

# String Phenomenology

Angel M. Uranga  
IFT UAM-CSIC



Asian Winter School  
January 2016

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# Plan

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> Panorama A

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- > Non-perturbative aspects

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- > **Panorama A**
- > **Non-perturbative aspects**
- > **Panorama B**
- > **Fluxes and applications**

# > Panorama A



# “String Phenomenology”

- 📌 String theory describes gravitational and gauge interactions in a unified framework, consistent at the quantum level
- 📌 What about the **Standard Model** ? (or whatever extension)

# “String Phenomenology”

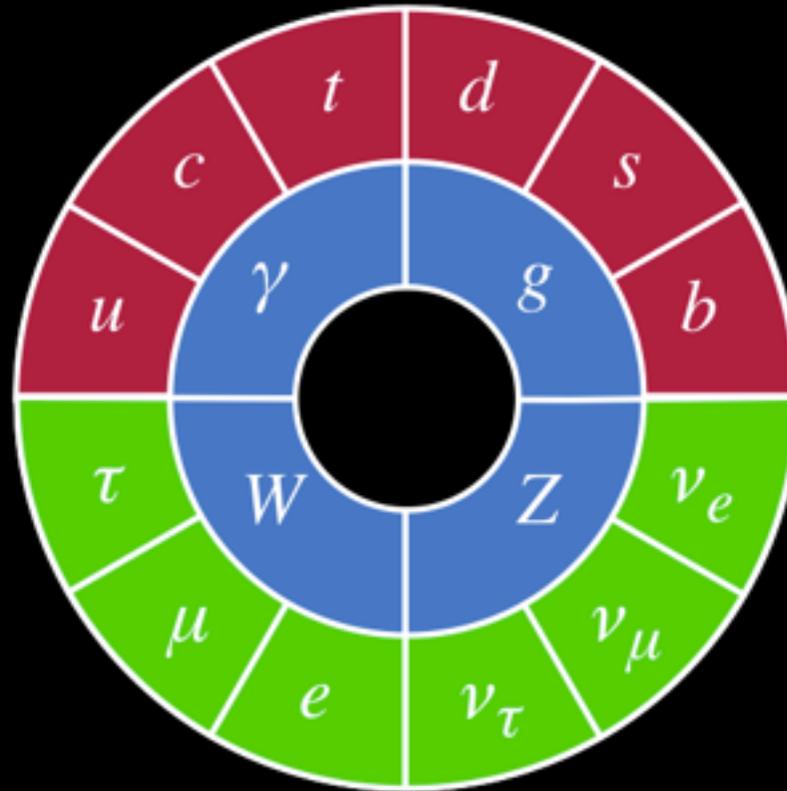
- 📌 String theory describes gravitational and gauge interactions in a unified framework, consistent at the quantum level
- 📌 What about the **Standard Model**? (or whatever extension)
- 📌 **Aim of String Phenomenology:**
  - Understand basic mechanisms realizing interesting properties  
Chirality, susy, gauge symm, moduli stab, flat potentials, instantons...
  - Determine classes of constructions incorporating subsets of them  
Non abelian gauge interactions, replicated charged fermions, Higgs scalars, Yukawas, moduli stabilization, tuning of cc, ...
  - Within each class, obtain explicit models as close to SM as possible with the hope of learning more about the high energy regime of SM in string theory

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- 📌 Old program, yet sustained progress, following formal developments...

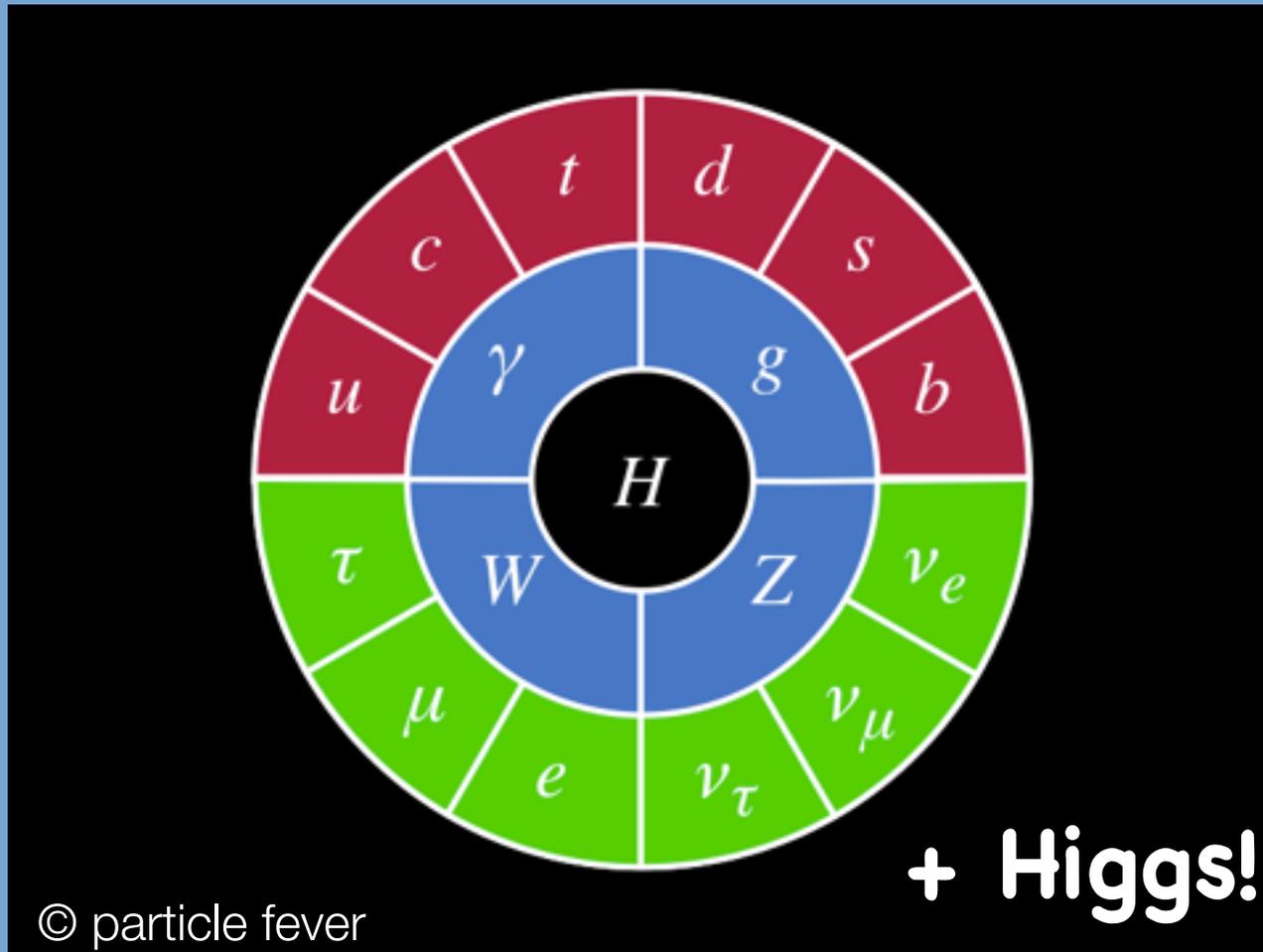
# Standard Model of Particles

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© particle fever

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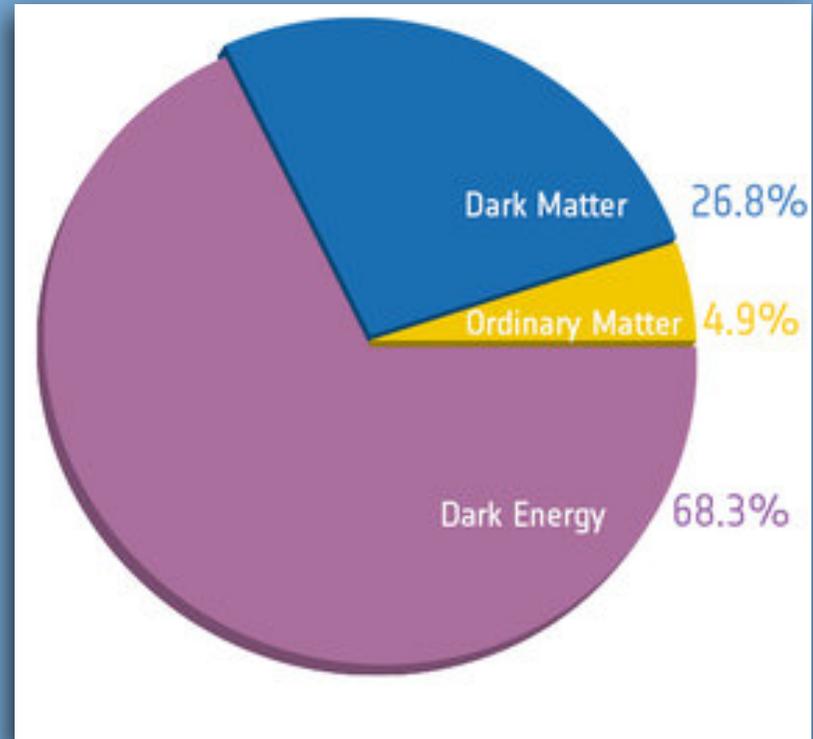
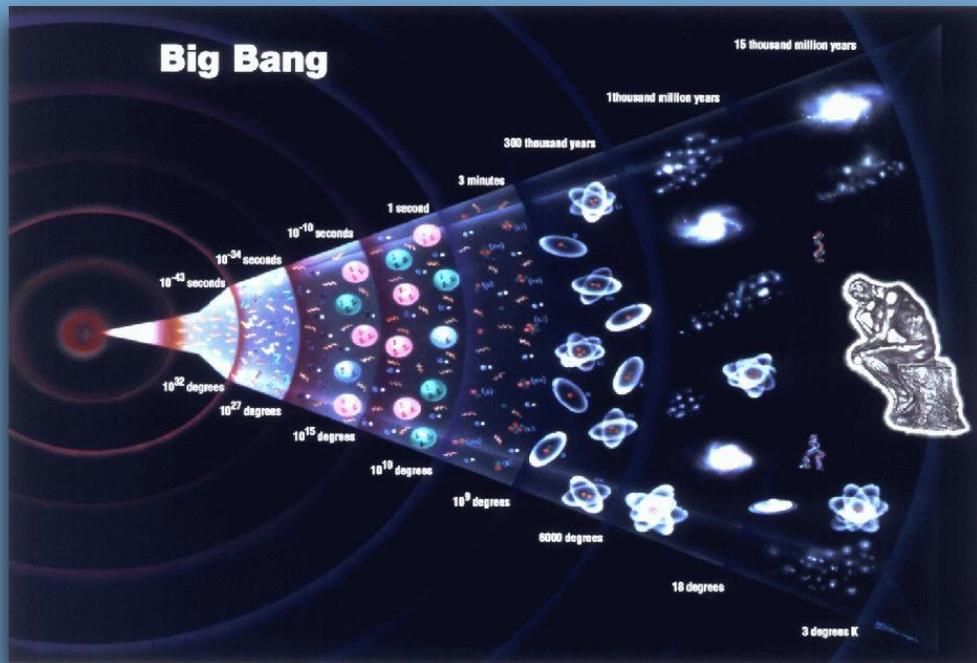
$$SU(3) \times SU(2) \times U(1)_Y$$

Matter fields		Y
$Q_L$	$3(3, 2)$	$1/6$
$U_R$	$3(\bar{3}, 1)$	$-2/3$
$D_R$	$3(\bar{3}, 1)$	$1/3$
$L$	$3(1, 2)$	$-1/2$
$E_R$	$3(1, 1)$	$1$
$\nu_R$	$3(1, 1)$	$0$

Higgs in  $(1, 2)$ ,  $Y = -1/2$

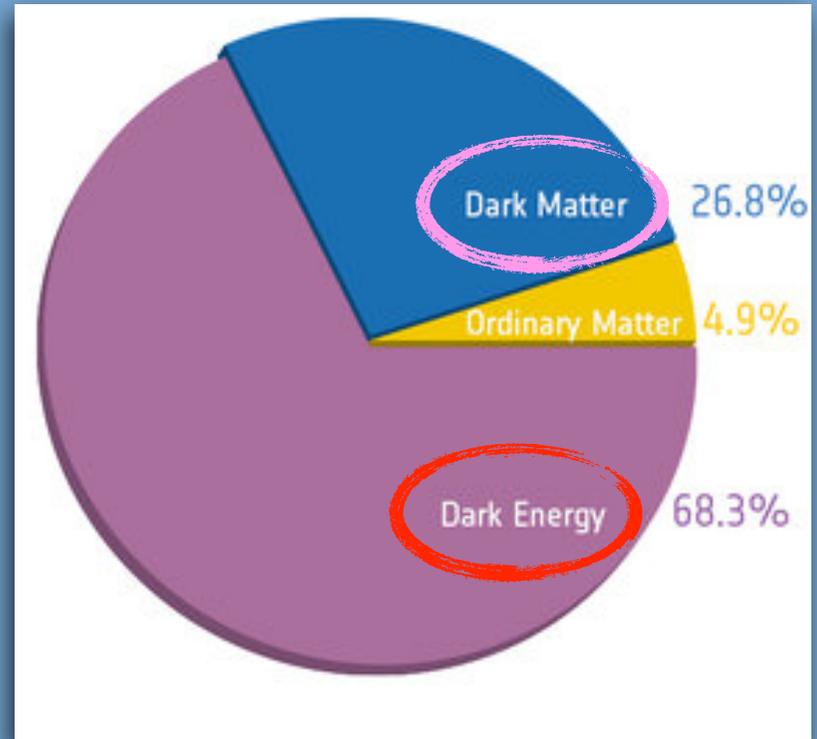
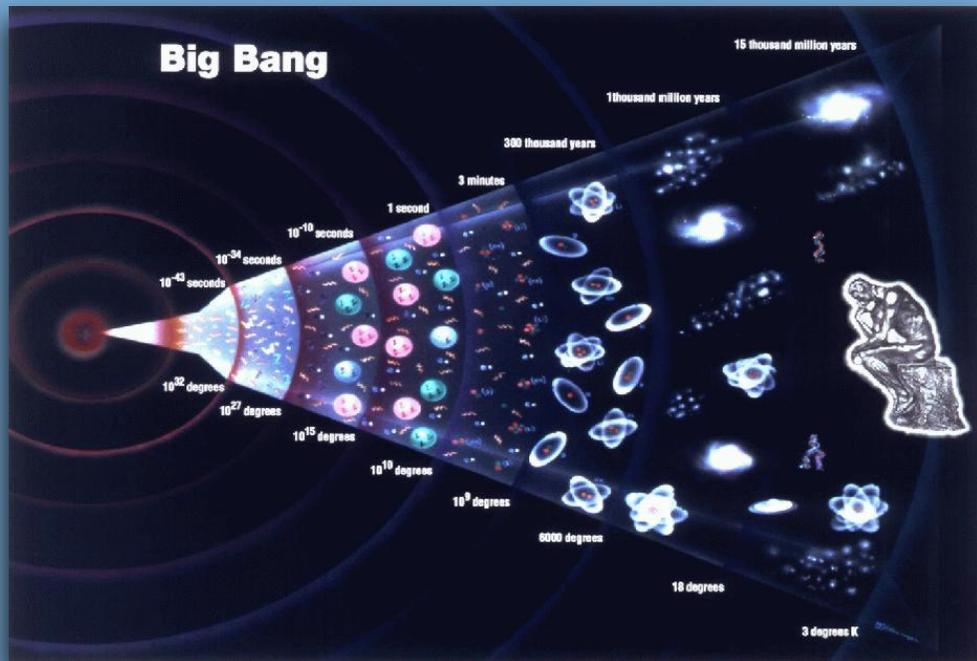
# Cosmological Standard Model

( $\Lambda$ CDM, "concordance model")



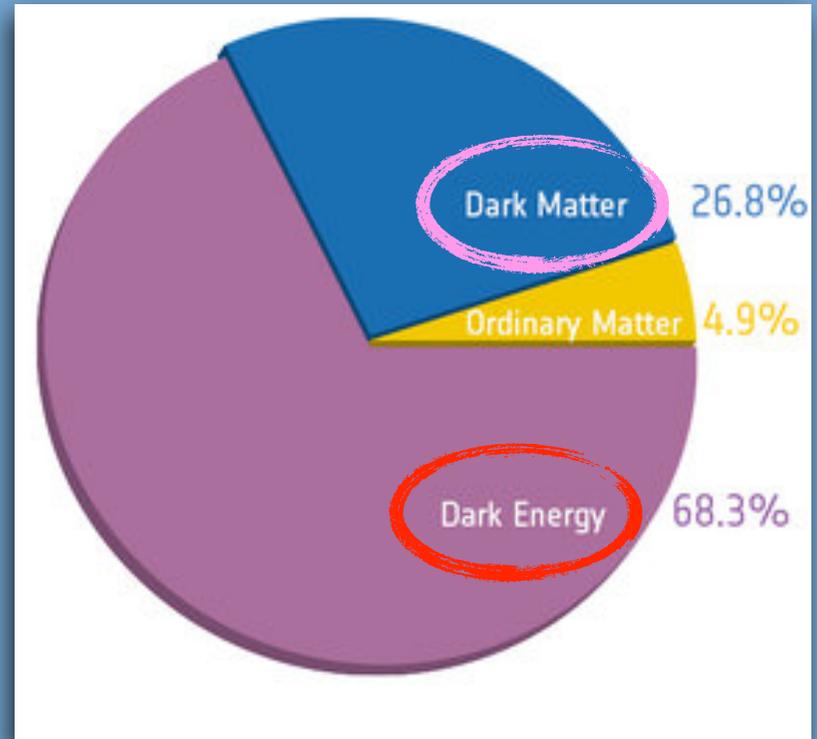
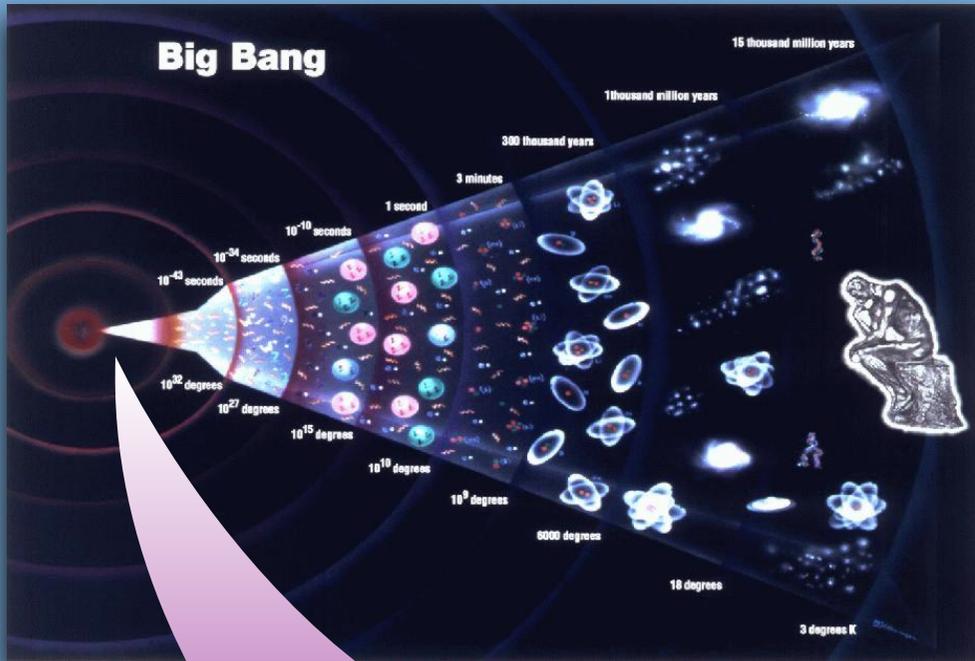
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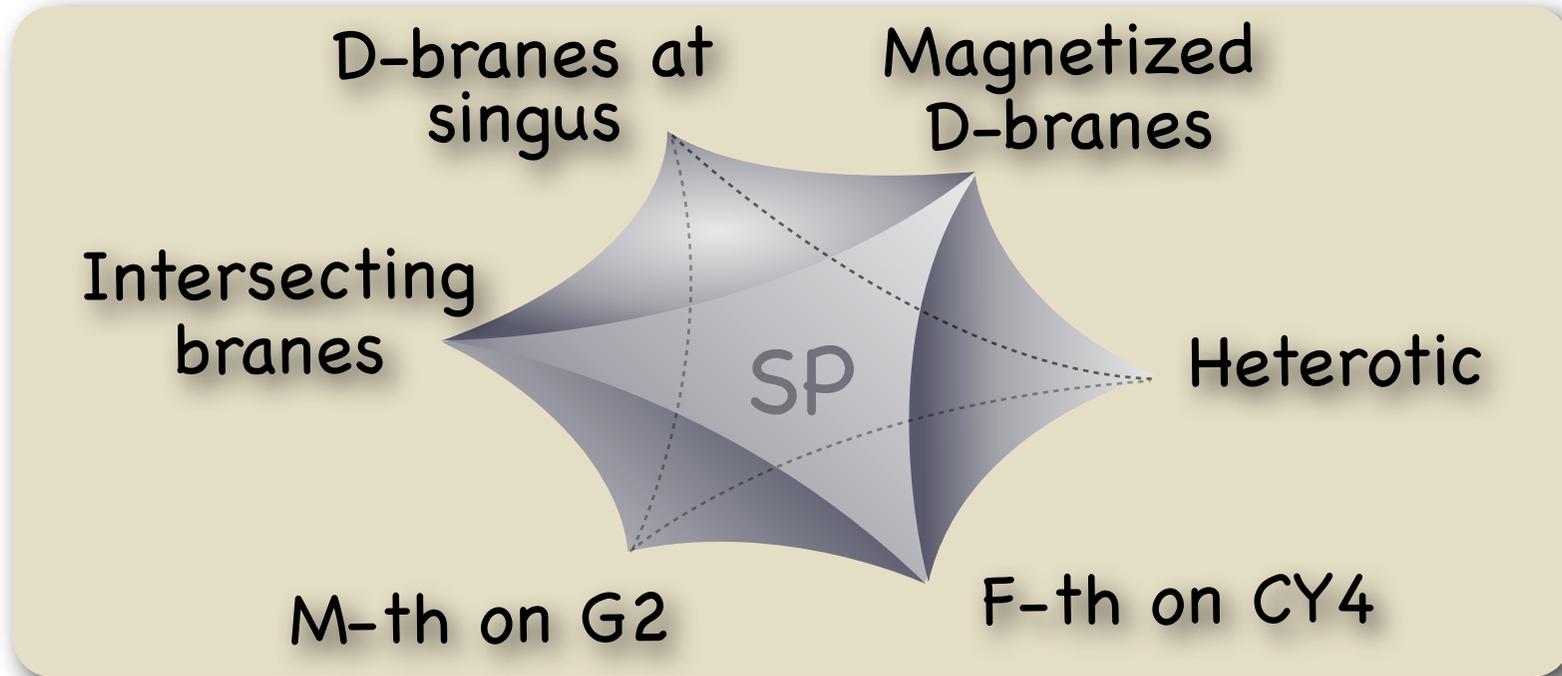
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inflation  
cf. Shiu's  
lectures

# "String Phenomenology"

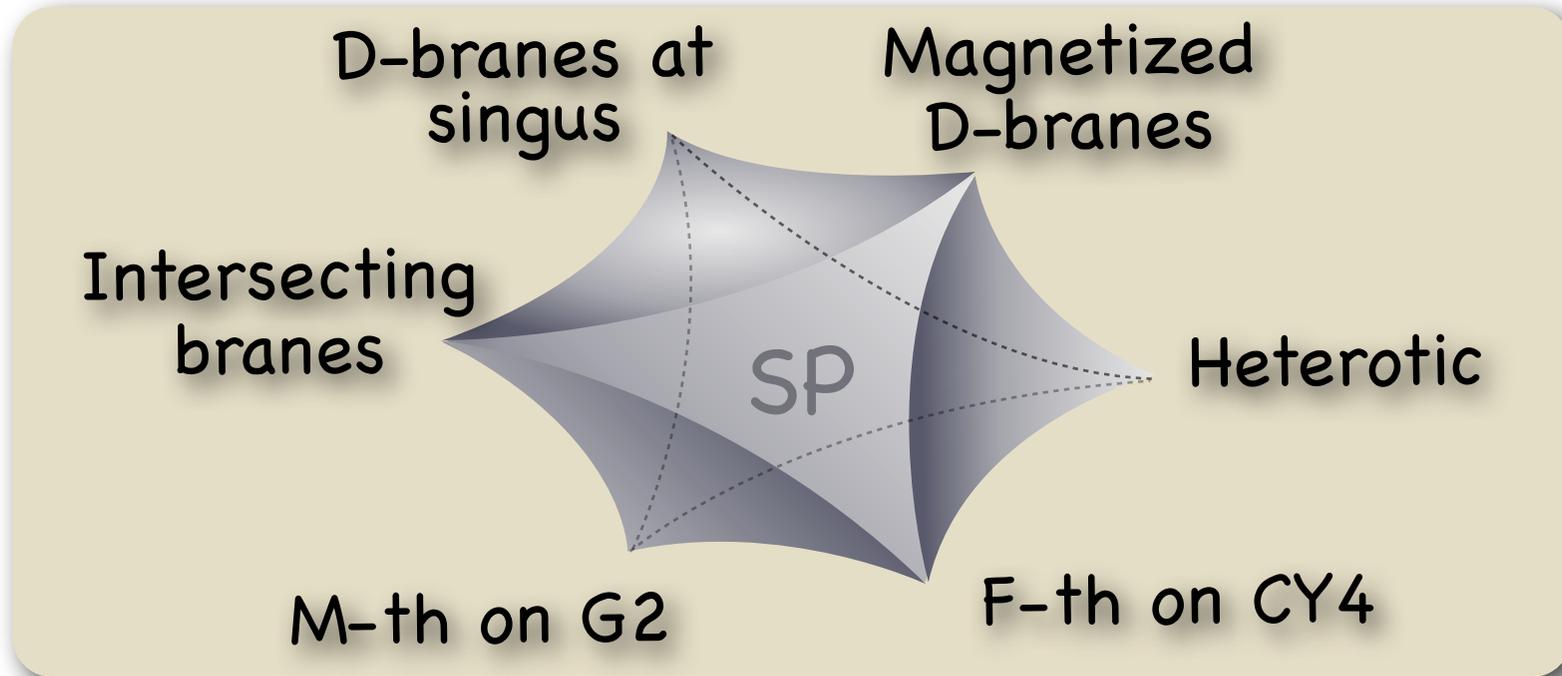
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Remarkably, generic 4d compactifications are chiral

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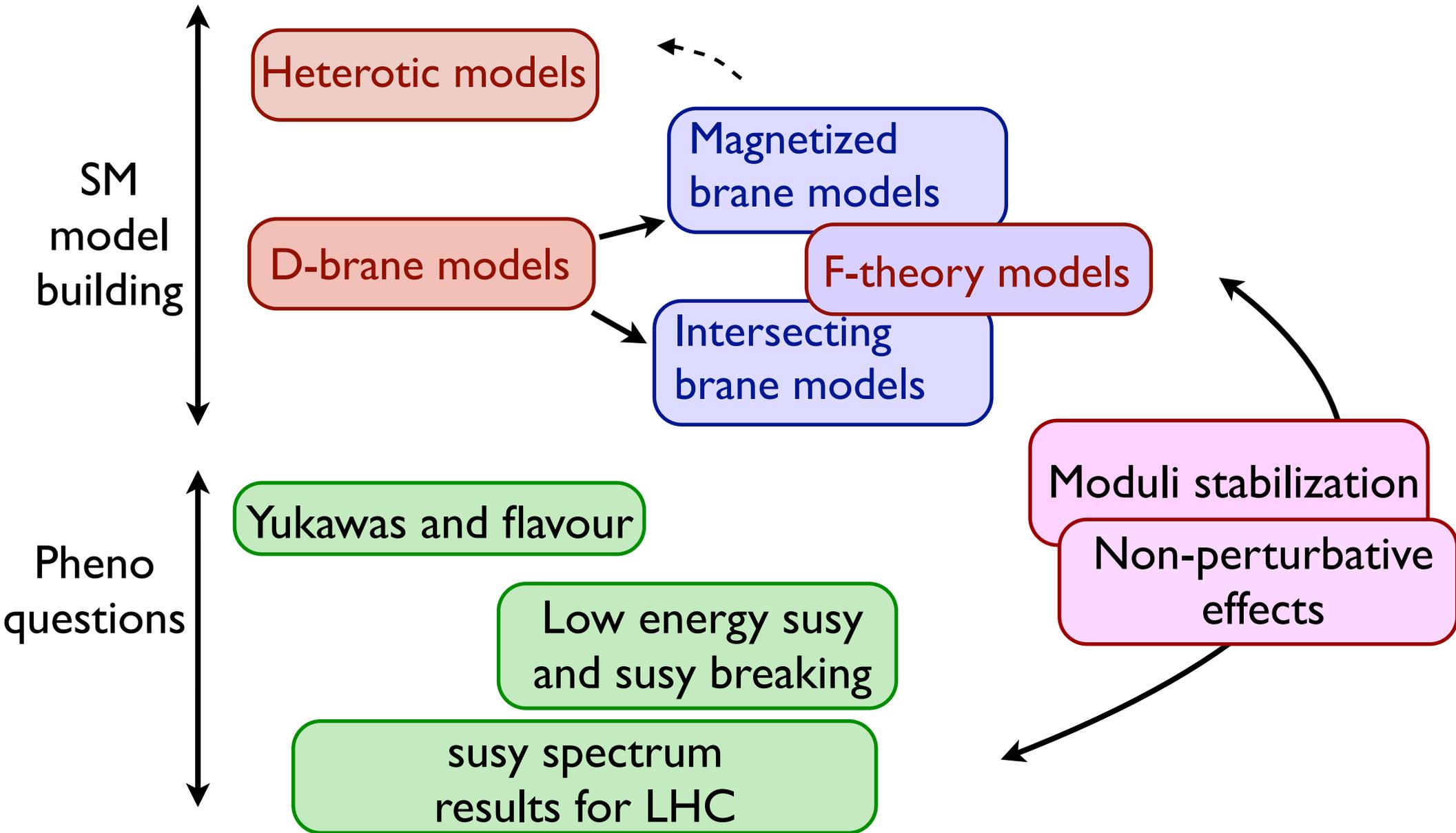
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- Overview with an eye on recent developments

Disclaimer: Necessarily incomplete

Follow particular path: not and ordered set

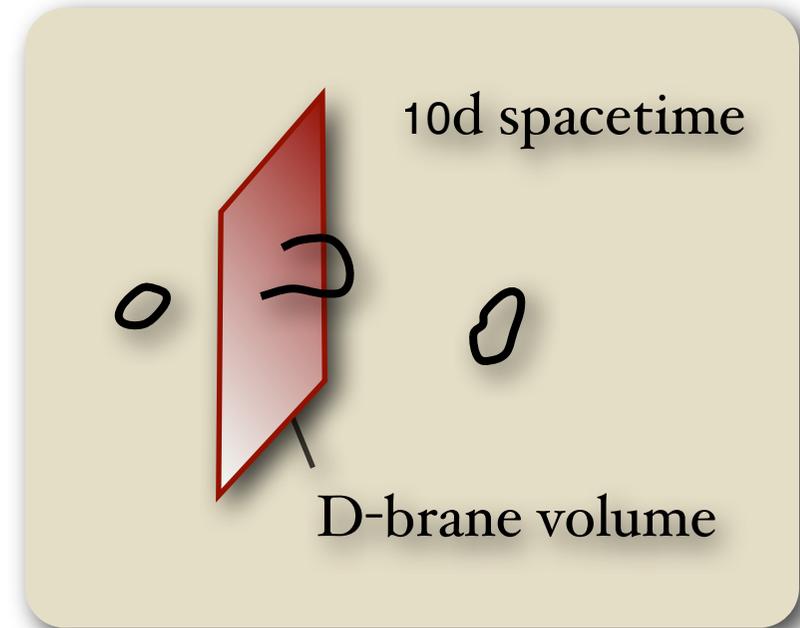
# Road map



# D-branes

Some of most successful applications are based on D-branes

High-dim. planes on which open strings end



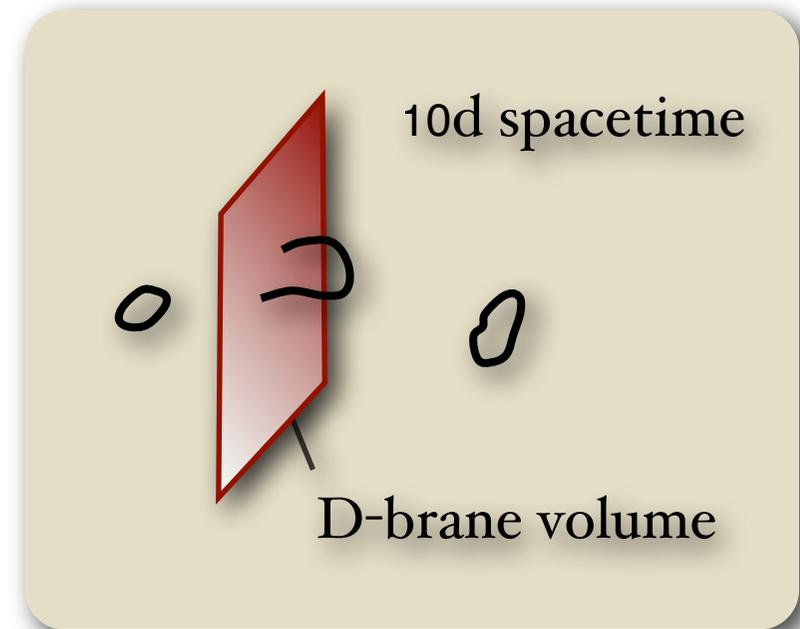
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 Brane world:

- Closed strings: gravity in 10d
- Open strings: gauge+matter on brane



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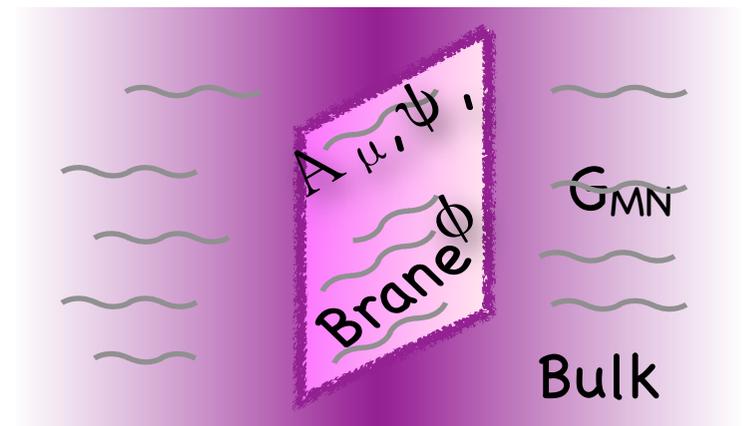
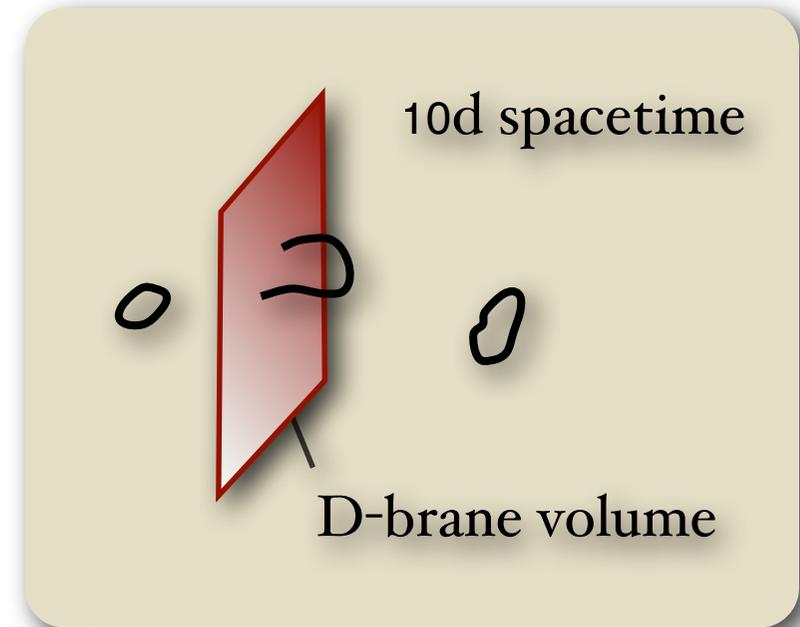
High-dim. planes on which open strings end

📌 Brane world:

- Closed strings: gravity in 10d
- Open strings: gauge+matter on brane

📌 Allows large extra dimensions

$$M_P^2 g_{SM}^2 = \frac{M_s^{11-p} V_\perp}{g_s}$$



# D-branes

 Geometrization of physical properties

E.g. Non-abelian gauge interactions  
on volume of coincident D-branes

Piling them up...

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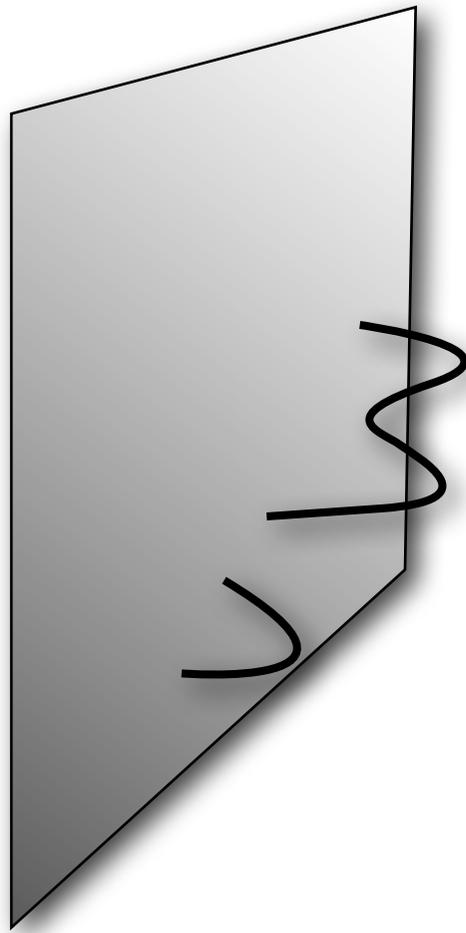
Piling them up...



... we engineer non-abelian gauge interactions!

# D-branes

$U(1)$  gauge field on the worldvolume of a single D-brane



# D-branes

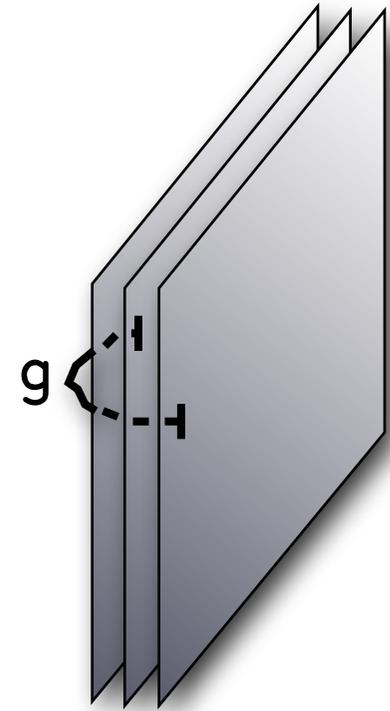
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**Electromagnetism**

# D-branes

Piling up e.g. 3 D-branes, we get 9 open string sectors

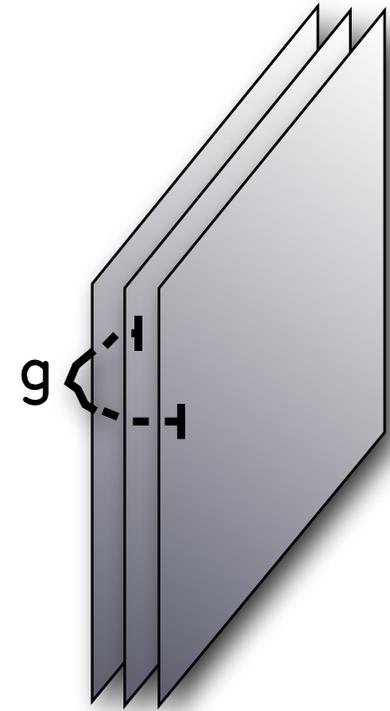


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“color-anticolor”

$$\begin{pmatrix} r\bar{r} & r\bar{g} & r\bar{b} \\ g\bar{r} & g\bar{g} & g\bar{b} \\ b\bar{r} & b\bar{g} & b\bar{b} \end{pmatrix}$$

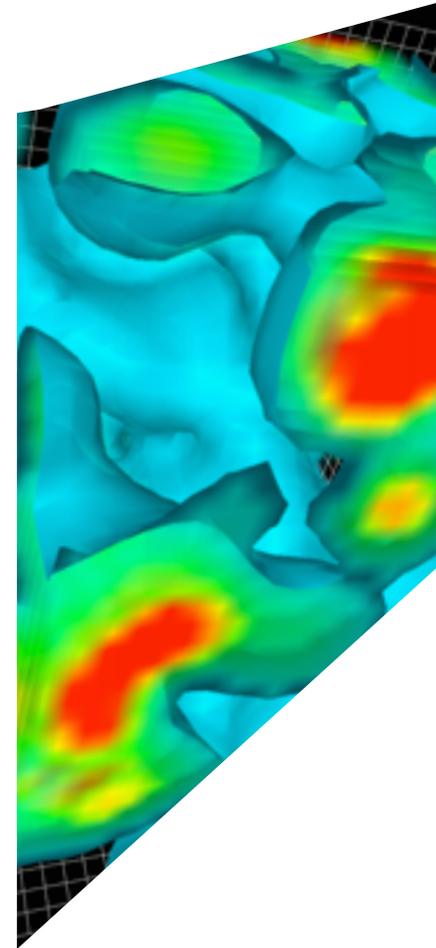


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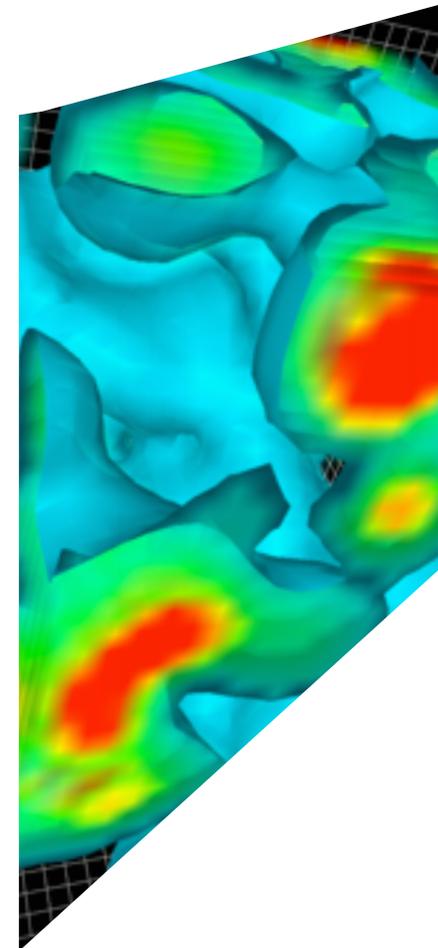
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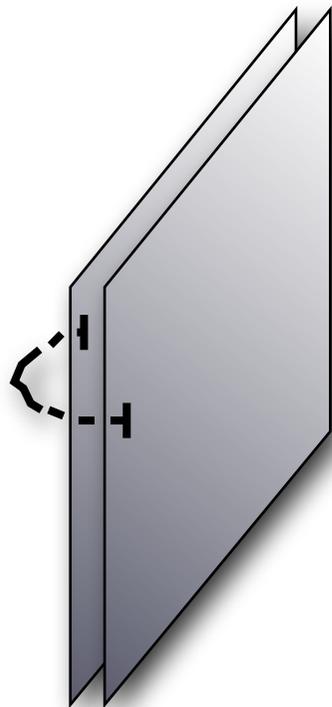
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U(3): 8 “gluons” plus a U(1)



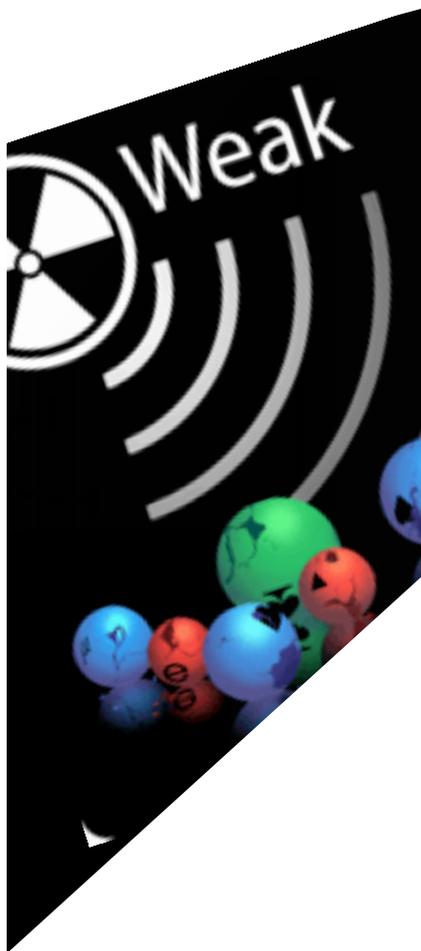
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Similarly, two extra D-branes

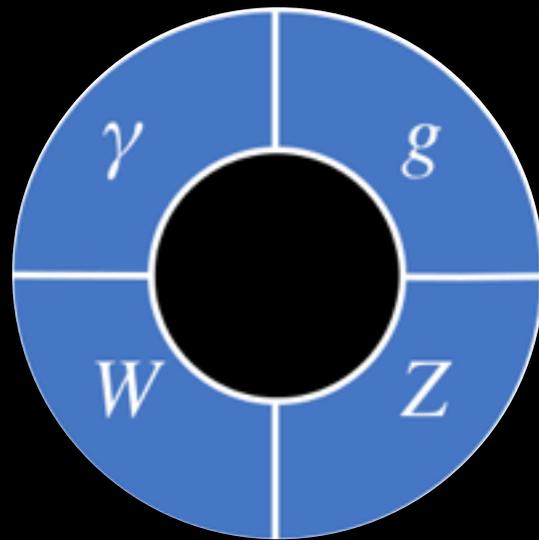


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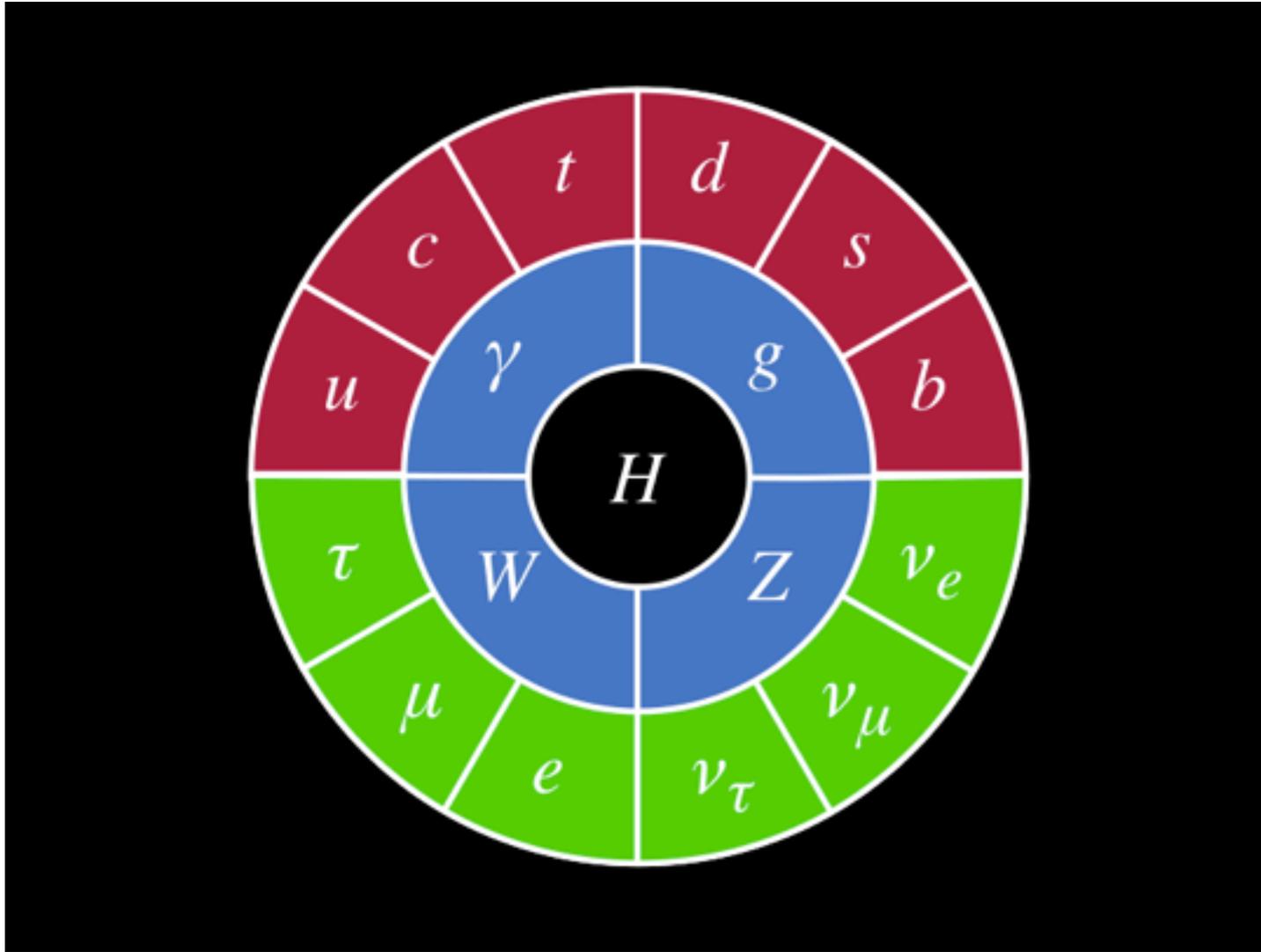


# D-branes



# D-branes

Matter? Higgs?

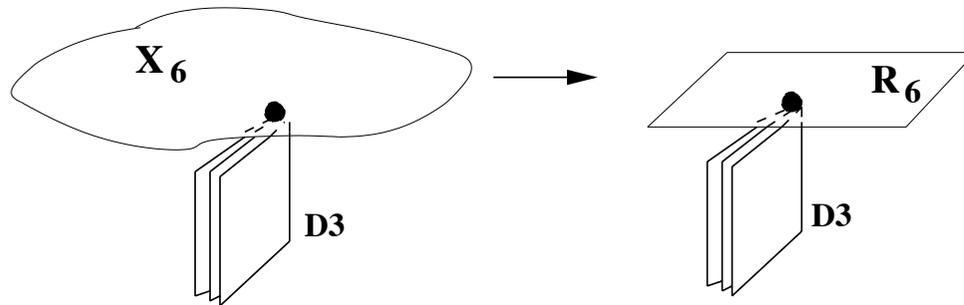


# D-branes

 Isolated D-branes in smooth geometries cannot lead to chiral gauge theories

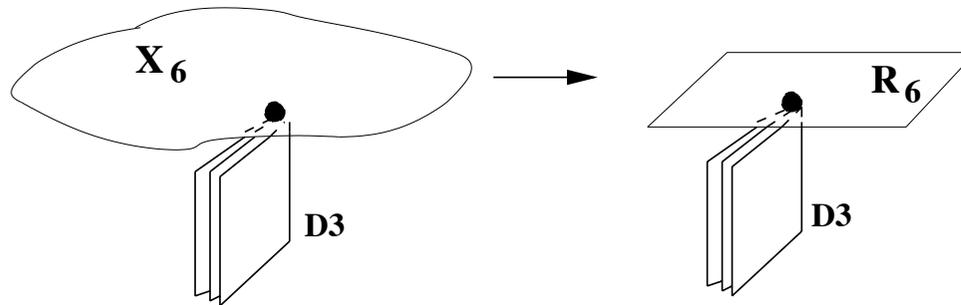
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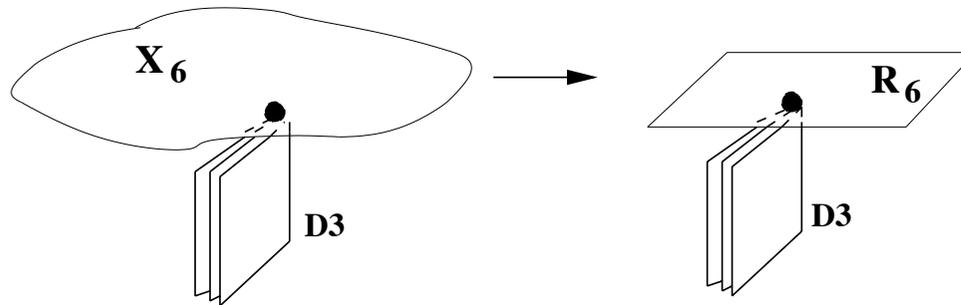


📌 **Setups for SM model building**

- D-branes at singularities
- Intersecting D-branes
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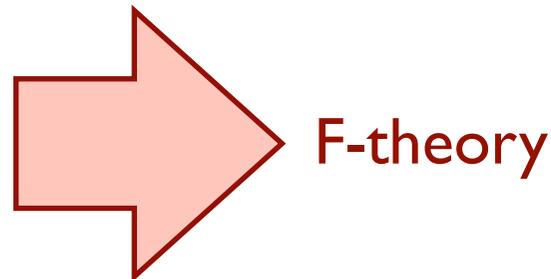


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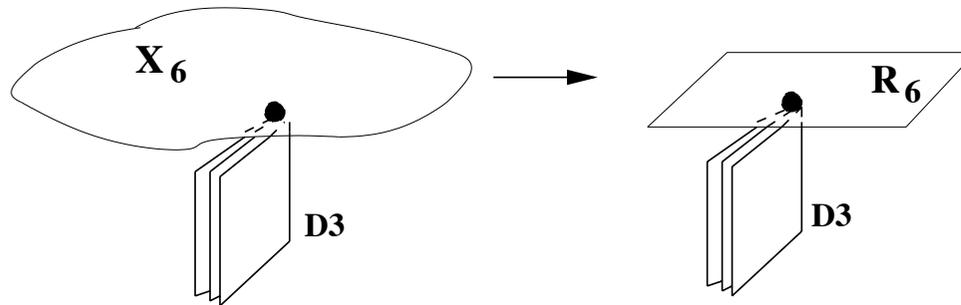
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F-theory

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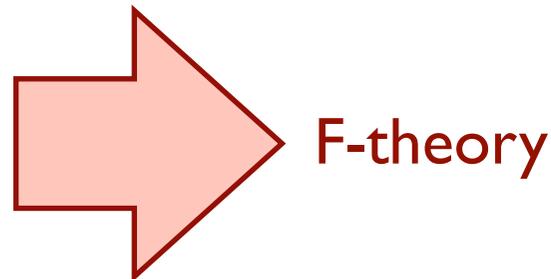


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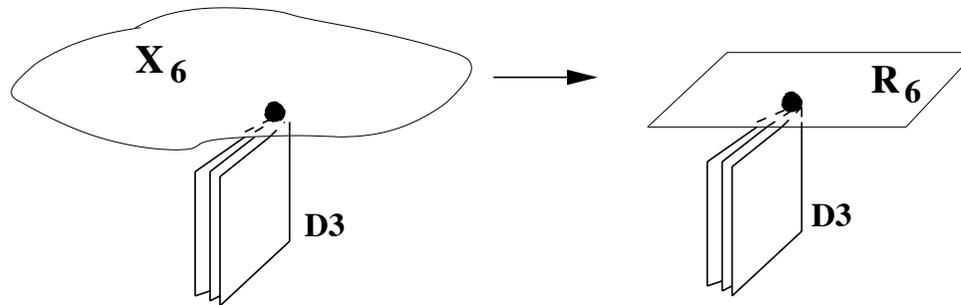
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Related to others by string dualities

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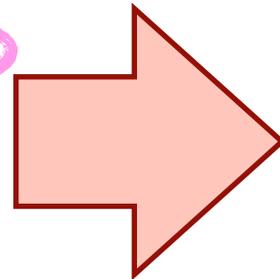
📌 Setups for SM model building

- D-branes at singularities

Now

- Intersecting D-branes

- Magnetised D-branes

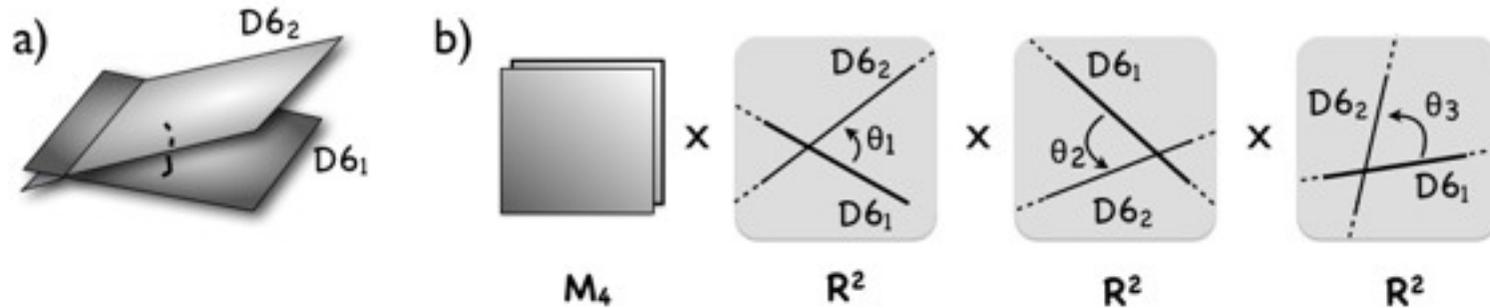


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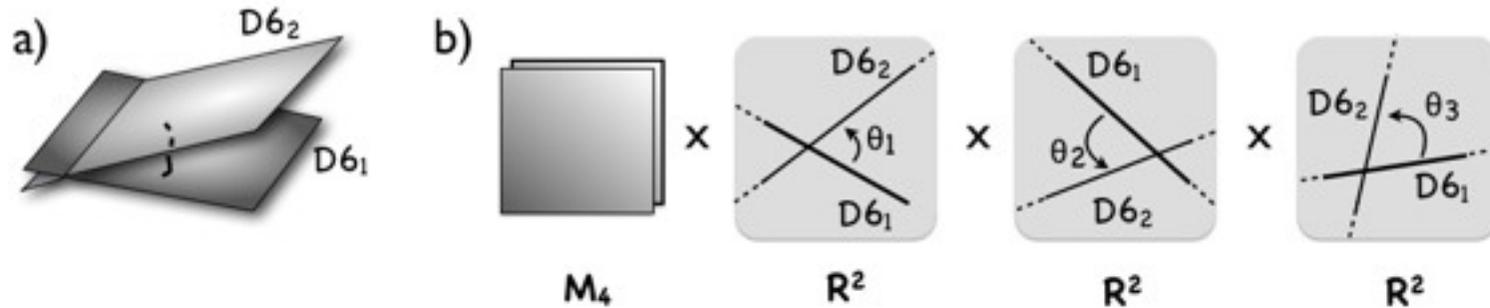
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Consider type IIA string theory with two stacks of D6-branes (hence 7d subspaces) intersecting over a 4d subspace of their volumes



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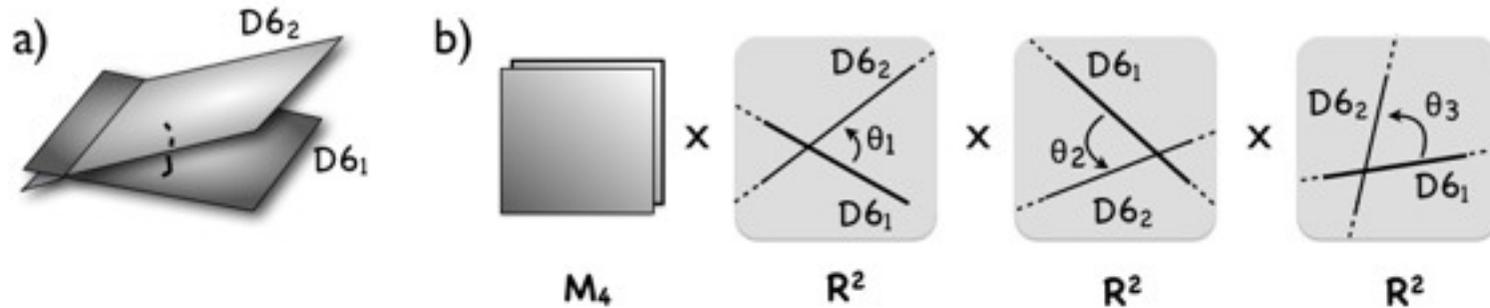


Three sectors of open strings

- $D6_1$ - $D6_1$ :  $U(N_1)$  on 7d plane 1
- $D6_2$ - $D6_2$ :  $U(N_2)$  on 7d plane 2
- $D6_1$ - $D6_2$ : 4d chiral fermion in  $(N_1, \bar{N}_2)$  on 4d intersection

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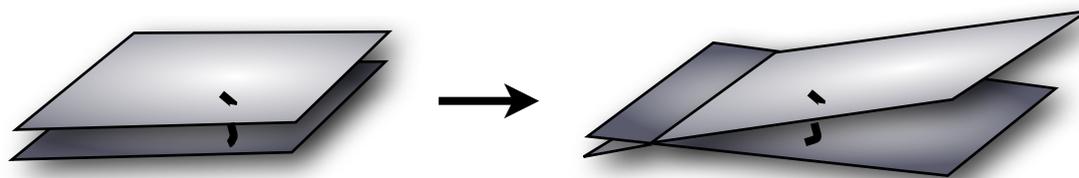


Three sectors of open strings

- D6<sub>1</sub>-D6<sub>1</sub>: U(N<sub>1</sub>) on 7d plane 1
- D6<sub>2</sub>-D6<sub>2</sub>: U(N<sub>2</sub>) on 7d plane 2
- D6<sub>1</sub>-D6<sub>2</sub>: 4d chiral fermion in (N<sub>1</sub>,  $\bar{N}_2$ ) on 4d intersection

View as 'unfolding'  $U(N_1+N_2) \rightarrow U(N_1) \times U(N_2)$

$\text{Adj} \rightarrow \text{Adj}_1 + \text{Adj}_2 + (N_1, \bar{N}_2) + \text{cc}$



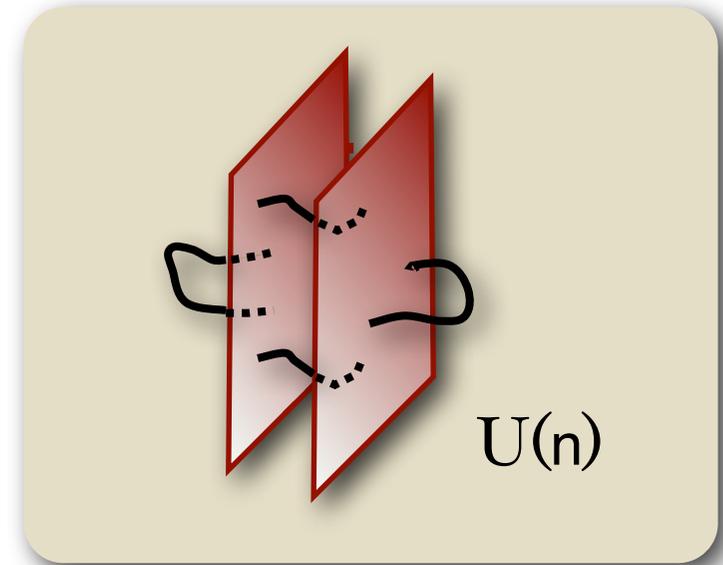
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- Non-abelian  $U(n)$  gauge interactions from “n” coincident D-branes

Matrix of open string sectors

$$\begin{pmatrix} 11 & 12 & \dots & 1n \\ 21 & 22 & \dots & 2n \\ \dots & \dots & \dots & \dots \\ n1 & n2 & \dots & nn \end{pmatrix}$$

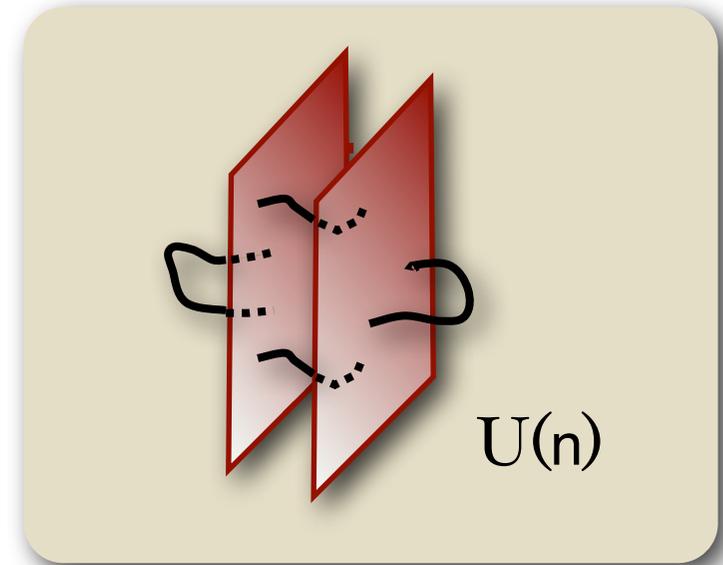


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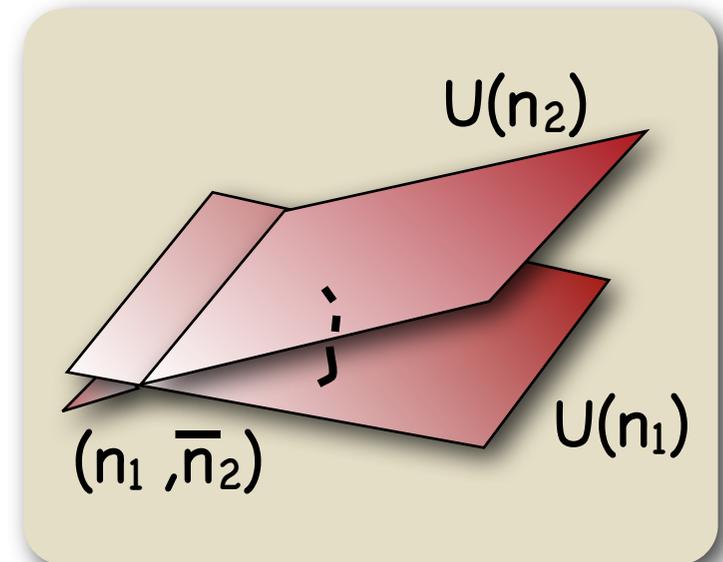
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- 4d Charged matter from intersection of stacks of D6-branes

Matrix of open string sectors

$$\begin{pmatrix} 11 & 12 & \dots & 1n_1 \\ 21 & 22 & \dots & 2n_1 \\ \dots & \dots & \dots & \dots \\ n_21 & n_22 & \dots & n_2n_1 \end{pmatrix}$$



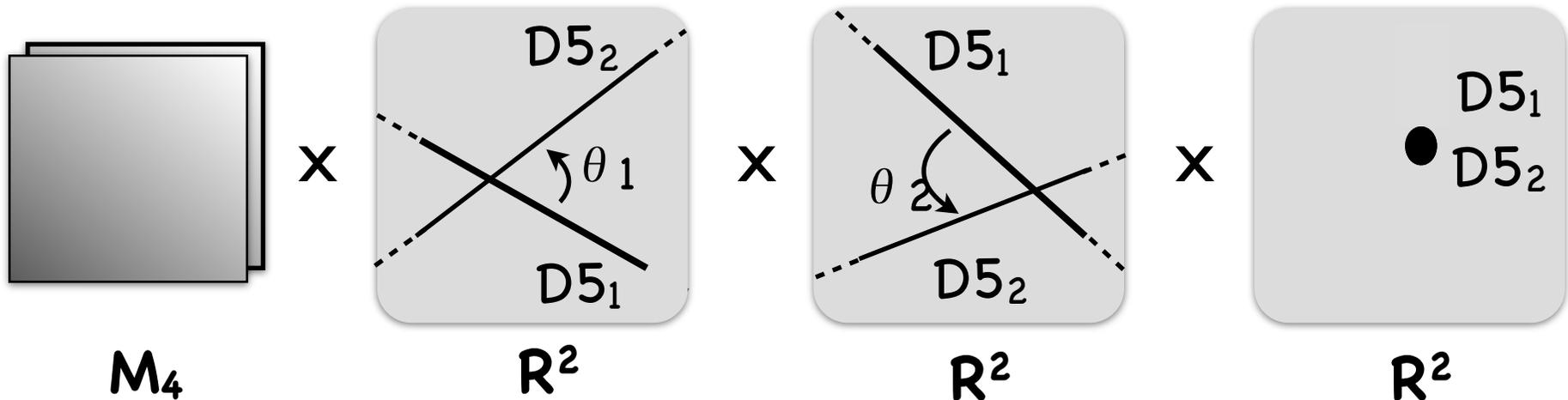
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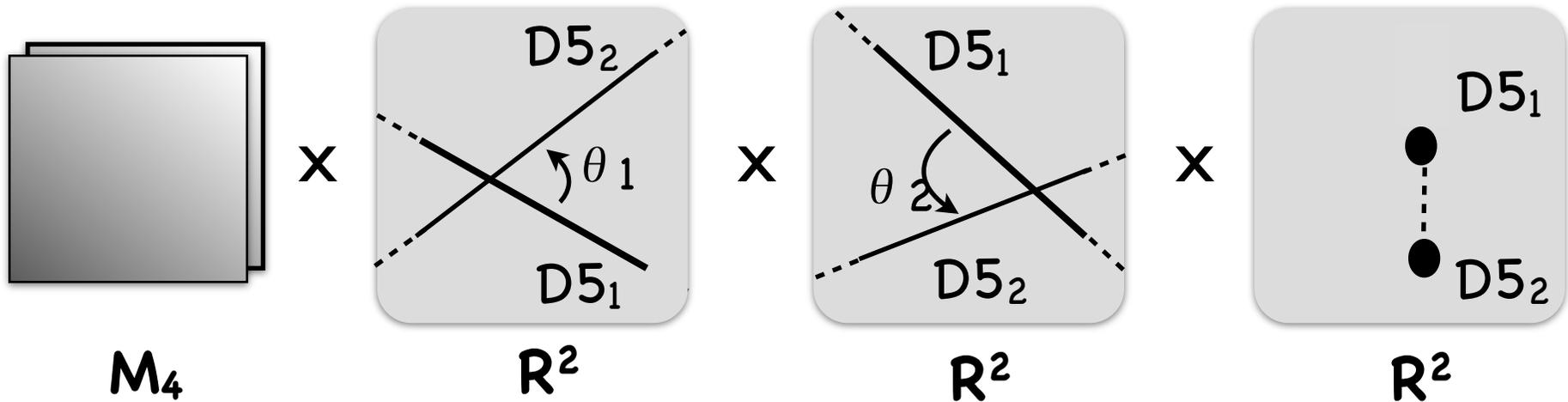
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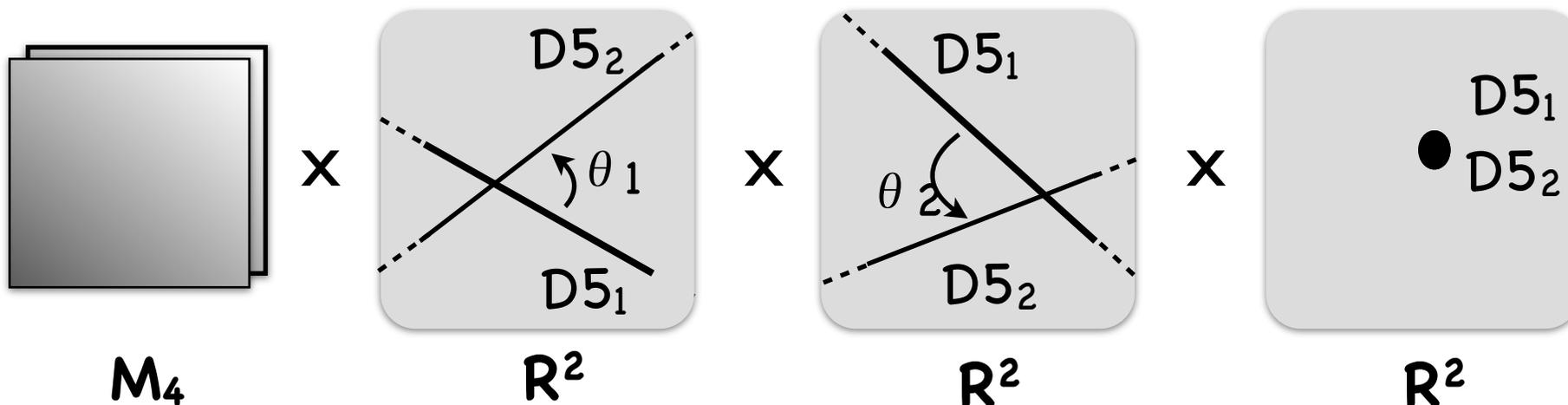
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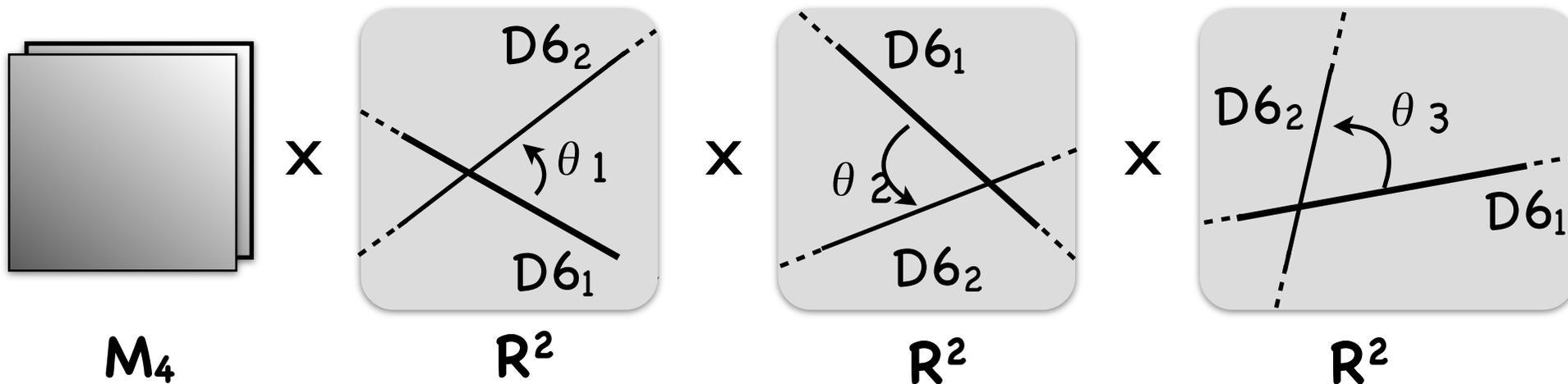
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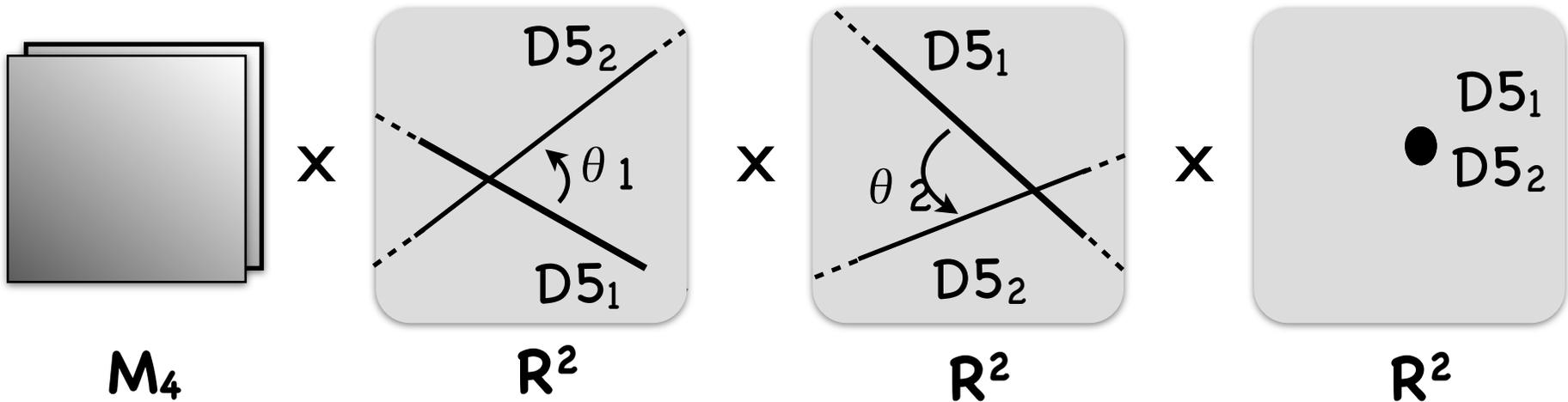
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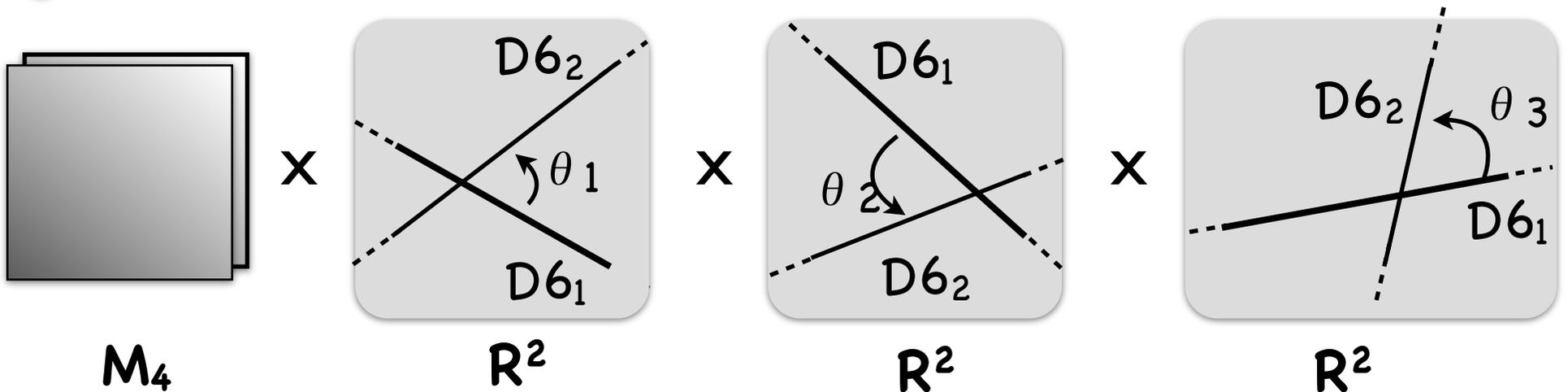
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# Scalars and supersymmetry

 In addition, the D6<sub>1</sub>-D6<sub>2</sub> sector contains 4d scalars, which are generically massive, and potentially light

$$M^2 = \frac{1}{2} (\theta_1 \pm \theta_2 \pm \theta_3) M_s^2$$

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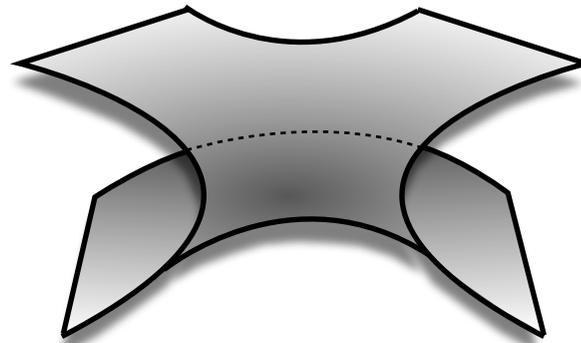
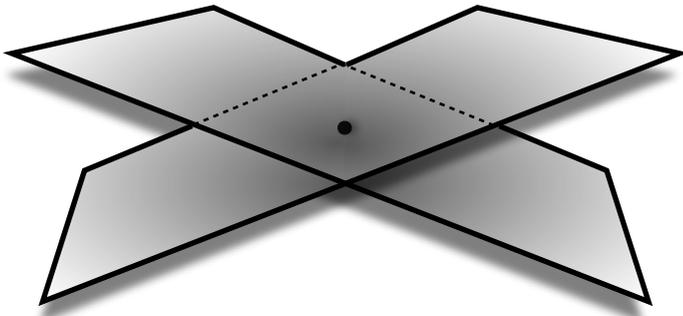
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Possible instability corresponds to possible brane recombination



Higgs effect

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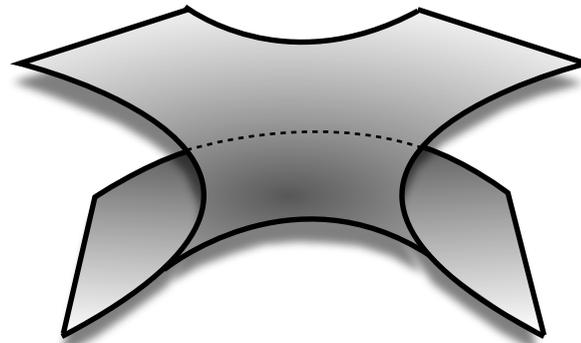
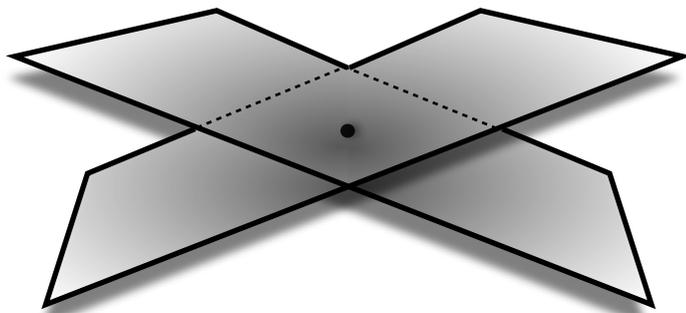
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Higgs effect

- Nice geometric interpretation in terms of volume minimization

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Geometric interpretation of structure of preserved supersymmetry

\*Local analysis. Susy of complete compact models requires O6-planes

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- Physically, condition to have one 4d massless complex scalar field

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- Geometrically, Special Lagrangian 3-cycle

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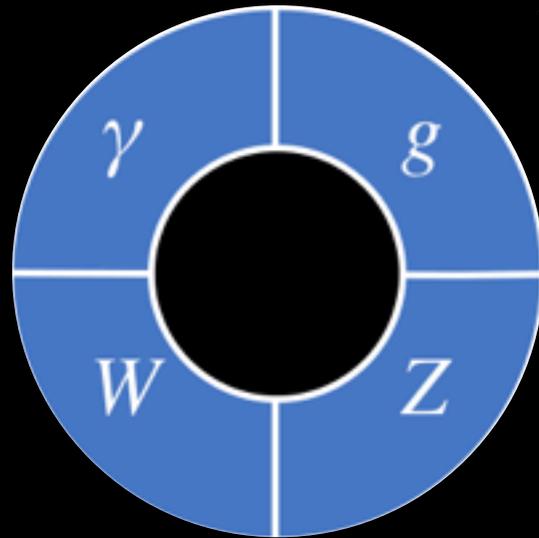
Geometric interpretation of structure of preserved supersymmetry

- Enhanced  $N=2$  supersymmetry for one zero angle e.g.

$$\theta_1 \pm \theta_2 = 0 \quad ; \quad \theta_3 = 0$$

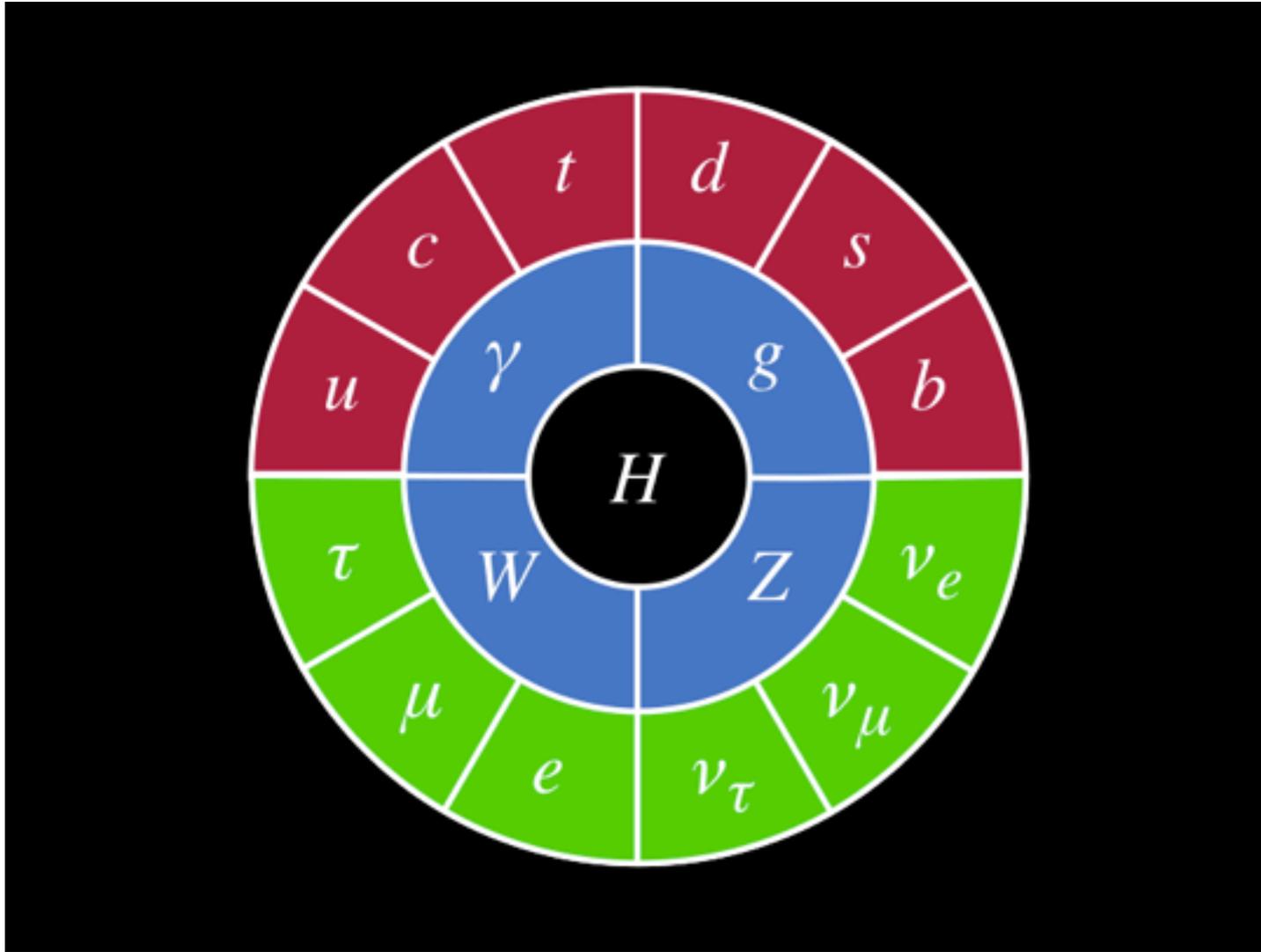
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# D-branes



# D-branes

Matter? Higgs?

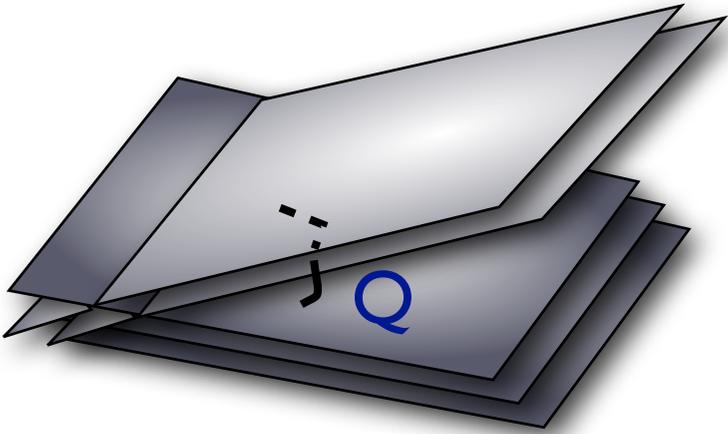


# Intersecting D-branes

- Matter & Higgs from intersections

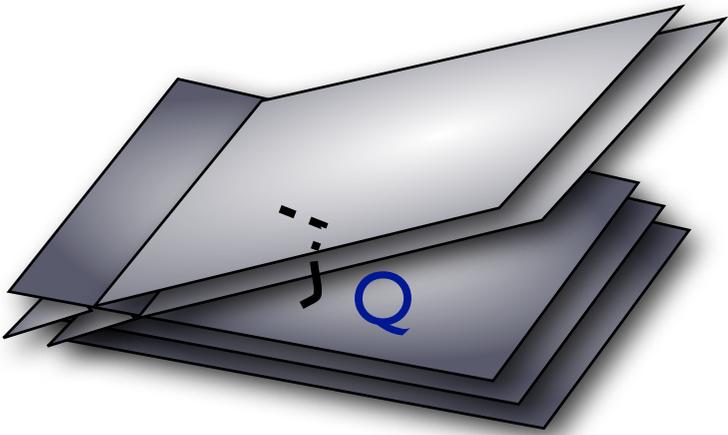
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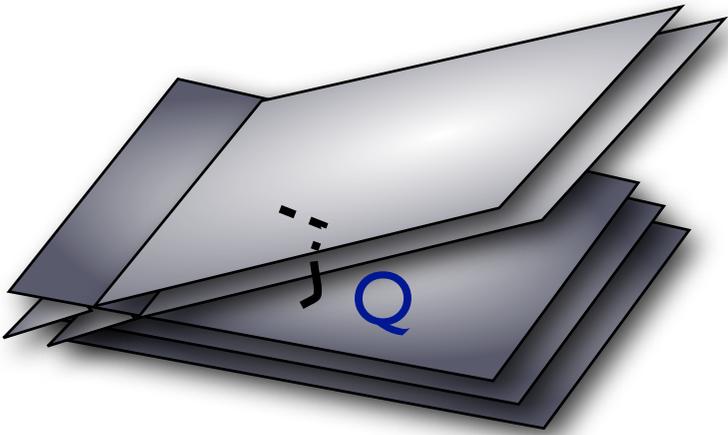
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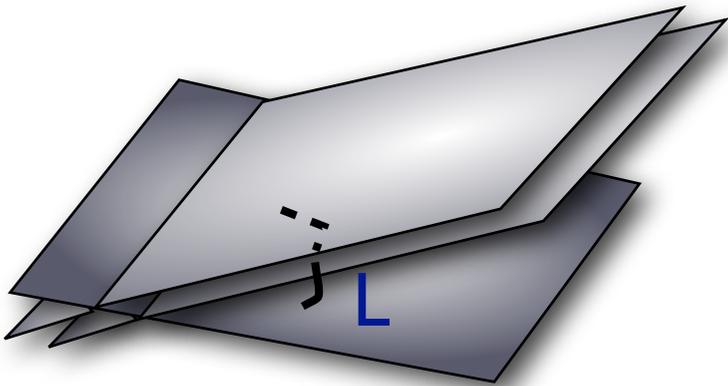
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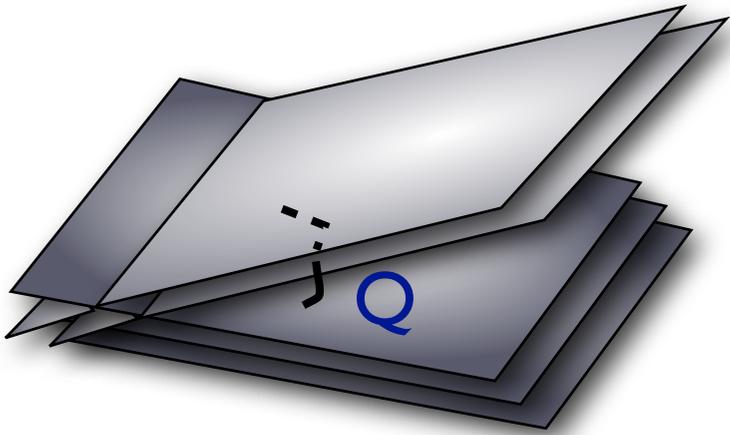


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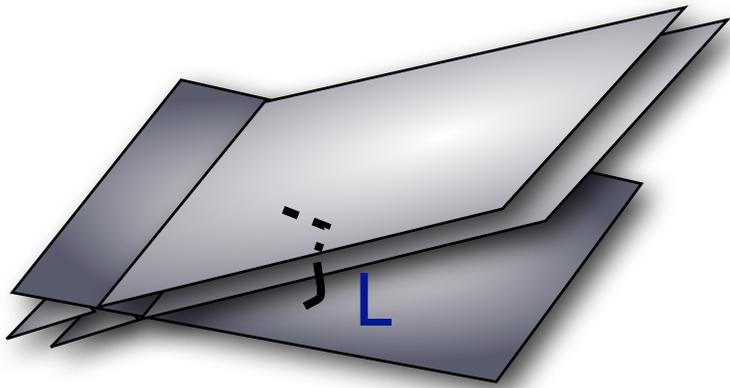


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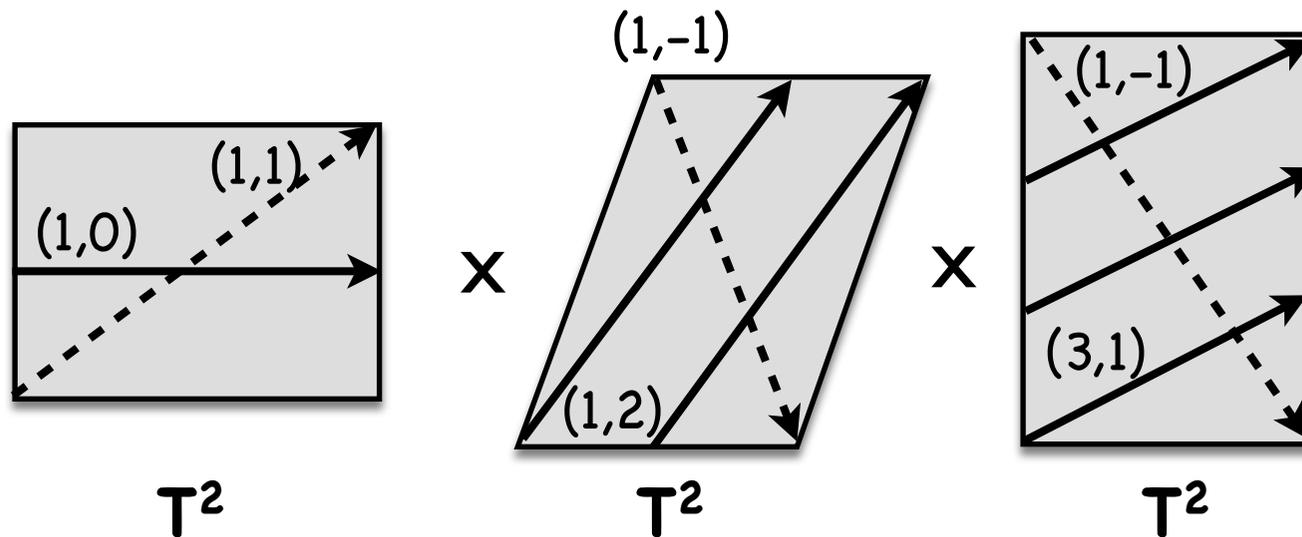
Leptones, Higgs:  
no color but EW

# Compactification

 To obtain 4d gravity, need to compactify

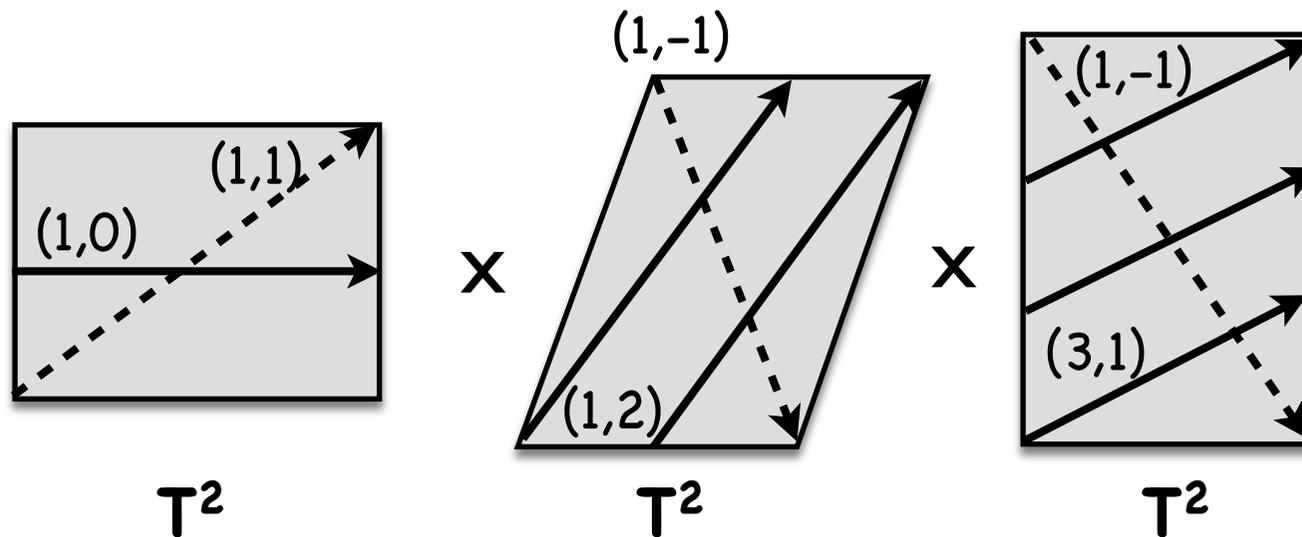
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- Configurations of D6-branes in sets of  $N_a$   $D6_a$ -branes wrapping 3-cycles  $\Pi_a$  described as products of 1-cycles  $(n_a^i, m_a^i)$  on each  $(T^2)^i$

# Compactification

- Closed string sector is maximally supersymmetric  
4d  $N=8$  supergravity multiplet  
(more realistic 4d  $N=1$  supergravity in non-toroidal models:  
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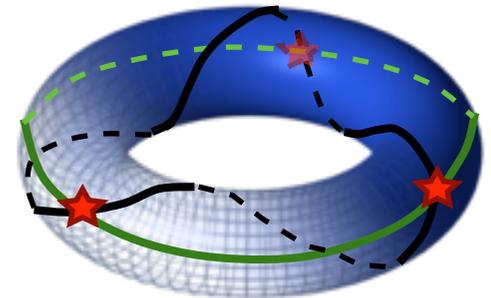
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Intersection number = geometric origin of family replication!

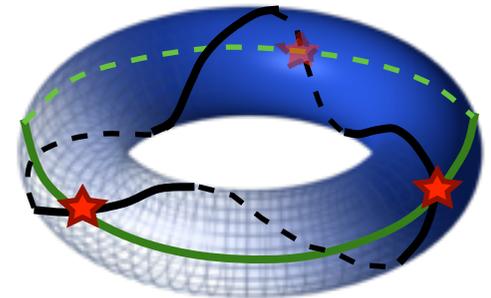
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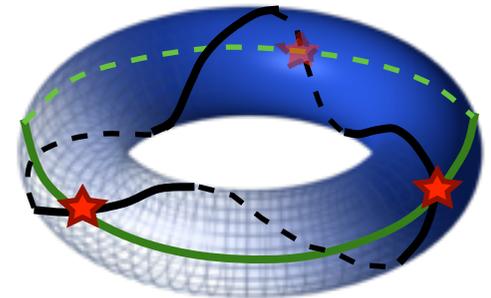
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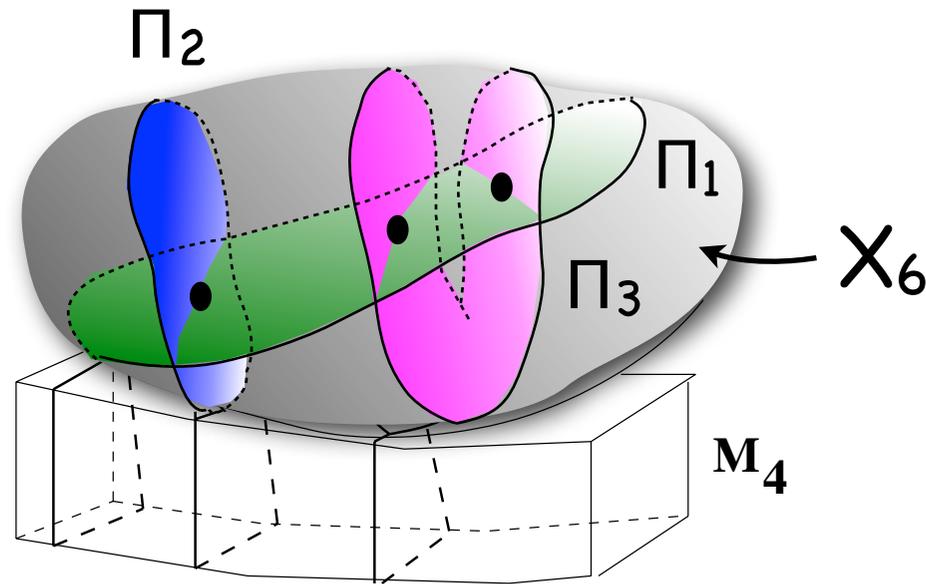
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- Non-chiral features of the spectrum (susy, scalars,...) depend on  
detailed configuration

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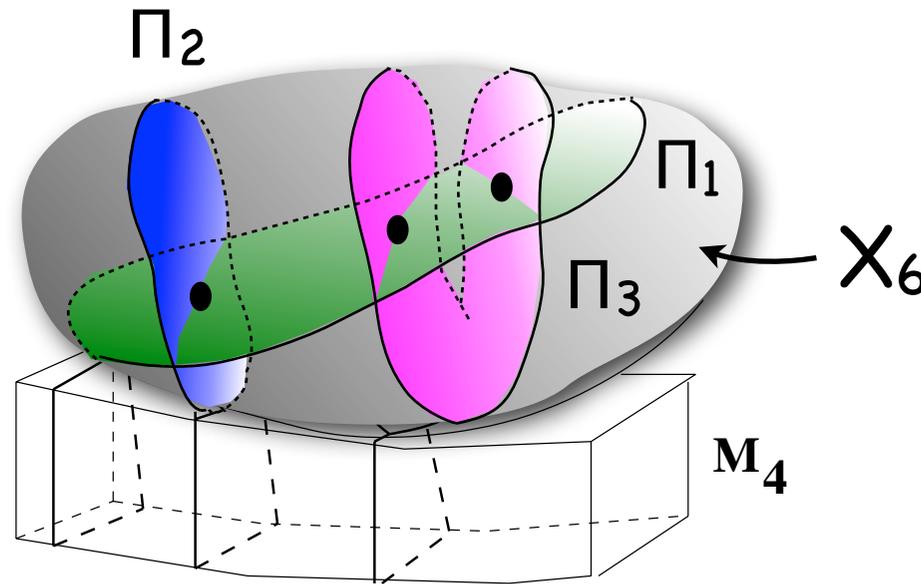
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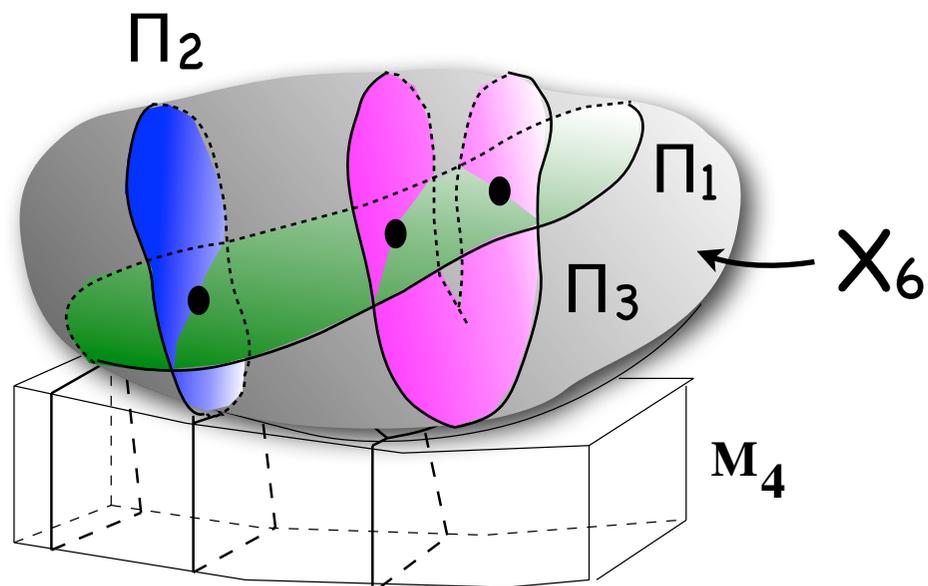


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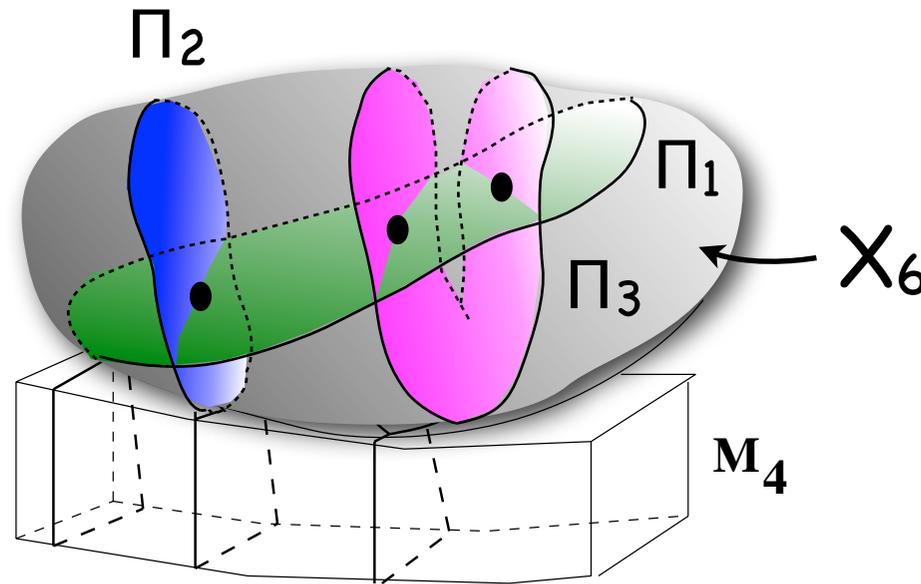
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General class of string compactifications with non-abelian gauge symmetry and replicated charged chiral fermions

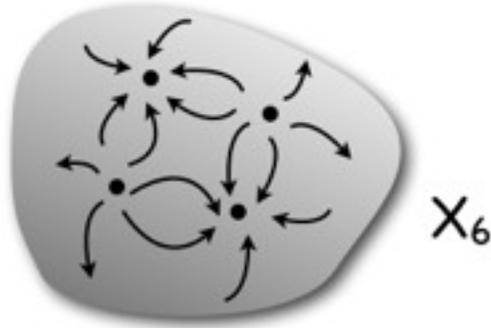
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# Tadpoles, anomalies and all that

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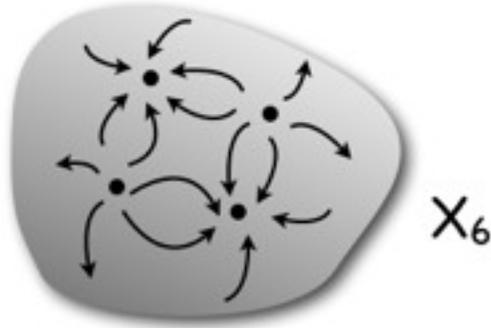
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Gauss' law

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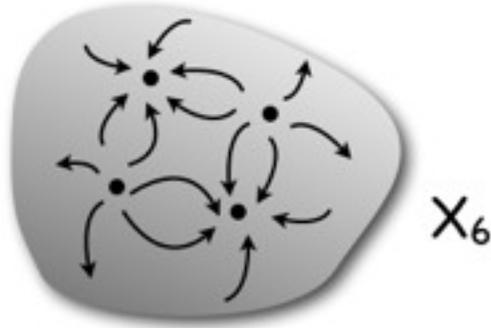


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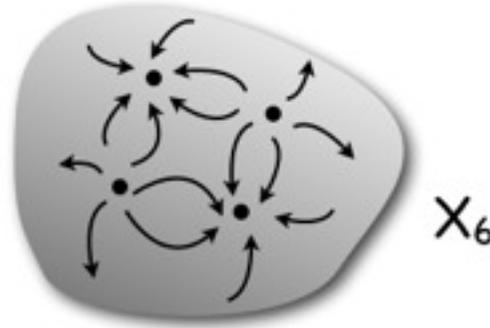


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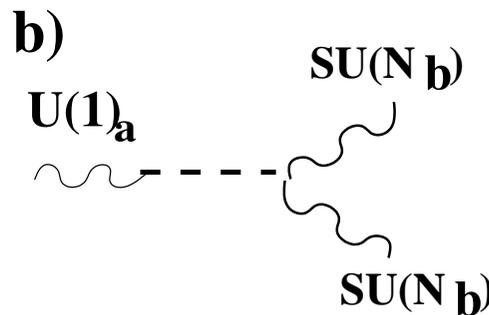
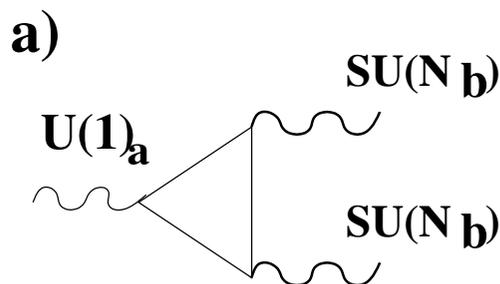
Gauss' law

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- $SU(N_a)^3$  non-abelian vanish identically

- $U(1)_a$ - $SU(N_b)^2$  mixed cancel via Green-Schwarz mechanism involving the D6-brane couplings

$$\sum_{k,a} \int_{4d} B_k \wedge \text{tr} F_a \quad \& \quad \sum_{k,a} \int_{4d} a_k \text{tr} (F_b \wedge F_b)$$



# Tadpoles, anomalies and all that

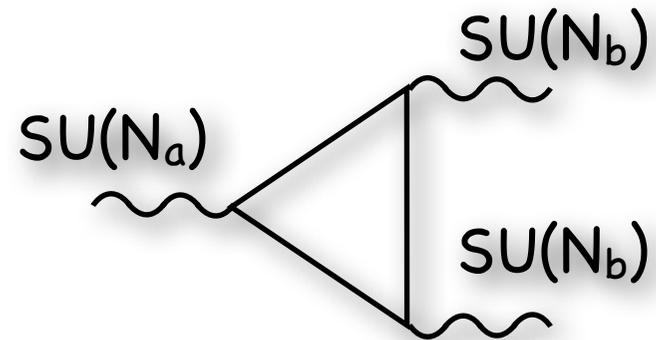
## RR tadpole cancellation

$$dF_2 = \sum_a N_a \delta(\Pi_a) \rightarrow \sum_a N_a [\Pi_a] = 0$$

## Non-abelian anomaly cancellation

$$SU(N_a)^3 : A_a = \sum_b I_{ab} N_b$$

$$0 = [\Pi_a] \cdot \sum_b N_b [\Pi_b] = \sum_b N_b I_{ab}$$

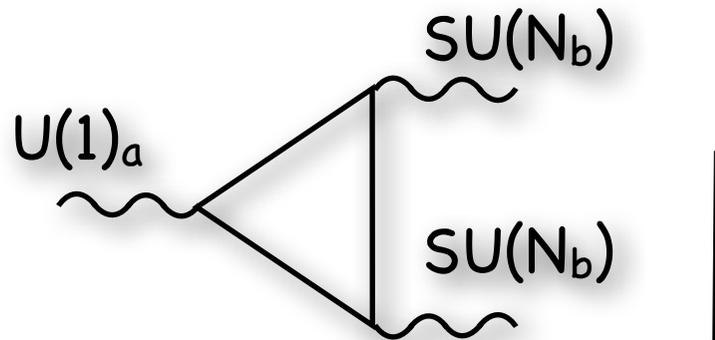


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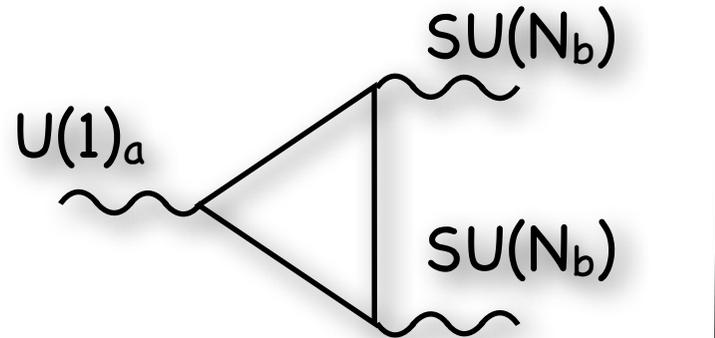
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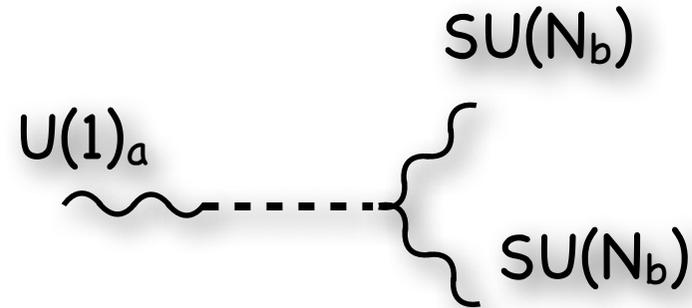
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# Tadpoles, anomalies and all that

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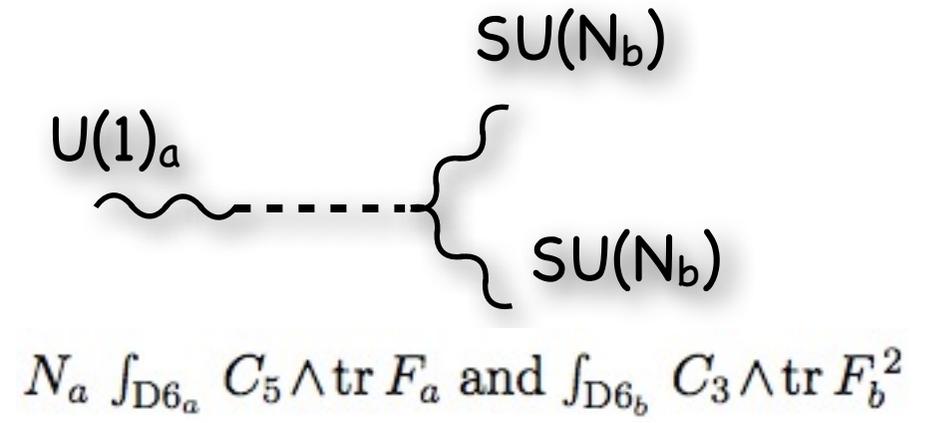
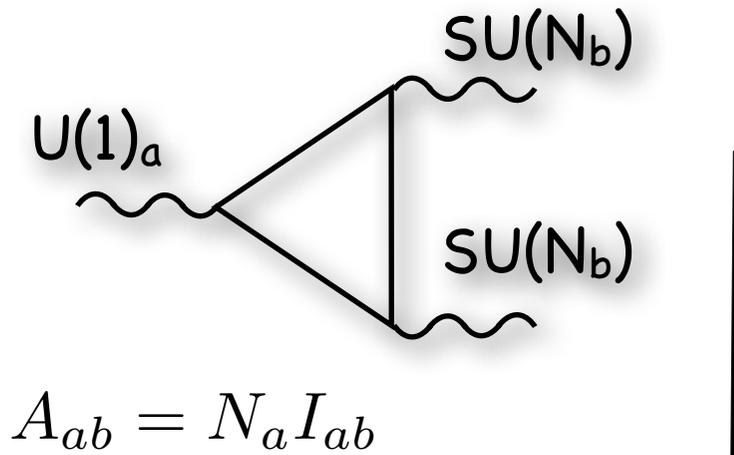
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$$N_a \int_{D6_a} C_5 \wedge \text{tr} F_a \text{ and } \int_{D6_b} C_3 \wedge \text{tr} F_b^2$$

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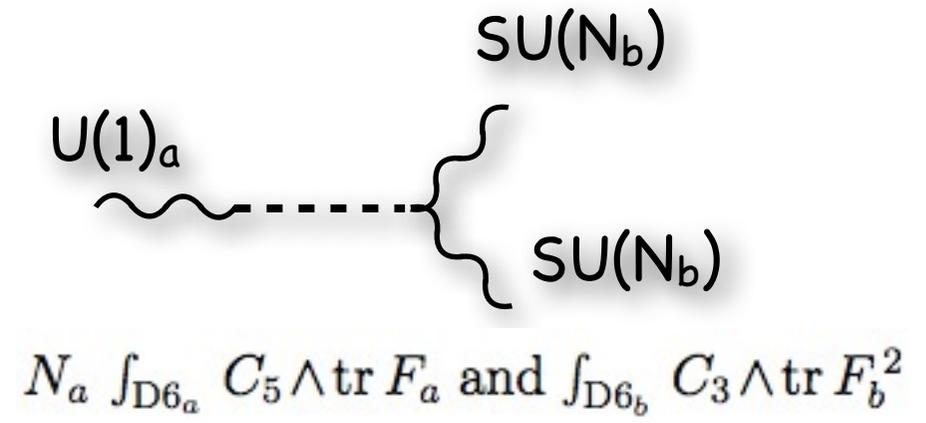
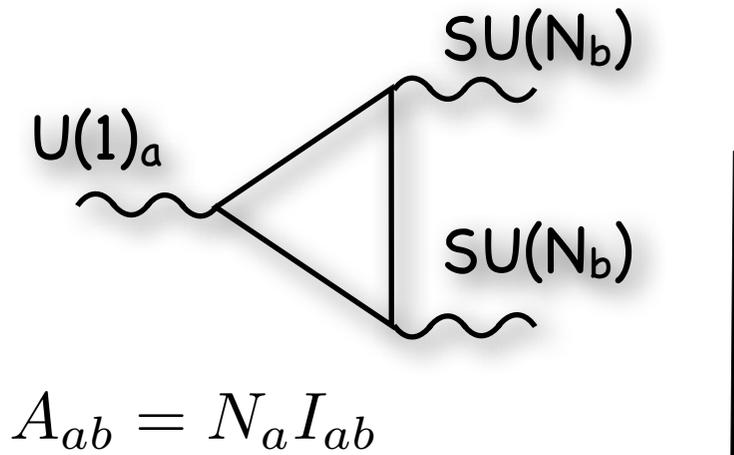


Introduce basis

$$(B_2)_k = \int_{[\Lambda_k]} C_5 \quad , \quad \phi_{\tilde{l}} = \int_{[\Lambda_{\tilde{l}}]} C_3 \quad \text{with } d\phi_{\tilde{l}} = -\delta_{k\tilde{l}} *_{4d} (B_2)_k$$

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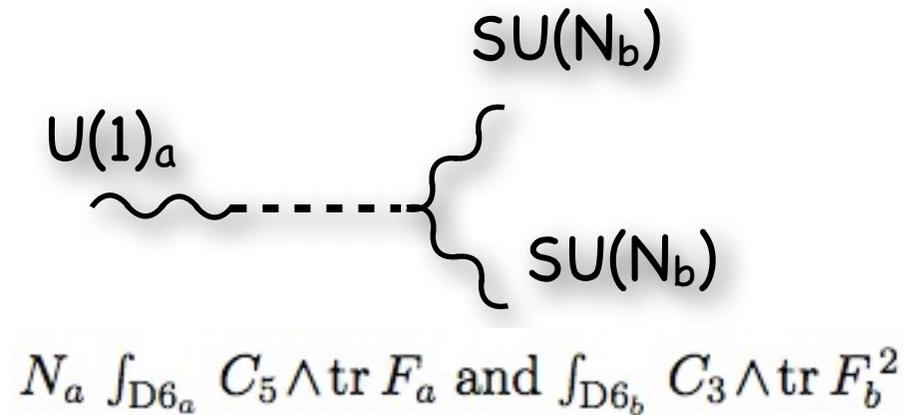
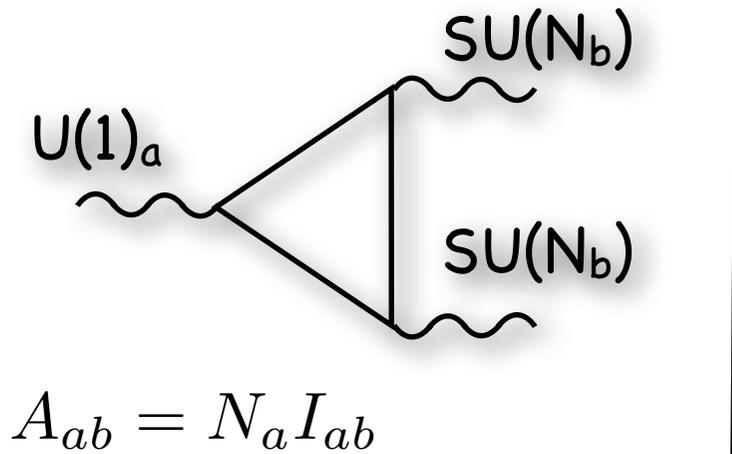
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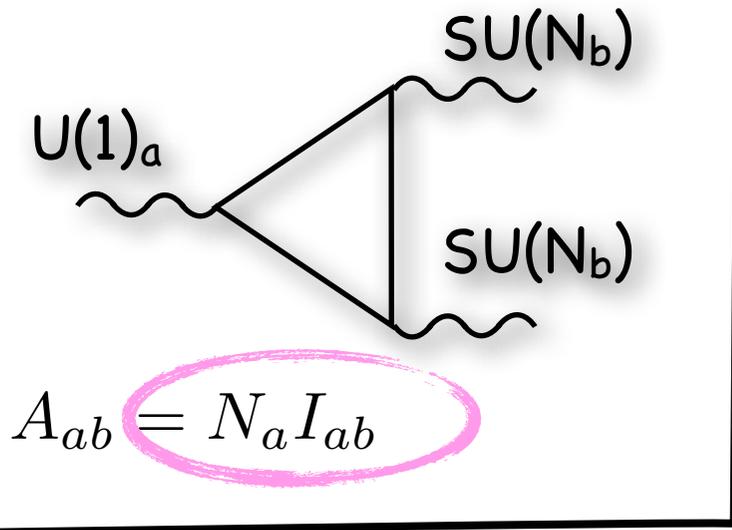
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 Due to BF couplings, all 'anomalous' and some 'non-anomalous' U(1)'s become massive, with mass of order the string scale

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## Consequences

- Impose that hypercharge generator remains massless
- Additional U(1)'s removed  
remain as global symmetries exact in perturbation theory
- Operators violating the latter can appear non-perturbatively  
D-brane instantons, see later

# Compact orientifolds

## Compact susy models using orientifold projections

Take type IIA on CY  $X$ , and mod out by  $\Omega R$ , where  $\Omega$  flips the string orientation, and  $R$  is a  $Z_2$  symmetry of the CY  $X$

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Supersymmetric D6/O6 systems with zero total charge/tension can still be non-trivial
- 📌 Technical modifications of configuration, tadpoles, anomalies,... skip!
  - Every D6-brane stack has an orientifold image  $D6'$
  - Intersections  $ab$  and  $ab'$
  - D6-brane stacks on top of O6 have  $SO/Sp$  group rather than  $U(N)$

# Towards the SM

## Simplest road to SM

[Ibáñez, Marchesano, Rabadán; Cremades, Ibáñez, Marchesano; '01]

Introduce four stacks of D6's a,b,c,d with

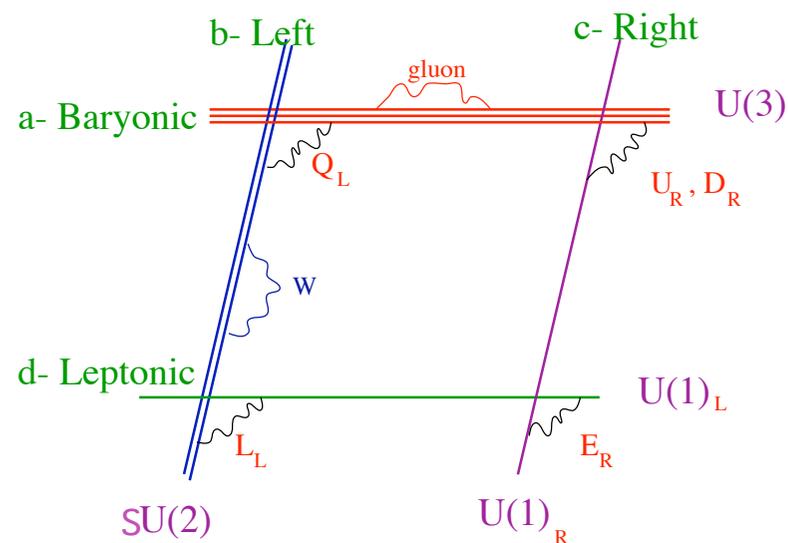
$$U(3)_a \times USp(2)_b \times U(1)_c \times U(1)_d$$

$$I_{ab} = 3 \rightarrow Q_L$$

$$I_{ac} = -3, I_{ac'} = 3 \rightarrow U_R, D_R$$

$$I_{db} = 3 \rightarrow L$$

$$I_{dc} = -3, I_{dc'} = -3 \rightarrow E_R, \nu_R$$



Spectrum of SM with hypercharge

$$Y = \frac{1}{6}Q_a - \frac{1}{2}Q_c - \frac{1}{2}Q_d$$

Look for models realizing these intersection numbers:  
 $\Rightarrow$  surprisingly easy!

# Towards the SM

## Explicit realization in toroidal models

$N_\alpha$	$(n_\alpha^1, m_\alpha^1)$	$(n_\alpha^2, m_\alpha^2)$	$(n_\alpha^3, m_\alpha^3)$
$N_a = 3$	(1,0)	(1,3)	(1,-3)
$N_b = 1$	(0,1)	(1,0)	(0,-1)
$N_c = 1$	(0,1)	(0,-1)	(1,0)
$N_d = 1$	(1,0)	(1,3)	(1,-3)

$$I_{ab} = I_{ab'} = 3 \quad I_{ac} = I_{ac'} = -3$$

$$I_{db} = I_{db'} = -3 \quad I_{cd} = 3; I_{cd'} = -3$$

# Towards the SM

## Explicit realization in toroidal models

$$\begin{array}{cccc}
 N_\alpha & (n_\alpha^1, m_\alpha^1) & (n_\alpha^2, m_\alpha^2) & (n_\alpha^3, m_\alpha^3) \\
 N_a = 3 & (1, 0) & (1, 3) & (1, -3) \\
 N_b = 1 & (0, 1) & (1, 0) & (0, -1) \\
 N_c = 1 & (0, 1) & (0, -1) & (1, 0) \\
 N_d = 1 & (1, 0) & (1, 3) & (1, -3)
 \end{array}$$

$$I_{ab} = I_{ab'} = 3 \quad I_{ac} = I_{ac'} = -3$$

$$I_{db} = I_{db'} = -3 \quad I_{cd} = 3; I_{cd'} = -3$$

Intersection	Matter fields		$Q_a$	$Q_c$	$Q_d$	Y	Y'	$3Q_a - Q_d$
$ab, ab'$	$Q_L$	$3(3, 2)$	1	0	0	1/6	1/3	3
$ac$	$U_R$	$3(\bar{3}, 1)$	-1	1	0	-2/3	2/3	-3
$ac'$	$D_R$	$3(\bar{3}, 1)$	-1	-1	0	1/3	-4/3	-3
$bd, b'd$	$L$	$3(1, 2)$	0	0	-1	-1/2	-1	1
$cd$	$E_R$	$3(1, 1)$	0	-1	1	1	0	-1
$cd'$	$\nu_R$	$3(1, 1)$	0	1	1	0	2	-1

# Towards the SM

## Explicit realization in toroidal models

$N_\alpha$	$(n_\alpha^1, m_\alpha^1)$	$(n_\alpha^2, m_\alpha^2)$	$(n_\alpha^3, m_\alpha^3)$
$N_a = 3$	(1,0)	(1,3)	(1,-3)
$N_b = 1$	(0,1)	(1,0)	(0,-1)
$N_c = 1$	(0,1)	(0,-1)	(1,0)
$N_d = 1$	(1,0)	(1,3)	(1,-3)

(need few extra branes, adding few extra matter)

## Supersymmetric for suitable choices of $T^2$ geometry

MSSM with pair of Higgs doublets in non-chiral bc sector

## Realizations in other setups

Orientifolds of type II Gepner models leads to some 200.000 models

## Towards more detailed analysis of properties

# Towards the SM

# Towards the SM

## Gauge couplings

A priori no natural unification at string scale

Each coupling depends on wrapped volume

$$\frac{1}{g_a^2} = \frac{V_{\Pi a}}{g_s}$$

# Towards the SM

## Gauge couplings

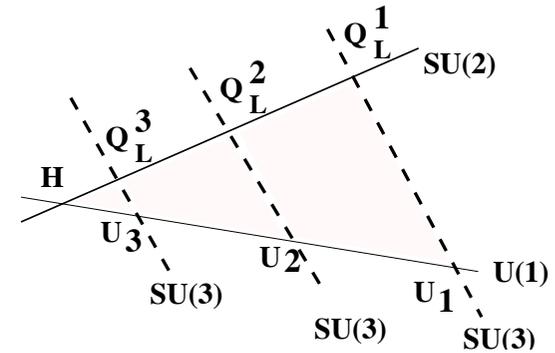
A priori no natural unification at string scale  
Each coupling depends on wrapped volume

$$\frac{1}{g_a^2} = \frac{V_{\Pi a}}{g_s}$$

## Yukawa couplings

Mediated by open string worldsheet  
instantons

$$Y_{jk} \simeq e^{-A_{Hjk} + i\phi_{jk}}$$



# Towards the SM

## Gauge couplings

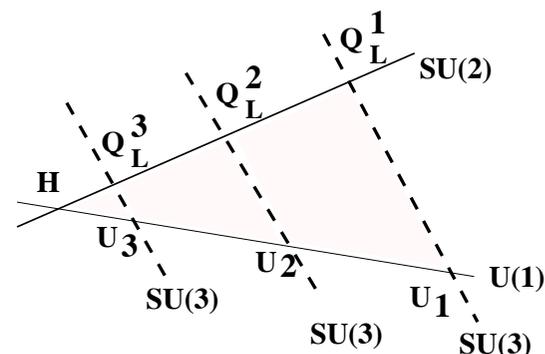
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## Yukawa couplings

Mediated by open string worldsheet  
instantons

$$Y_{jk} \simeq e^{-A_{Hjk} + i\phi_{jk}}$$



Exponential dependence potentially explains fermion mass hierarchy  
Extensive computations of general formula

Complicated functions of open and closed string (Kahler) moduli  
Need to face moduli stabilization

Geometrical interpretation possibly useful in search for textures

# Towards the SM

# Towards the SM



## String scale

- In susy models, can 'choose' string scale  
(until susy breaking is specified)
- In non-susy models, can alleviate hierarchy by large volume

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Forbidden in perturbation theory by  $U(1)_a$  baryon number  
Violated by instantons, just like in SM

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## Susy breaking and soft terms

# Towards the SM

- 📌 String scale
  - In susy models, can 'choose' string scale (until susy breaking is specified)
  - In non-susy models, can alleviate hierarchy by large volume

- 📌 Proton decay *Lectures 2,3*
  - Forbidden in perturbation theory by  $U(1)_a$  baryon number
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- 📌 Susy breaking and soft terms

# Towards the SM

- 📌 String scale
  - In susy models, can 'choose' string scale (until susy breaking is specified)
  - In non-susy models, can alleviate hierarchy by large volume

- 📌 Proton decay *Lectures 2,3*
  - Forbidden in perturbation theory by  $U(1)_a$  baryon number
  - Violated by instantons, just like in SM

- 📌 Susy breaking and soft terms *Lecture 4*



U,D

$\gamma$

E,  $\nu_R$

Q

Z

Thanks!

W

Q

L