

SUPERSYMMETRIC QUANTUM MECHANICS AND MORSE THEORY

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HOMEWORK 2: DIFFERENTIAL FORMS

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1 Differential forms

Let us consider Euclidean space \mathbb{R}^3 with coordinates x^1, x^2, x^3 endowed with Riemann metric metric $g(\partial_i, \partial_j) = g_{ij} = \delta_{ij}$. For a vector field

$$\vec{v} = v^i(x)\partial_i \in Vect(\mathbb{R}^3) \quad (1.1)$$

we can associate two differential forms

$$\omega_{\vec{v}}^2 = \frac{1}{2} \sum_{ijk} \epsilon_{ijk} v^i dx^j \wedge dx^k \in \Omega^2(\mathbb{R}^3), \quad \omega_{\vec{v}}^1 = \sum_{ij} g_{ij} v^j dx^i \in \Omega^1(\mathbb{R}^3). \quad (1.2)$$

Show that the following relations are identities:

1. (5 points) Hodge dual

$$\omega_{\vec{v}}^2 = *\omega_{\vec{v}}^1, \quad \omega_{\vec{v}}^1 = *\omega_{\vec{v}}^2. \quad (1.3)$$

2. (5 points) Gradient formula

$$\omega_{\vec{\nabla}f}^1 = df. \quad (1.4)$$

3. (5 points) Curl formula

$$d\omega_{\vec{v}}^1 = \omega_{\vec{\nabla} \times \vec{v}}^2. \quad (1.5)$$

4. (5 points) Divergence

$$*d\omega_{\vec{v}}^2 = \vec{\nabla} \cdot \vec{v}. \quad (1.6)$$

5. (10 points) Maxwell equations: Let us consider 4d Minkowski space $\mathbb{R}^{3,1} = \mathbb{R}^3 \times \mathbb{R}$ with coordinates $t = x^0, x^1, x^2, x^3$ and metric

$$g(\partial_i, \partial_j) = g_{ij} = \delta_{ij}, \quad g(\partial_i, \partial_0) = 0, \quad g(\partial_0, \partial_0) = -1 \quad (1.7)$$

For electromagnetic field in terms of electric field 3-vector $E^i(x)$ and magnetic field 3-vector $B^i(x)$ we can associate a 2-form

$$F = g_{ij} E^i dx^0 \wedge dx^j + \epsilon_{ijk} B^i dx^j \wedge dx^k = dx^0 \wedge \omega_E^1 + \omega_B^2 \quad (1.8)$$

Show that Maxwell equations in vacuum can be written in the form

$$dF = d * F = 0. \quad (1.9)$$

with $*$ being Hodge dual in four-dimensional space. We assumed the choice of units where the speed of light $c = 1$.

2 Differential forms as functions on supermanifold

In class we established a relation between differential forms and functions on supergaminold

$$F : \Omega^*(\mathbb{R}^n) \rightarrow C^\infty(\Pi T\mathbb{R}^n) : \omega \rightarrow F_\omega. \quad (2.1)$$

Let us choose x^1, \dots, x^n and ψ^1, \dots, ψ^n as even and odd coordinates on $\Pi T\mathbb{R}^n$. The differential p -form

$$\omega = \sum_{i_1 \dots i_p} \frac{1}{p!} \omega_{i_1 \dots i_p} dx^{i_1} \dots dx^{i_p} \in \Omega^p(\mathbb{R}^n) \quad (2.2)$$

maps to a function

$$F_\omega(x, \psi) = \sum_{i_1 \dots i_p} \frac{1}{p!} \omega_{i_1 \dots i_p} \psi^{i_1} \dots \psi^{i_p} \in C^\infty(\Pi T\mathbb{R}^n). \quad (2.3)$$

1. (5 points) Show that the differential d and vector field substitution ι_v are first order differential operators on $C^\infty(\Pi T\mathbb{R}^n)$

$$F_{d\omega} = D_d F_\omega, \quad F_{\iota_v \omega} = D_{\iota_v} F_\omega \quad (2.4)$$

and evaluate D_d and D_{ι_v} in (x, ψ) coordinates.

2. (5 points) Using the differential operator representation verify the properties of external derivative d and ι_v

- Differential

$$d^2 = 0. \quad (2.5)$$

- Antysymmetry

$$\iota_v \iota_w \omega = -\iota_w \iota_v \omega. \quad (2.6)$$

- Derivation: Let $\omega \in \Omega^k(\mathbb{R}^n)$ and $\mu \in \Omega^m(\mathbb{R}^n)$

$$d(\omega \wedge \mu) = d\omega \wedge \mu + (-1)^k \omega \wedge d\mu, \quad (2.7)$$

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$$\iota_v(\omega \wedge \mu) = \iota_v\omega \wedge \mu + (-1)^k \omega \wedge \iota_v\mu. \quad (2.8)$$

3. (5 points) We can use Cartan formula to describe Lie derivative via

$$\mathcal{L}_v = \{d, \iota_v\} \quad (2.9)$$

In terms of the functions on $C^\infty(\Pi T\mathbb{R}^n)$ the D_{L_v} is the graded commutator of two first order differential operators, so it is also a first order differential operator

$$D_{L_v} = \{D_d, D_{\iota_v}\} \quad (2.10)$$

Evaluate the D_{L_v} in coordinates x, ψ .

4. (5 points) Using differential operator representation verify the properties of Lie derivative

- Derivation property: Let $\omega \in \Omega^k(\mathbb{R}^n)$ and $\mu \in \Omega^m(\mathbb{R}^n)$

$$\mathcal{L}_v(\omega \wedge \mu) = \mathcal{L}_v\omega \wedge \mu + \omega \wedge \mathcal{L}_v\mu.$$

- Lie algebra :

$$(\mathcal{L}_v\mathcal{L}_w - \mathcal{L}_w\mathcal{L}_v)\omega = \mathcal{L}_{[v,w]}\omega.$$

- Lie algebra representation:

$$(\mathcal{L}_v\iota_w - \iota_w\mathcal{L}_v)\omega = \iota_{[v,w]}\omega.$$

5. (10 points) Let us restrict our attention to the differential forms μ, ω that go to zero at infinity sufficiently fast to the integrals

$$\langle \omega, \mu \rangle = \int_{\mathbb{R}^n} \omega \wedge * \mu \quad (2.11)$$

are finite and we can define the Hodge-dual operators. Verify the Hodge dual of a differential p -form ω as a function on supermanifold $\Pi T\mathbb{R}^n$

$$(-1)^{\frac{1}{2}(n-p)(n+p-1)} F_{*\omega}(\psi, x) = \int_{\mathbb{R}^{0|n}} d^n \eta \ e^{\sum \delta_{ij} \psi^i \eta^j} F_\omega(x, \eta) \quad (2.12)$$

with integration over $d^n\eta$ being Grassmann integration.

Hint: Let α be a monomial expression of Grassmann variables, not depending on η^i then

$$\int_{\mathbb{R}^{0|n}} d^n\eta \alpha\eta^1\dots\eta^n = \alpha. \quad (2.13)$$

6. (10 points) The hodge dual d^* of external derivative d is the second order differential operator D_{d^*} on $C^\infty(\Pi T\mathbb{R}^n)$. Use the Hodge star formula (2.12) to evaluate this operator in x^i, ψ^i coordinates

$$F_{d^*\omega} = D_{d^*}F_\omega. \quad (2.14)$$

3 Laplacian with superpotential

Given a smooth function h on closed smooth manifold M with Riemann metric g we can construct a differential operator

$$d_h = e^{-h}de^h = d + dh\wedge : \Omega^p(M) \rightarrow \Omega^{p+1}(M). \quad (3.1)$$

- (5 points) Show that d_h is the differential.
- (5 points) Show that

$$H^k(\Omega^*(M), d_h) = H_{DR}^k(M). \quad (3.2)$$

- (10 points) Let us for simplicity assume that $M = \mathbb{R}^n$ with Euclidean metric $g_{ij} = \delta_{ij}$ and we restrict our attention to the forms, that vanish sufficiently fast at infinity, so the integrals over \mathbb{R}^n are finite. Describe the dual operator d_h^* defined via

$$\langle d_h\omega, \mu \rangle = \langle \omega, d_h^*\mu \rangle, \quad \mu \in \Omega^k(M) \quad (3.3)$$

in terms of d^* and ι_v .

Hint: It might be helpful to represent the differential forms as a functions on supermanifold $\Pi T\mathbb{R}^n$.

- (10 points) Let us for simplicity assume that $M = \mathbb{R}^n$ with Euclidean metric $g_{ij} = \delta_{ij}$ and we restrict our attention to the forms, that vanish sufficiently fast at infinity, so

the integrals over \mathbb{R}^n are finite. Evaluate the Laplacian

$$\Delta_h = d_h d_h^* + d_h^* d_h \tag{3.4}$$

in the form of differential operator acting on $C^\infty(\Pi T\mathbb{R}^n)$.