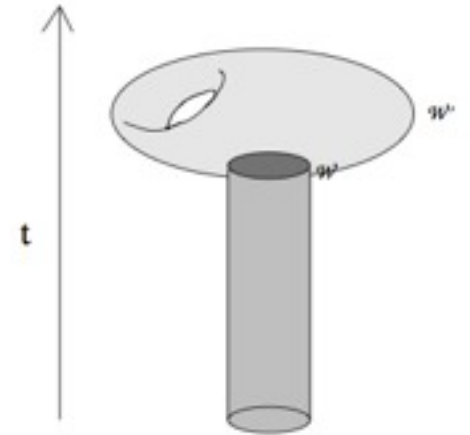
Monodromy

Multiple branches connected by domain walls changing (p+2)-form flux. They are D(6-p) on (4-p)-cycle

Fluxes, 3-forms and axion monodromy



Many other realizations



Fluxes, 3-forms and axion monodromy

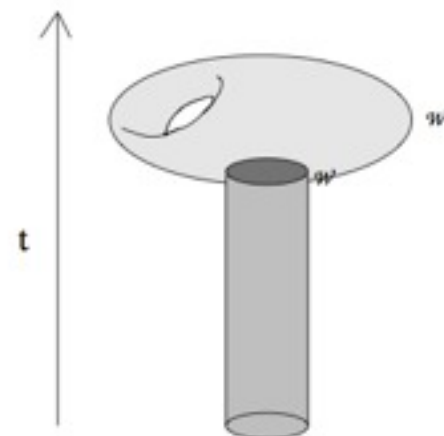


Many other realizations

Ex: IIB with NSNS flux on A-cycle

$$F_4 = \int_B F_7 \quad ; \quad \int_A H_3 = n$$

$$\int_{4d \times X_6} \overline{H}_3 C_0 F_7 \rightarrow \int_{4d} n C_0 F_4$$



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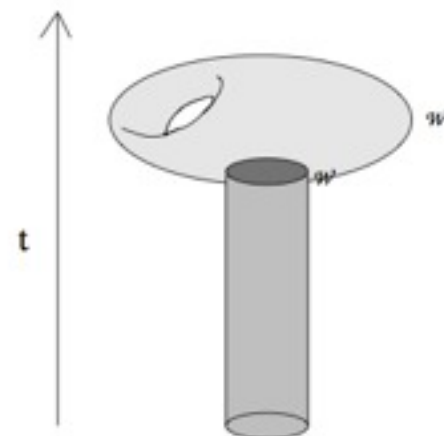
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10d IIB axion has a monodromy. Origin of energetic cost?

GVW sup_o $W = \int_{X_6} (F_3 - \tau \overline{H}_3) \wedge \Omega$



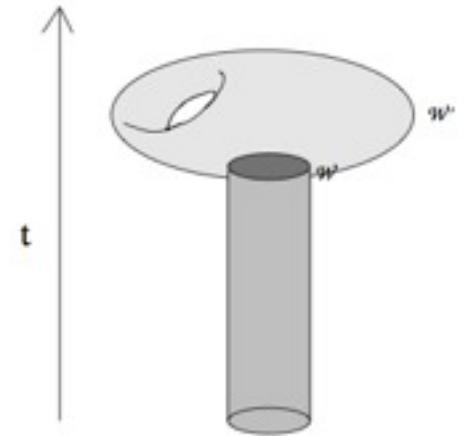
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Period of C_0 changes n units of F_3 flux on A

4d Domain Wall is $D5$ on B

4d instanton is $D(-1)$ (cosine modulation)

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- Neutrino masses, non-perturbative operators
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 Expect continuous progress and new results and useful input from LHC & cosmo

To all organizers & hosts & participats of
The 10th Asian Winter School on Strings, Particles and Cosmology



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ありがとう

Arigato * Thank You