

Feeling fivebranes in matrix models

by Monte Carlo Simulation

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Collaboration with
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1. Introduction

M(atrrix) Theory (BFSS model)

[de Wit–Hoppe–Nicolai '88,
Banks–Fischler–Shenker–Susskind '96]

Defined by

- D0–brane quantum mechanics
- Matrix regularization of super–membrane theory

$$X_M(\tau, \sigma_1, \sigma_2) \rightarrow \hat{X}_M(\tau)$$

$N_c \times N_c$ matrix

Non–perturbative formulation of M–theory

- Correct supergraviton interactions
- M2 is realized as a classical configuration.

1. Introduction

Gauge/gravity duality for BFSS

[Itzhaki-Maldacena-Sonnenschein-Yankielowicz '98]

Thermal
BFSS
(D0 QM)



IIA SUGRA
around black 0-branes

$$\frac{\lambda}{U^3 N_c^{\frac{4}{7}}} \sim \langle e^\phi \rangle^{\frac{4}{7}}$$
$$\frac{U^3}{\lambda} \sim \left(\frac{\alpha'}{R^2} \right)^2$$

U : energy scale

R : curvature radius

$$\lambda = g_{YM}^2 N_c$$

Hawking temperature: $T \sim \frac{1}{L} \left(\frac{u_0}{L} \right)^{\frac{5}{2}}$

L : length scale

1. Introduction

Gauge/gravity duality for BFSS

Success in computational simulation:

Hanada–Hyakutake–Ishiki–Nishimura '14

Correct T and g_s expansion with α' correction from BFSS

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When we go back to M–theory description,

How does **M5** play a role in Matrix model?

cf) transverse M5

... M5 charges in BMN model

[Maldacena–Sheikh Jabbari–Raamsdonk '02]

longitudinal M5 ... **BD model** [Berkooz–Douglas '96]
($\rightarrow$ **D4** in IIA)

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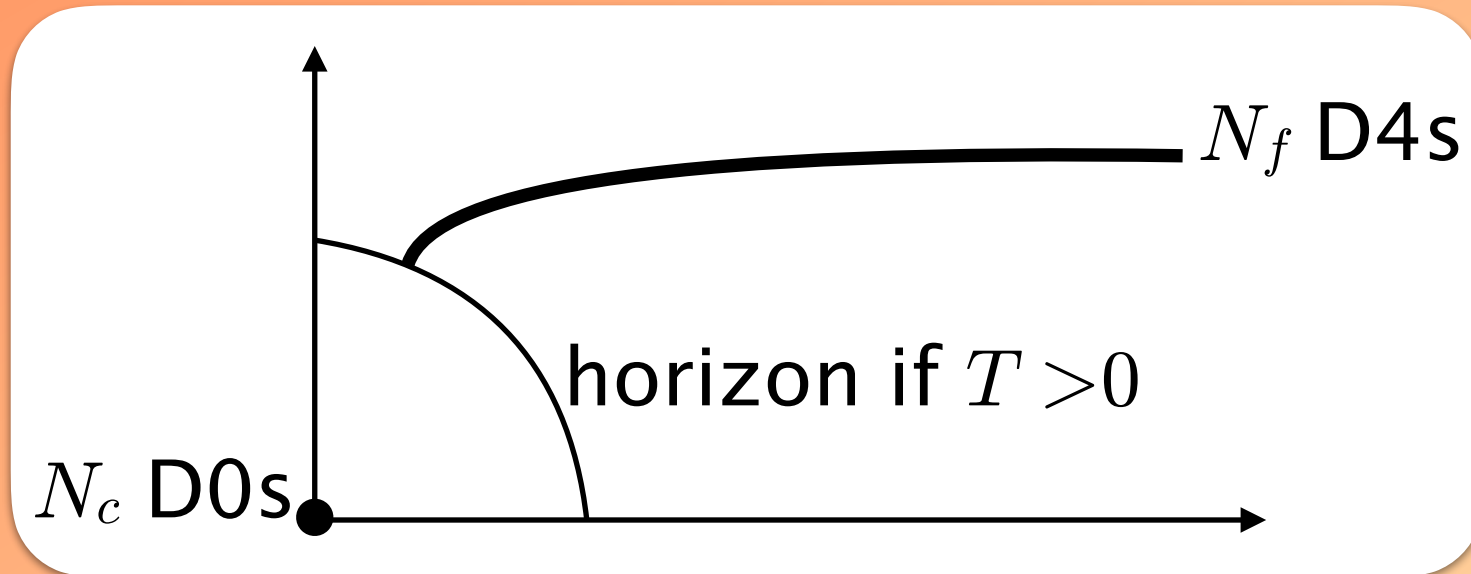
Let's compute **BD model**
by Monte Carlo simulation.

2. BD model

D0 quantum mechanics with probe D4

[Berkooz–Douglas '96]

$SO(9)$ symmetry breaks down to $SO(5) \times SO(4)$.



quench approx.
 $N_f \ll N_c$

field contents:

D0–D0 strings

X^M : $N_c \times N_c$ Hermitian

Ψ : 2 $N_c \times N_c$ symplectic Majoranas

D0–D4 strings

Φ_ρ : $N_c \times N_f$ complex

χ : $N_c \times N_f$ Weyl

(in the meaning of $SO(5)$ spinor)

$M=1, \dots, \overbrace{5}^{SO(5)}, \overbrace{6, \dots, 9}^{SO(4)}$

$\rho, \dot{\rho} = 1, 2 : SU(2) \times SU(2) \sim SO(4) \text{ sym.}$

3. Monte Carlo simulation of BD model

[Filev–O'Connor '15]

Gauge/gravity agreement in Mass–Condensate plot

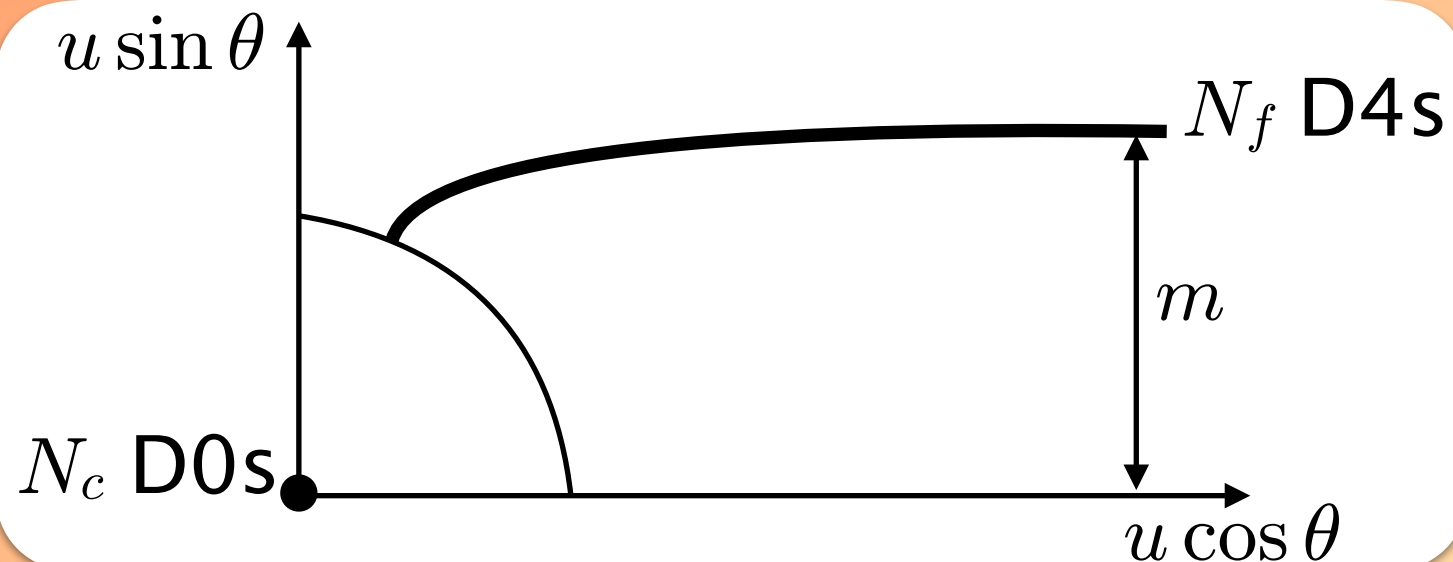
Gauge theory side

$$c = \frac{\partial F}{\partial m} \quad (F = -\beta^{-1} \ln Z)$$

Gravity side

classical sol. of DBI eom:

$$\sin[\theta(u)] = \tilde{m} \frac{u_0}{u} + \tilde{c} \frac{u_0^3}{u^3} + \dots$$



3. Monte Carlo simulation of BD model

[Filev-O'Connor '15]

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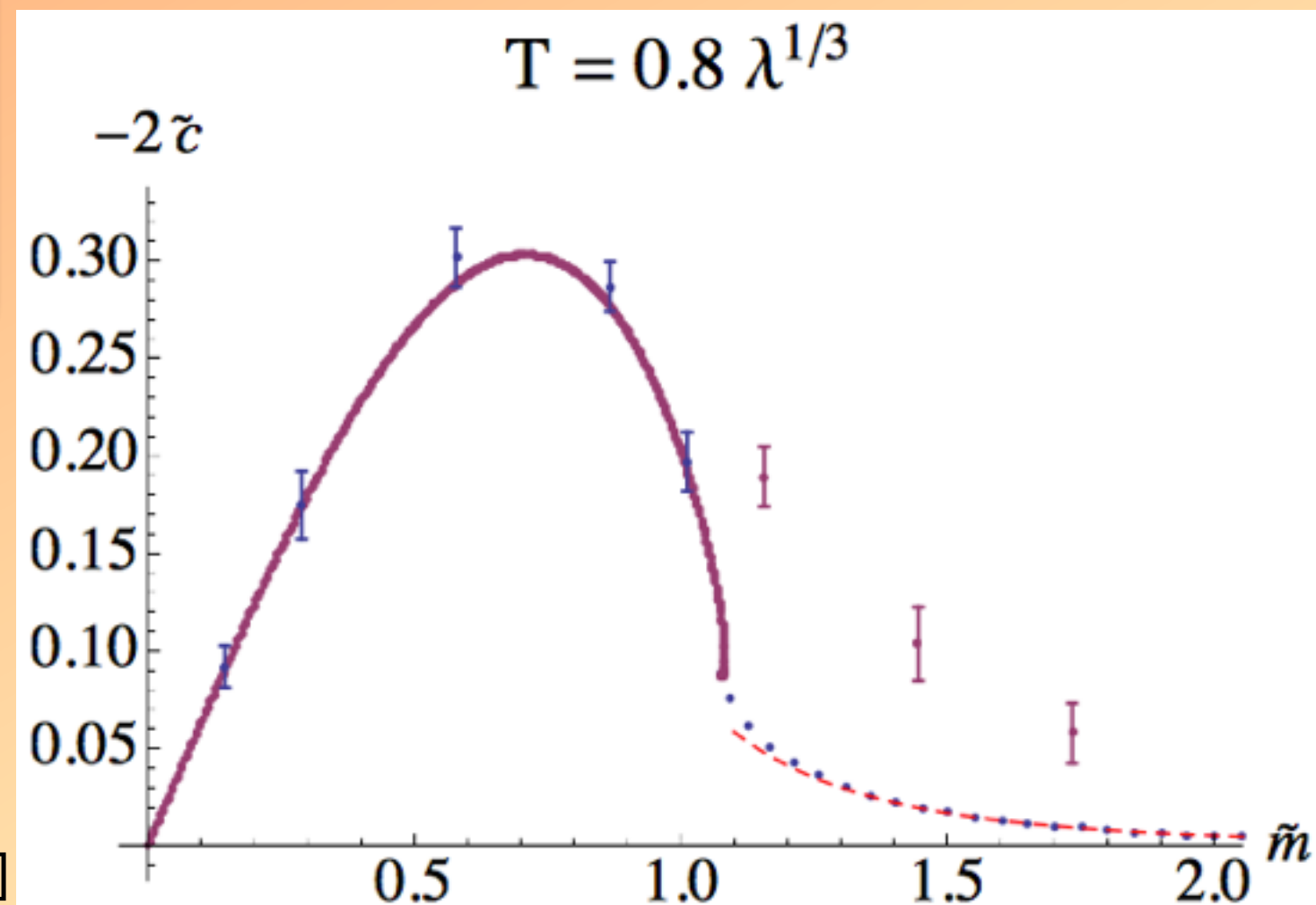
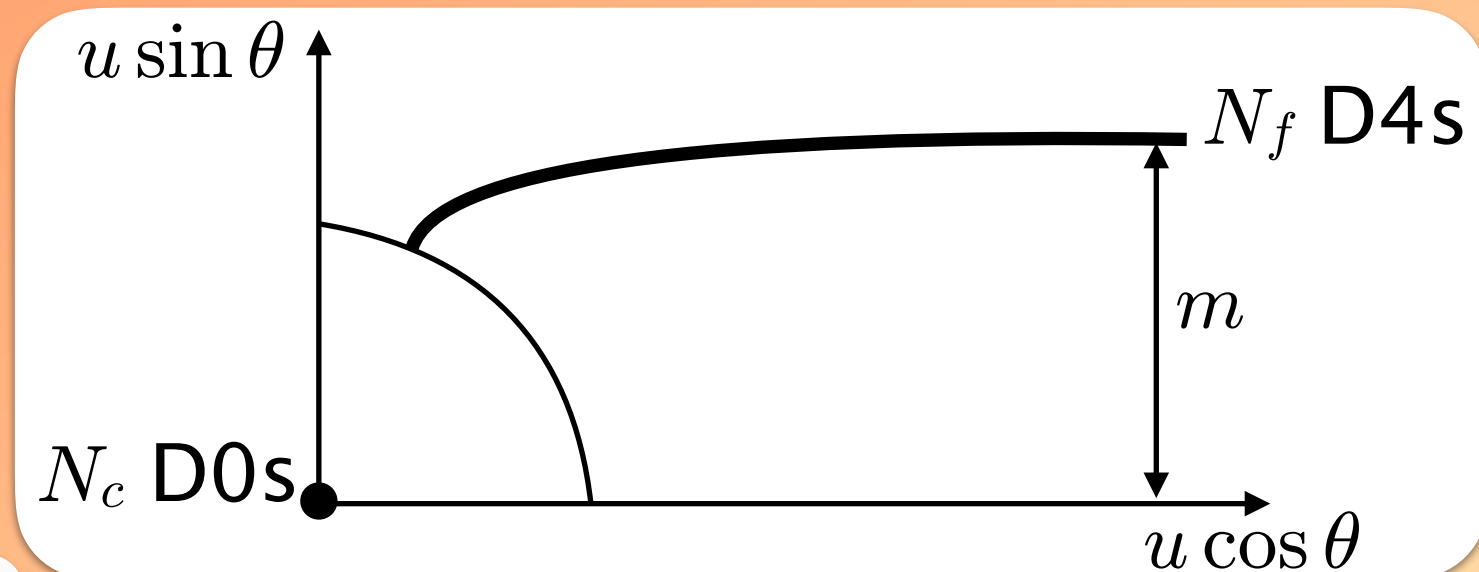
$$\sin[\theta(u)] = \tilde{m} \frac{u_0}{u} + \tilde{c} \frac{u_0^3}{u^3} + \dots$$

Agreement at small mass!
(deconfined phase)

solid red line: gravity side

points w/ error bars: BD model

[Filev-O'Connor '15]



4. More BD model simulation

– High temperature expansion

[YA–Filev–Kovacik–O’Connor to appear]

- Able to analyze BD model by 0–mode dynamics
- Further check in high temperature regime

– Deformed BD model (KLY model) [Kim–Lee–Yi ’02]

- M–theory with M5 in pp–wave background
(analog of BMN model)
→ # of SUSY in KLY model is same as BD model.
- Difficulty in BFSS and BD model would be gone.
(due to flat–direction)
→ smaller N_c can be investigated more easily

5. Summary

- Monte Carlo simulation of **BD model** (D0 QM with probe D4)
- **Agreement between gauge & gravity** sides for condensate c
- We are working on BD model and the deformed one to get more (nontrivial) gauge/gravity evidence.

Backups

1. Introduction

Gauge/gravity duality for BFSS

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$$\langle e^\phi \rangle^{\frac{4}{7}} \sim \frac{\lambda}{U^3 N^{\frac{4}{7}}}$$
$$\left(\frac{\alpha'}{R^2} \right)^2 \sim \frac{U^3}{\lambda},$$

U : energy scale
 R : curvature radius
 $\lambda = g_{YM}^2 N$

near-horizon black 0-brane solution:

$$ds^2 = - \left(\frac{u}{L} \right)^{\frac{7}{2}} \left(1 - \left(\frac{u_0}{u} \right)^7 \right) dt^2 + \left(\frac{L}{u} \right)^{\frac{7}{2}} \left\{ \frac{du^2}{1 - \left(\frac{u_0}{u} \right)^7} + u^2 d\Omega_8^2 \right\}$$

Hawking temperature: $T \sim \frac{1}{L} \left(\frac{u_0}{L} \right)^{\frac{5}{2}}$

L : length scale

1. Introduction

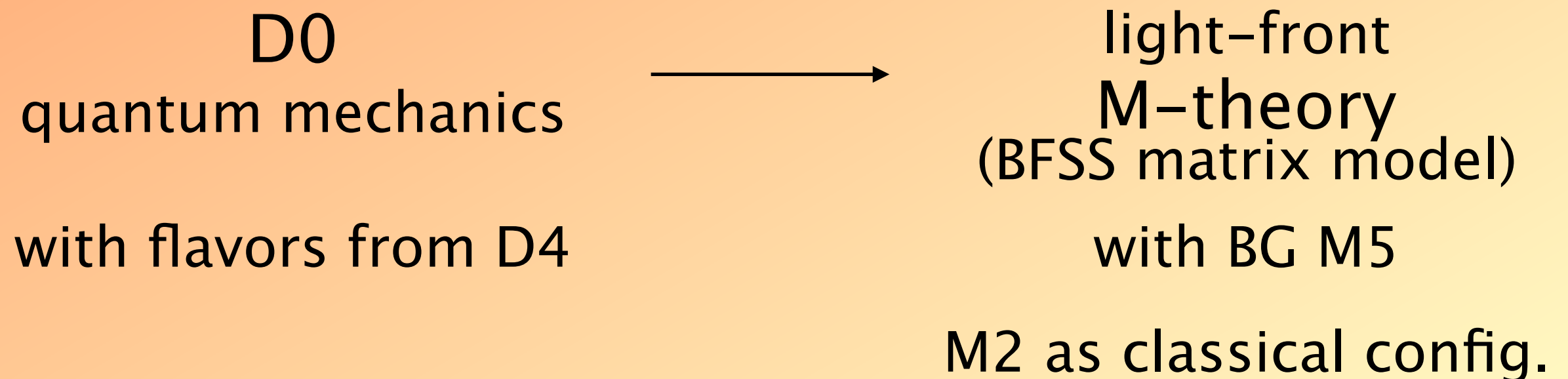
My motivation:

Matrix model as
a non-perturbative formulation of string/M theory

How about M5 in BFSS?

longitudinal M5

transverse M5

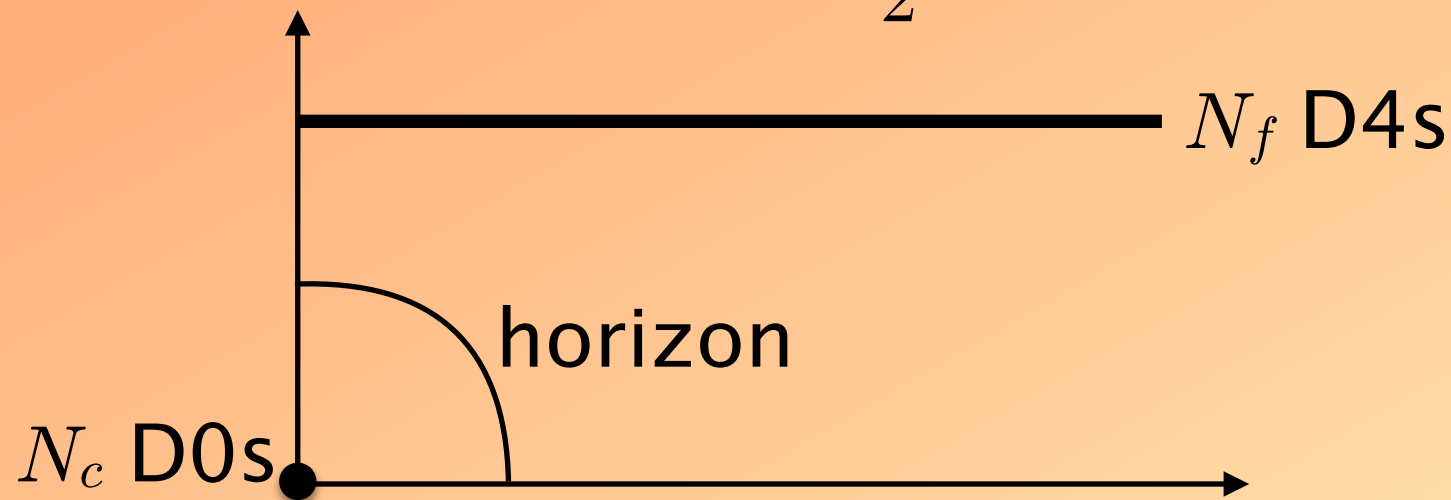


2. BD model

D0 quantum mechanics in background D4

$SO(9)$ symmetry breaks down to $SO(5) \times SO(4)$.

$$S = N_c \int_0^\beta d\tau \left[\text{Tr} \left(\frac{1}{2} D_\tau X^M D_\tau X^M + \frac{1}{4} [X^M, X^N]^2 \right) + \text{tr} \left(D_\tau \bar{\Phi}^\rho D_\tau \Phi_\rho + \bar{\Phi}^\rho (X^a - m^a)^2 \Phi_\rho - \bar{\Phi}^\rho [\bar{X}^{\sigma\dot{\rho}}, X_{\rho\dot{\rho}}] \Phi_\sigma - \frac{1}{2} \bar{\Phi}^\rho \Phi_\sigma \bar{\Phi}^\sigma \Phi_\rho + \bar{\Phi}^\rho \Phi_\rho \bar{\Phi}^\sigma \Phi_\sigma \right) + \text{fermions} \right]$$



$$N_f \ll N_c$$

X^M : $N_c \times N_c$ matrix field

Φ_ρ : $N_c \times N_f$ matrix field

$M=1, \dots, 9$

$a=1, \dots, 5$: $SO(5)$ sym.

$\rho, \dot{\rho} = 1, 2$: $SU(2) \times SU(2) \sim SO(4)$ sym.

2. BD model

D0 quantum mechanics in background D4

$$S = N_c \int_0^\beta d\tau \left[\text{Tr} \left(\frac{1}{2} D_\tau X^a D_\tau X^a + \frac{1}{2} D_\tau \bar{X}^{\rho\dot{\rho}} D_\tau X_{\rho\dot{\rho}} + \frac{1}{2} \lambda^{\dagger\rho} D_\tau \lambda_\rho + \frac{1}{2} \theta^{\dagger\dot{\rho}} D_\tau \theta_{\dot{\rho}} \right) \right. \\ \left. + \text{tr} (D_\tau \bar{\Phi}^\rho D_\tau \Phi_\rho + \chi^\dagger D_\tau \chi) \right. \\ \left. + \text{Tr} \left(\frac{1}{4} [X^a, X^b]^2 + \frac{1}{2} [X^a, \bar{X}^{\rho\dot{\rho}}] [X^a, X_{\rho\dot{\rho}}] \right) \right. \\ \left. + \frac{1}{2} \text{Tr} \sum_{A=1}^3 \mathcal{D}^A \mathcal{D}^A + \text{tr} (\bar{\Phi}^\rho (X^a - m^a)^2 \Phi_\rho) \right. \\ \left. + \text{fermion coupling terms} \right]$$

$$\mathcal{D}^A = \sigma_\rho^A{}^\sigma \left(\frac{1}{2} [\bar{X}^{\rho\dot{\rho}}, X_{\sigma\dot{\rho}}] - \Phi_\sigma \bar{\Phi}^\rho \right)$$