

# SUPERSYMMETRIC QUANTUM MECHANICS AND MORSE THEORY

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**HOMEWORK 3:  $d = 0$  SSUPERSYMMETRY**

Due date: July 13

# 1 Linking number

**Introduction:** For pair of loops (curve with no boundary)  $C_1$  and  $C_2$ , given in parametric form

$$C_1 = S^1 \rightarrow \mathbb{R}^3 : t \mapsto \gamma_1^i(t), \quad C_2 = S^1 \rightarrow \mathbb{R}^3 : t \mapsto \gamma_2^i(t) \quad (1.1)$$

the formula for the linking number is well known

$$\text{lk}(C_1, C_2) = \frac{1}{4\pi} \int_0^{2\pi} dt_1 \int_0^{2\pi} dt_2 \sum_{ijk} \frac{\epsilon_{ijk} \dot{\gamma}_1^i(t_1) \dot{\gamma}_2^j(t_2) (\gamma_1(t_1) - \gamma_2(t_2))^k}{|\gamma_1(t_1) - \gamma_2(t_2)|^3}. \quad (1.2)$$

On the lecture we briefly discussed the relation between the linking number and intersection number. The goal of this problem is to derive the formula above from the intersection theory.

**Construction:** The Poincare lemma for  $\mathbb{R}^3$  implies that

$$H_1(\mathbb{R}^3) = H^1(\mathbb{R}^3) = 0, \quad (1.3)$$

so for any loop  $C \subset \mathbb{R}^3$  there exists a surface  $S \subset \mathbb{R}^3$  such that

$$\partial S = C. \quad (1.4)$$

For two loops  $C_1$  and  $C_2$  in  $\mathbb{R}^3$  we can define linking number via the intersection number

$$\text{lk}(C_1, C_2) = C_1 \cdot S, \quad (1.5)$$

where  $S$  is surface, homotopic to a disc, such that  $\partial S = C_2$ . We can use the integral formula for the intersection number

$$C \cdot S = \int_{\mathbb{R}^3} \eta_C \wedge \eta_S. \quad (1.6)$$

The 1-form  $\eta_C$  is defined via

$$\int_C \omega = \int_{\mathbb{R}^3} \eta_C \wedge \omega, \quad \forall \omega \in \Omega^1(\mathbb{R}^3). \quad (1.7)$$

The Poincare dual form  $\eta_C$  for loop  $C$  is trivial in cohomology

$$[\eta_C] = 0 \in H^2(\mathbb{R}^3) = 0, \quad (1.8)$$

hence we can write it as

$$\eta_C = d\eta_S. \quad (1.9)$$

Formally, the 1-form  $\eta_S$  is given by

$$\eta_S = d^{-1}\eta_C, \quad (1.10)$$

but the inverse of external derivative is not a differential operator on forms. We can use homotopy to construct differential operator version of  $d^{-1}$

$$\eta_C = (1 - \Pi_0)\eta_C = (dK + Kd)\eta_C = dK\eta_C \implies \eta_S = K\eta_C, \quad (1.11)$$

where homotopy

$$K = \int_0^\infty ds \, d^* e^{-s\Delta}. \quad (1.12)$$

1. (10 points) Evaluate the Poincare dual  $\eta_C$  for the closed loop given in parametric form

$$C = \gamma : S^1 \rightarrow \mathbb{R}^3 : t \mapsto \gamma(t). \quad (1.13)$$

2. (10 points) Evaluate the  $d^*\eta_C$  for  $\eta_C$  from part 1.
3. (10 points) Evaluate the Hodge-Laplacian  $\Delta = dd^* + d^*d$  for 1-form on  $\mathbb{R}^3$  written in components

$$\omega = \sum_{i=1}^3 \omega_i(x) dx^i \quad (1.14)$$

*Hint:* You can use either of 3 methods:

- Straightforward evaluation in terms of differential forms.
  - In HW 2 we related  $d$  and  $d^*$  on  $\mathbb{R}^3$  to divergence, gradient and curl, so we can use vector calculus in 3d to evaluate Hodge-Laplacian.
  - In HW 2 we related  $d$  and  $d^*$  on  $\mathbb{R}^3$  to certain differential operators on  $\Pi T\mathbb{R}^3 = \mathbb{R}^{3|3}$ , so we can use differential operators representation to evaluate Hodge-Laplacian.
4. (10 points) Use your answers for  $d^*\eta_C$  and Hodge-Laplacian action on forms to evaluate the  $\eta_S = K\eta_C$ . Use integral representation for  $\delta$ -function in  $\eta_C$

$$\delta^3(x) = \prod_{k=1}^3 \delta(x^k) = \frac{1}{(2\pi)^3} \int_{\mathbb{R}^3} d^3p \, e^{i \sum p_k x^k} \quad (1.15)$$

5. (10 points) Evaluate the intersection number

$$C \cdot S = \int_{\mathbb{R}^3} \eta_C \wedge \eta_S \quad (1.16)$$

by performing the  $\mathbb{R}^3$  integral first,  $p$ -integral second and  $s$ -integral last.

## 2 $d = 0$ $N = 4$ supersymmetry

Let  $W(z)$  be the holomorphic ( $\bar{z}$ -independent) polynomial with complex coefficients. Let us consider a integral over the complex plane

$$Z_W = \frac{1}{2\pi i} \int_{\mathbb{C}} dz d\bar{z} e^{-W(z)\overline{W(z)}} \partial_z W(z) \overline{\partial_z W(z)} \quad (2.1)$$

We can introduce four Grassmann-odd variables  $\psi^z, \bar{\psi}^z, \psi^{\bar{z}}$  and  $\bar{\psi}^{\bar{z}}$  to rewrite the integral as partition function

$$\begin{aligned} Z_W &= \frac{1}{2\pi i} \int d^4\psi dz d\bar{z} e^{-S(z,\psi)}, \\ S(z,\psi) &= W(z)\overline{W(z)} - \partial_z W \psi^z \bar{\psi}^z - \overline{\partial_z W(z)} \psi^{\bar{z}} \bar{\psi}^{\bar{z}} \end{aligned} \quad (2.2)$$

1. (10 points) Show that the action  $S(\psi, z)$  is invariant under the four Grassmann-odd symmetries.

*Hint:* You can use one copy of  $N = 2$  transformation for  $z$  and another copy for  $\bar{z}$  so that

$$\begin{aligned} \delta z &= \epsilon^z \bar{\psi}^z + \bar{\epsilon}^z \psi^z \\ \delta \bar{z} &= \epsilon^{\bar{z}} \bar{\psi}^{\bar{z}} + \bar{\epsilon}^{\bar{z}} \psi^{\bar{z}} \end{aligned} \quad (2.3)$$

2. (10 points) Show that the integration measure  $d^4\psi dz d\bar{z}$  also invariant under the same symmetries (to linear order in symmetry parameters)

3. (10 points) An integral of symmetry-invariant action with respect to symmetry-invariant measure localizes. What are the localization points for integral (2.2)?

4. (10 points) Evaluate the (Gaussian) integrals around the localization points.

5. (10 points) The partition function  $Z_W$  is invariant under the shift  $W(z) \rightarrow W(z) + tf(z)$

with  $f(z)$  being holomorphic polynomial with complex coefficients of degree less than degree of  $W$ . We can use such invariance to deform the  $W$  of degree  $n$  to the standard form

$$W(z) = az^n, \quad a \in \mathbb{C} \setminus 0 \quad (2.4)$$

Assuming  $W(z)$  has form (2.4) evaluate the integral

$$Z_W = \frac{1}{2\pi i} \int_{\mathbb{C}} dz d\bar{z} e^{-W(z)\overline{W(z)}} \partial_z W(z) \overline{\partial_z W(z)} \quad (2.5)$$

*Hint:* It might be helpful to use radial coordinates  $r, \phi$  on  $\mathbb{C}$

$$z = re^{i\phi}, \quad \bar{z} = re^{-i\phi}, \quad r \in [0, +\infty), \quad \phi \in [0, 2\pi). \quad (2.6)$$